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A
S Y S T E M
O F
M E C H A N I C S,

BEING THE
SUBSTANCE OF LECTURES UPON THAT BRANCH OF
NATURAL PHILOSOPHY,

By the Rev. T. PARKINSON, M. A.

FELLOW OF CHRIST'S COLLEGE, CAMBRIDGE.

C A M B R I D G E,

Printed by J. ARCHDEACON Printer to the UNIVERSITY;
For J. & J. MERRILL, and J. DEIGHTON, in Cambridge; T. CADELL, and P. ELMSLY,
in the Strand, B. WHITE, in Fleetstreet, and G. WILKIE, in St. Paul's
Churchyard, London.

MDCCLXXXV.



TO THE
TUTORS, AND OTHER MEMBERS OF THE UNIVERSITY,
INTERESTED IN THE
ADVANCEMENT OF PHILOSOPHICAL KNOWLEDGE,
THE FOLLOWING WORK,
FOR THE USE OF THE ACADEMIC STUDENT,
IS,
WITH GREAT DEFERENCE AND RESPECT,
INSCRIBED
BY THEIR VERY HUMBLE SERVANT,
THOMAS PARKINSON.

CHRIST'S COLLEGE,
NOV. 25, 1784.

TO THE

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THE want of systematic treatises on mechanics and hydrostatics hath long been considered as equally troublesome to the tutor, and discouraging to his pupils, and first induced a desire to facilitate the attainment of these branches of science by a collection, and methodical arrangement, of their scattered parts. The present coincides, very nearly, with the propositions, or heads of lectures, upon this subject, adopted by the generality of tutors in the university, and is recommended by an obvious connexion and regular order of dependancy. In all natural science, the analytic method of reasoning necessarily precedes the synthetic, and the converse of this order would terminate in uncertainty and chimera. The existence and delineation of those properties of matter, which are conceived to generate the phenomena of pressure and motion, are therefore concisely and generally premised in the first six chapters, and supply data, derived from experience and uncontroverted facts, for synthetic demonstrations. The general properties of motion are described in the chapter upon solidity, because the clearest conceptions of motion are derived from impulse, which arises from solidity; and the laws of motion are introduced in the chapter upon the inertia of matter, because they are inseparably connected with

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with it. The composition and resolution of forces are immediately subsequent; because every proposition upon this subject is deducible from the second law of motion, and may be esteemed a corollary to it. The connection of the other parts of the treatise is, I trust, equally natural and obvious. To have divested this elementary branch of philosophy of that dry, uninviting aspect, so frequently and feelingly lamented by the young pupil, was certainly most desirable; but the same complaint is applicable to the elements of all science, and could not, I believe, be removed without an entire change of the formality of precise definitions and propositions into a more undefined and less scientific form of composition, nor consequently without sacrificing real advantage to less material amusement. The following performance, therefore, claims little more than the inferior merit of facilitating the progress of the student by a selection, from the works of others, which may supersede the necessity of applying to a multitude of books, and an arrangement coinciding with his lectures. The whole is written in the same language for the sake of uniformity; demonstrations are dilated or contracted as was deemed expedient; and sometimes, though as seldom as possible, new proofs are given. It is no easy task to compose a system equally accommodated to every description of understandings, and I dare not hope to have accomplished it: for the benefit of those, who may be dissatisfied with the discussion of any subject here; accurate references to the page or chapter of the best writers where it is treated, are always inserted. To be of service to the ignorant and uninformed was the chief motive for undertaking this work, and the sole object of attention in the execution of it: this last consideration may serve to obviate some objections, which naturally will occur to more enlightened readers already conversant with the subject. Lest I appear to presume upon the experience that may be supposed to attend my situation, I beg leave to explain
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myself. When several demonstrations of the same proposition are exhibited, let it not be inferred that one was deemed insufficient to establish its truth and required an auxiliary: when any important subject is treated with prolixity, and extended beyond the limits usually prescribed to an elementary treatise, let it not be attributed to negligence or inadvertency: they are both the effects of design, and I am warranted, by my own experience at least, to assert their utility. It must be remembered, that clear and adequate ideas upon a new subject are not communicated by a transient impression, and that the capacity and comprehension of an uninformed mind are only expanded and improved by repeated exertions and assiduity. Science implies something more than a mere ability to go through a demonstration; and a clear comprehensive knowledge of a question, which may be competent for the solution of problems and dissipation of doubts and objections, is perhaps only to be attained by long reflection, and studiously contemplating it on every side; and these are certainly much assisted by different demonstrations, which place the subject in different points of view, and by tracing its affinity with other truths similar to, or deducible from, it. A single demonstration may easily be too long, and too anxious a desire to be perspicuous may occasion perplexity and confusion; but it is not easy to give too many demonstrations of a fundamental proposition, or too many problems and corollaries resulting from it. This species of prolixity appears to me to be extremely useful, as it may, at least, teach a young student the art of thinking, an art which can only be acquired, and is found indeed to be no easy acquisition. For these reasons, the reader will observe some leading subjects extended beyond the bounds absolutely necessary, and some parts will apparently be superfluous: I shall hold him justified in thinking so, when he shall be so informed on the subject as to pronounce them useless to him.

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I cannot conclude this preface without expressing my sincere thanks to Mr. Fisher and Mr. Vince, of Caius College; to the first for his general theorem of the wedge; to the latter for his friendly revision of the whole work. Had I consulted these friends more frequently, many errors and imperfections, for I believe there are many, might possibly have been avoided. My thanks are also justly due to those gentlemen who have kindly honoured this undertaking with their approbation and support, and I willingly embrace this public opportunity of acknowledging my obligation.



INTRO.

INTRODUCTION

TO THE STUDY OF

NATURAL PHILOSOPHY.

1. **N**ATURAL phenomena, in the widest acceptation of the terms, denote any effects in the material part of the creation addressed to one or more of the senses; and natural philosophy is the history of these phenomena, and an investigation of the causes employed in their production. A phenomenon may itself be a natural cause productive of numberless effects, and each of these may also be a cause of others, &c.; for, of that infinite variety of events observable in the material world, none are induced per saltum, but effect is dependent upon effect in contiguous succession. As matter is totally inactive, and incapable of communicating motion to itself, all its motions, and powers of producing a change of motion, in the various operations of nature, are derivative: but the instruments immediately directing the movements of the several parts of the system elude the inquiry of human ability, and whether any inexplicable effect be owing to the Creator's immediate fiat, or some secondary material power, cannot be known; for the action of a pure spirit upon matter cannot be comprehended: but many subordinate instruments in the government of nature are conspicuous, matter being impressed by its great Creator with several attributes, which appear, and are conceived by some philosophers, to reside in it, ministerial to the continuance of existence and preservation. The rules by which these attributes are directed in their operations, are called natural principles or laws, because in similar circumstances, they are invariably the same; and the attributes themselves are generally called powers or forces,

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from the similitude of their effects to those produced by animal exertions : such are gravity, cohesion, elasticity, magnetism, electricity. The uniform and regular action and utility of some of these natural powers are very observable; but the last powers seem to be still in a state of analysis, and the laws, by which their influence is directed, very imperfectly ascertained. The attraction of cohesion, or of that power by whose influence the minute particles of matter tend to each other at small distances, is the cement which prevents the dispersion of the component parts of matter, and, as far as we are competent to decide, administers to the growth of bodies in the animal, vegetable, and fossil kingdoms. The figures and motions of the great bodies composing the solar system are preserved by an uninterrupted exertion of the attraction of gravity. Nor is the opposite quality of elasticity or repulsion less regular or important. The particles of air are endued with this repulsive power which is essential to the preservation of all bodies contiguous to it, and a diminution of it would proportionably diminish its salutary influence, rendering it noxious to animal and vegetable life; and many hard bodies are obviously possessed of this quality probably to preserve their specific nature by protecting their constituent particles from the effects of attrition, &c. The operations of other attributes of matter, however desultory and accidental apparently, are, upon the fullest information, discovered to be restrained within prescribed limits, whose observance is productive of harmony, and violation of disorder. Plants and animals are always produced from their proper seed; the reflection and refraction of light are effected according to inviolate laws; matter, by its vis inertiae, resists the action of any material power, and when in motion moves, according to determinate rules, &c. To describe the various phenomena exhibited in the production and change of motion in the material world, whether explicable or not, promiscuously, is the business of the natural historian; and to select those which are explicable, and investigate the secondary powers, or qualities, conceived to be resident in matter, by whose instrumentality they appear to be effected, is the province and proper occupation of the natural philosopher.

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2. The historical part of philosophy, or description of natural phenomena, is immediately transcribed from the works of nature; and, from the inadequacy of our ideas of matter, an investigation of the secondary material powers producing them must be derived from the same source. Some phenomena being observed to be invariably coexistent or successive, a connexion is understood, or presumed, to obtain between them; and, the object of philosophy being an evolution of this connexion, the works of nature where only it is discoverable, must be attentively explored, the relation noted, and extended, by a cautious selection of all the circumstances of similitude in other phenomena, to a common principle or law of nature. The following process is therefore adopted by the best philosophers, and is the only certain method of philosophizing.

First, By repeated and accurate observations upon matter, or upon some general similar phenomena in the material world, the existence and qualities of any individual cause are demonstrated from the phenomena evidently and invariably attached to it; the cause of this cause, and the cause of this last, &c. and the effects successively more subordinate and particular resulting from the phenomena, considered as mechanical causes, are then investigated by again exploring the works of nature, and this process is repeated to the limit of human ability. How far the chain of these secondary material powers, either progressive or regressive, from those that are more general to their subordinate effects, and vice versa, is extended, is undiscoverable; but the philosopher's researches are soon limited by the occurrence of bodies inconceivably minute, or removed to immeasurable distances. An examination of the few powers that are known is, however, a rational and not unprofitable employment; for they exhibit proofs of unbounded power, consummate wisdom, and paternal benevolence in the great Creator; and by their connexion with many useful arts, are made subservient to the wants and infirmities of his creatures. *Secondly*, The intensity at a given distance, and law of variation at different distances, of any material power being ascertained by the mensuration of its effects, it is assumed as a principle, and from its influence, all similar phenomena are demonstrated to result by

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mathematical or other scientific methods of reasoning. For many effects are, upon examination, found to bear strong marks of similitude, which to a negligent observer exhibit no likeness; and, these being divested of all adventitious particularities, and ranged under the same cause, a few principles are discovered to pervade the whole system of matter, producing innumerable phenomena. The first of these is called the analytic, and the second the synthetic method of philosophizing.

3. Examples illustrating the analytic method of reasoning.

EXAMP. I. The spherical figure of the earth proved analytically.

P H E N O M E N A.

PHENOM. I. The altitude of either pole of the equator is always equal to the latitude of the observer.

PHENOM. II. In an eclipse of the moon a section of the earth's shadow is always terminated by a circular arc.

Many other phenomena might be adduced in proof of the globular figure of the earth, but it is established by these, because no other afford a solution of them. This figure was impressed at the original formation of the earth, and is preserved by an unremitted exertion of the attraction of gravity; and here the analysis terminates, the cause of gravity being undiscovered.

4. EXAMP. II. The diurnal revolution of the earth proved analytically.

P H E N O M E N A.

PHENOM. I. The sun and most of the planets revolve round their axes, and a similar motion of the earth is analogous to these.

PHENOM. II. The sun, planets and fixed stars appear to revolve uniformly in circles, whose planes are perpendicular to, and centers in, a line passing through the center of the earth, and perform one revolution in twenty-four hours nearly.

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PHENOM. III. The gravity of bodies is least under the equator, and increases as they recede from it.

If this motion of the heavenly bodies were real, they would be retained in their orbits by forces tending to their corresponding centers in the axis of the world (Newt. sect. 2. prop. 2.), which are imaginary points in which no attractive powers reside; and some of these motions could not be effected without the influence of forces infinitely greater than any mechanical powers yet discovered. They are therefore apparent only, and result from the rotation of the earth round an axis; and this is confirmed by phenom. 2. which is easily explicable by, and inexplicable without, such rotation.

5. EXAMP. III. The properties of light investigated analytically.

P H E N O M E N A.

PHENOM. I. If the sun's rays be admitted, through a small aperture, into a dark room, and refracted through a glass prism, an oblong image consisting of several colours will be formed upon a sheet of paper placed at a proper distance.

PHENOM. II. If another prism be placed behind the former, so that their axes be perpendicular, the image will be refracted on one side, of the same length and breadth as before.

PHENOM. III. Any one colour refracted through any number of prisms, always exhibits a circular image of that colour.

If the sun's rays were equally refrangible, its image would be circular, and, if differently refrangible, oblong; and consequently (phenom. 1st and 2d) light is composed of heterogeneous parts differently refrangible. If the sun's rays splitted or dilated by refraction, the image would be circular and increased by every refraction, therefore (phenom. 2d) they do not split or dilate; and (phenom. 3d) the properties of homogeneous rays are innate and unalterable by refraction.

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6. The conclusions thus derived by analysis are analogical, as they result from a comparison between present and past phenomena, and will be just only when these comparisons are repeatedly made, and the similitude of the phenomena perfectly established. In example the 1st. the altitude of the pole being found equal to the latitude of the observer, and a section of the earth's shadow to be a circular arc, and these phenomena being invariably the same in innumerable trials, an assurance that they result from some fixed causes and will never vary, arises and commands our assent.

Philosophy therefore, when real and not fantastical, has for its basis uniform and uncontroverted experience; and hence appears the necessity of adhering to the first rule of philosophizing.

R U L E I.

No more causes of natural events ought to be admitted than are real, and sufficient to explain the phenomena.

7. Of those bodies which are neither concealed by their minuteness, nor remote distance, few can undergo an analytical discussion; and philosophy would be very limited, were it confined to those only which have actually been the objects of experiments; but it is rendered universal by observing the following rules.

R U L E II.

Effects of the same kind are to be ascribed to the same cause.

The same cause occasions the descent of bodies in different places of the earth, respiration in different animals, the sensation of heat from a culinary fire and the sun, solutions of bodies in different menstruums, and of water in air, &c.

R U L E III.

Those qualities of matter which do not vary, and are found in all bodies that admit of experimental examination, ought to be considered as qualities of all bodies in general.

Thus

Thus the attraction of gravity is discovered, by experiments and astronomical observations, to pervade all matter with which we are acquainted: all bodies near the earth's surface, and the moon, gravitate towards the earth; the waters of the sea gravitate towards the moon; the sun, planets, and comets, gravitate towards each other mutually; and of this principle no bodies accessible to experiments can be divested, and therefore, from this rule, it is inferred to be an universal attribute of matter.

9. As much false philosophy hath been disseminated by fanciful hypotheses and mere metaphysical considerations, unsupported by the reality of facts, the only sources of natural science, the utility of the following rule is apparent.

R U L E. IV.

In experimental philosophy, propositions collected from phenomena by induction are to be deemed, notwithstanding contrary hypotheses, either accurately true or very nearly so; until other phenomena occur by which they may be rendered either more accurate, or liable to exception.

10. The existence, quantity at a given distance, and law of variation at different distances, of any natural power being demonstrated by the analytic method of reasoning, the only way by which they can be known; they are assumed as established general principles, and their adequacy to the production of other phenomena, similar to those used in the analysis, is proved synthetically.

Examples illustrating the synthetic method of reasoning.

EXAMP. I. The influence of the attraction of gravity, its quantity and law of variation, being ascertained analytically, it is assumed as an established general principle, and demonstrated geometrically to be an adequate cause of the motions of pendulums, projectiles, precession of the equinoxes, irregularities of the moon's motion, rising and ebbing of the sea, and innumerable other phenomena.

EXAMP.

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EXAMP. II. The weight of the air being ascertained by experiments, it is supposed to be allowed, and affords an easy solution of the phenomena of pumps, syringes, barometers, &c.

EXAMP. III. The equality of the angles of incidence and reflection of light, of the constant ratio obtaining between the sines of incidence and refraction, and the unequal refrangibility of different colours, are discovered by experiments; and, being supposed to be universally true, they afford easy explanations of the figure and magnitude of images formed by reflection and refraction, of the rainbow, &c.

EXAMP. IV. The spherical figure of the earth, its diurnal motion, and obliquity of the ecliptic are discovered by repeated observation, and, being assumed as known principles, they afford an easy solution of the doctrine of the sphere, of the art of dialling, of day and night and their inequalities, of heat and cold, and many other phenomena.

II. Of these two methods of reasoning, the latter may strictly be denominated science; for the existence of any natural power, its magnitude and variation being presupposed, it is demonstrated geometrically to be competent to the production of various effects, and, were this power changed or annihilated, the truth of this reasoning would remain unaltered. Whether this science be chimerical, or accord with actual existence, depends upon the truth of the presupposed data, or analytic conclusions, which admit of various degrees of conviction from bare presumption to certainty. The existence of some undiscovered and still unknown quality of matter, collected only from a few experiments, amounts only to a presumption; but the degree of probability increases with their number, and, by repeated and uniformly concurring trials, establishes at length an evidence as unquestioned as the existence of matter or the senses. That the qualities of matter will remain unchanged and continue always to produce the same effects, depends upon the will of that divine Agent, who ordained that the
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various events in the material world should be effected by the intervention of simple and general material causes, acting with invariable constancy and usefulness, and only suffered to transgress their prescribed bounds according to the dictates of infinite wisdom, power and goodness. Presuming therefore upon the identity of the senses and the uniformity of natural operations, the study of philosophy may be deemed scientific; for, being founded upon the basis of uniform and uncontroverted experience, and geometric demonstration, its certainty will continue with them, and consequently can only be affected or subverted by a change or subversion of the constitution of nature.

12. Of all magnitudes, numbers are the most simple, and their minute differences most easily discriminated, and therefore more susceptible of clear ideas than any other magnitudes; and the relations subsisting between powers or forces, and the changes or velocities produced by their exertions, are said to be known, when reduced to numeral expressions, though the mode of producing these effects be unknown, and perhaps undiscoverable. The next, in degree of simplicity and clearness of comprehension, are lines, surfaces and solids, being permanent and easily compared by juxtaposition. The conception of other magnitudes is more remote and difficult. Forces, velocities, and times, having no permanent representatives like numbers and geometric magnitudes, and disappearing when not actually subsisting, are best understood from their effects; and this study is therefore much simplified by finding magnitudes of easier conception, and more easily measurable and permanent, as numbers, lines and surfaces, whose relation is the same with that of times and velocities, and natural powers. Philosophy may consequently be reduced to an investigation of the relation subsisting between lines, surfaces, &c. which vary as, and may be denominated the representatives of, natural powers and their effects; and a knowledge of ratios, or of the rules by which these relations are increased and diminished, is a necessary lemma to this study and here premised.

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RATIOS.

R A T I O S.

13. DEF. *W*HATEVER is capable of increase or decrease, is called magnitude or quantity, as numbers, lines, velocities, forces, &c.

14. DEF. Ratio is the mutual relation of two magnitudes of the same kind in respect of their greatness or smallness; the first is called the antecedent and the second the consequent of the ratio.

15. Ratio is therefore a comparison, which always implies a similitude, and is limited by the definition to the relative greatness or smallness of the quantities compared. The magnitude of the ratio of equal quantities is equal to nothing; for the existence of ratios results from the inequality of the quantities compared, though not measured by it; and it is evident, that the magnitude of a ratio may increase or decrease through every stage of assignable quantity.

E X A M P L E S.

FIG. I. EXAMP. I. If the distance of E from B be equal to nothing, the diameter QB , and chord QE coincide, and their ratio is a ratio of equality, and its magnitude equal to nothing; but as the arc BE increases from nothing to a semicircle, the magnitude of this ratio increases from nothing through every stage of assignable magnitude, and, when E and Q coincide, is unassignably great.

EXAMP. II. The ratio of the tangent TA to the sine SN of an arc SA , or of the secant CT to the radius CS , is a ratio of equality, and its magnitude equal to nothing, when the arc vanishes, and S and A coincide. But as the arc SA increases from nothing to a quadrant, the magnitude of this ratio increases from nothing through

through every stage of assignable magnitude, and, when S and Q coincide, is unassignably great. If the magnitude of this ratio (when e. g. $SA = 30^\circ$) be expounded by any finite line L , and twice this ratio by twice this line, &c. when S and A coincide, L is equal to nothing; and as S recedes from A , L increases and becomes infinitely great when S arrives at Q .

Ratios therefore are magnitudes, and like all other magnitudes the object of ratios, and capable of addition, subtraction, multiplication and division. They are positive or negative, according to their different effects in addition. The sum of two positive, or of two negative ratios, will constitute a positive, or negative ratio, greater than either of them; and, according as a positive ratio is greater or less than a negative one, their sum will be positive or negative; and, if they be equal, their sum is nothing. The magnitudes of the ratio of $L : M$ and of $M : L$, or of $3 : 2$ and of $2 : 3$, are clearly equal, but of different denominations; and any third ratio, of the same denomination with the ratio of $L : M$, will be equally increased by the addition of the ratio of $L : M$ and subtraction of the ratio of $M : L$. If the ratio obtaining between L and M , of which L is the greatest, be called positive, and represented by any line or number a , twice this ratio will be represented by twice a , and n times this ratio by $n \times a$. The ratio also of $M : L$ will be represented by $-a$, and $n \times$ the ratio of $M : L$ by $n \times -a$; and a line or number, representing any third affirmative ratio, will be equally increased by the addition of $n \times a$ and subtraction of $n \times -a$. As magnitudes are only measurable by magnitudes, sui generis, a line by a line, a surface by a surface, and a ratio by a ratio: to ascertain the quantity of any ratio, some more simple ratio may be used as a criterion. Thus logarithms are a series of numbers expressing the relation which subsists between any given ratio, considered as a criterion, and all other ratios with which it is compared.

16. DEF. *In a series of magnitudes of the same kind, either increasing or decreasing, the ratio of the extremes is said to be compounded of the ratios of the intermediate terms.*

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FIG. II. In the series of magnitudes A, B, C, D , &c. the ratio of the first to the last, or of $A : D$, is said to be compounded of the ratios of $A : B, B : C, C : D$; and any ratio, as that of $PQ : PR$, is said to be resolvable into, and equal to, the ratios of $PQ : Pa, Pa : Pb, Pb : Pc \dots PL : PR$; for the ratio of $PQ : PR$ is a real magnitude, and like other magnitudes divisible into its component parts, which are the ratios of $PQ : Pa, Pa : Pb, Pb : Pc$, &c.: therefore the sum of any number of continued ratios, where the antecedent of any ratio is the consequent of the preceding, is equal to the ratio of the first and last terms.

FIG. II. 17. Cor. When any ratio, as that of $PQ : PR$, is to be divided into any other ratios by an arbitrary insertion of other quantities Pa, Pb , &c.; these are not necessarily intermediate, or contained between PQ and PR , the ratio of $PQ : PR$ being equal to the ratios of $PQ : Pv$ and of $Pv : PR$. For the ratio of $PQ : Pv$ is equal to the ratios of $PQ : PR$ and of $PR : Pv$ (16), and consequently the ratio of $PQ : PR$ is equal to the ratio of $PQ : Pv$ diminished by the ratio of $PR : Pv$, or added to the ratio of $Pv : PR$ (15).

ADDITION OF RATIOS.

18. PROP. To add the ratios of $A : B, C : D, E : F$, &c. together.

FIG. III. When the two ratios of $A : B$ and $C : D$ are to be added, let A and C, B and D be respectively the sides of two rectangular parallelograms X and Z ; and (Euclid. VI. 23.) the ratios of $A : B$ and $C : D$, when compounded, are equal to the ratio of $X : Z$, or of $A \times C : B \times D$, these quantities being respectively equal to X and Z . When the ratios of $A : B, C : D, E : F$ are to be added, the sum of the two first is equal to the ratio of $AC : BD$ by the process above; and, by making $A \times C$ and $E, B \times D$ and F , respectively the sides of two rectangles, it appears, by the same process, that the sum of the ratios of $AC : BD$ and of $E : F$ is equal to the

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the ratio of $ACE : BDF$. Whatever be the number of ratios to be added the process is similar, and gives the following

R U L E.

Multiply the antecedents together for a new antecedent, and the consequents for a new consequent.

19. EXAMP. I. The sum of the ratios of $1:2, 3:4, 5:6, 7:8$ is equal to the ratio of $1 \times 3 \times 5 \times 7 : 2 \times 4 \times 6 \times 8$, or of $105 : 384$. This also appears from the following analogies and article (16); for (38)

$$\begin{aligned} 1 : 2 &:: 105 : 210 \\ 3 : 4 &:: 210 : 280 \\ 5 : 6 &:: 280 : 336 \\ 7 : 8 &:: 336 : 384. \end{aligned}$$

The sum of the ratios of $1:2, 3:4, 5:6, 7:8$, is equal to the sum of the ratios of $105:210, 210:280, 280:336, 336:384$, or (16) to the ratio of $105:384$, as above.

20. EXAMP. II. When two bodies move with uniform motions, philosophical writers say, that the spaces described S and s are to each other in a ratio compounded of the velocities and times; the meaning of which is, that the ratio of the spaces is equal to the sum of the ratio of the velocities when changed, and the ratio of the times when changed, or that the ratio of $S:s$ is equal to the two ratios of $V:v$ and of $T:t$, or to the ratio of $V \times T : v \times t$, supposing v and t to represent any corresponding values of the velocity and time.

21. Cor. 1. If there be any number of quantities A, B, C, D , &c. of which $A:B::R:r$ and $B:C::S:s$ and $C:D::T:t$, A will be to D as $RST:rst$; for $A:D$ in a ratio compounded of the ratios of $A:B, B:C, C:D$ (16), or their equals $R:r, S:s, T:t$, or as $RST:rst$ (18).

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S C H O L I U M.

22. When the terms of any ratios are forces, times, velocities, &c. lines or numbers are supposed to be taken whose ratios are the same with them, and rectangular parallelograms whose sides are these lines, or the products of these numbers actually multiplied together, are always implied in the multiplication of such magnitudes. The sum of the ratios of $L:M$ (denoting forces), of $N:P$ (denoting velocities), of $Q:R$ (denoting spaces), and of $S:T$ (denoting times), is equal to the ratio of $L \times N \times Q \times S : M \times P \times R \times T$, which signify two products of numbers whose factors are as $L:M, N:P, Q:R, S:T$.

SUBTRACTION OF RATIOS.

23. PROP. *To subtract the ratio of $C:D$ from that of $A:B$.*

Let the ratio of $A:B$ be equal to the ratios of $C:D$ and $x:y$; and (18) $A:B :: xC:yD$, and $x:y :: \frac{A}{C} : \frac{B}{D}$ (Euc.v.4.) $:: AD:BC$ (Euc.v.10.); from hence we have the following

R U L E S.

RULE I. *Divide the antecedent of the subtrahend by the antecedent of the ratio to be subtracted for a new antecedent, and the consequent by the consequent for a new consequent.*

Or, II. *Invert the terms of the ratio to be subtracted, and proceed as in addition.*

24. EXAMP. I. The ratio of $6:5$ subtracted from the ratio of $3:2$ is equal to the ratio of $\frac{3}{6} : \frac{2}{5}$ or of $15:12$, which is thus confirmed. The ratio of $3:2$ is equal to the ratio of $6:4$, which is equal to the ratio of $6:5$ and of $5:4$ (16); and if the ratio of

$6:5$

6 : 5 be taken away, the remainder is the ratio of 5 : 4 or of 15 : 12 (Euc. V. 15.).

25. EXAMP. II. The ratio of 2 : 3 diminished by the ratio of 4 : 5 is equal to the ratio of $\frac{2}{4} : \frac{3}{5}$ or of 10 : 12, which is thus confirmed. The ratio of 2 : 3 is equal to the ratio of 4 : 6 (Euc. V. 15.), that is, of 4 : 5 and 5 : 6 (16); and, if the ratio of 4 : 5 be taken away, the remainder is the ratio of 5 : 6 or of 10 : 12 (Euc. V. 15.).

MULTIPLICATION OF RATIOS.

26. PROP. *To multiply the ratio of A : B by any number m.*

The ratio of $A : B$ added to itself is equal to the ratio of $A^2 : B^2$, and this added to the ratio of $A : B$ is equal to the ratio of $A^3 : B^3$, and this repeated m times will clearly give the ratio of $A^m : B^m$ (18), and we have this rule.

R U L E.

Involve the terms of the ratio to a dimension equal to the multiplier.

27. EXAMP. I. Four times the ratio of $A : B$ is equal to the ratio of $A^4 : B^4$, which is confirmed by the following process; for (38)

$$\begin{aligned} A : B &:: A^4 : A^3 B \\ A : B &:: A^3 B : A^2 B^2 \\ A : B &:: A^2 B^2 : A B^3 \\ A : B &:: A B^3 : B^4. \end{aligned}$$

Therefore four times the ratio of $A : B$ is equal to the ratios of $A^4 : A^3 B$, $A^3 B : A^2 B^2$, $A^2 B^2 : A B^3$, $A B^3 : B^4$, or the ratio of $A^4 : B^4$ (16).

28. EXAMP.

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28. EXAMP. II. Five times the ratio of $2:3$ is equal to the ratio of $2^5:3^5$ or $32:243$, which also appears from the process above and (16). For (38)

$$\begin{aligned} 2:3 &:: 32:48 \\ 2:3 &:: 48:72 \\ 2:3 &:: 72:108 \\ 2:3 &:: 108:162 \\ 2:3 &:: 162:243. \end{aligned}$$

Therefore the ratio of $2:3$ multiplied into five is equal to the ratios of $32:48$, $48:72$, $72:108$, $108:162$, $162:243$; or to the ratio of $32:243$ (16).

DIVISION OF RATIOS.

29. PROP. To divide the ratio of $A:B$ by any number m .

Let $A = x^m$ and $B = y^m$; and the ratio of $A:B$ is equal to the ratio of $x^m:y^m$ or to m times the ratio of $x:y$ (18), which is therefore $\frac{1}{m}$ part of the ratio of $A:B$, and equal to the ratio of $A^{\frac{1}{m}}:B^{\frac{1}{m}}$; because $x = A^{\frac{1}{m}}$ and $y = B^{\frac{1}{m}}$. Whence we have the following

R U L E.

Extract that root of the terms of the ratio which is expressed by the divisor.

30. EXAMP. The ratio of $32:162$ divided by four is equal to the ratio of $\sqrt[4]{32}:\sqrt[4]{162}$ or of $2:3$. And this is confirmed by the process used in example (28), where the ratio of $32:162$ is equal to the ratio of $2:3 \times 4$, and consequently the ratio of $2:3$ is $\frac{1}{4}$ th of the ratio of $32:162$.

31. PROP. When the difference of two magnitudes is very small compared with the magnitudes themselves, their ratio is multiplied or divided

divided by any number m , by increasing or diminishing their difference m times.

DEM. Let the two magnitudes be A and $A \pm y$, whose difference y is very small compared with A ; and m times the ratio of $A : A \pm y$ is equal to the ratio of $A^m : \overline{A \pm y}^m$, or of $A^m : A^m \pm m A^{m-1} y + m \cdot \frac{m-1}{2} A^{m-2} y^2$, &c. or (dividing the antecedent and consequent by A^{m-1}) to the ratio of $A : A \pm m y$; because the terms involving $\frac{y^2}{A}$, $\frac{y^3}{A^2}$, &c. are evanescent compared with the two first, and may be neglected.

By a similar process $\frac{1}{m}$ th part of the ratio of $A : A \pm y$ appears to be equal to the ratio of $A : A \pm \frac{y}{m}$. Q. E. D.

32. EXAMP. Twice the ratio of $11 : 10$ is equal to the ratio of $\overline{11}^2 : \overline{10}^2$ or of $121 : 100$; and, according to this proposition, it is the ratio of $11 : 9$, or of $121 : 99$.

Twice the ratio of $101 : 100$ is equal to the ratio of $\overline{101}^2 : \overline{100}^2$ or of $10201 : 10000$; and according to this proposition it is equal to the ratio of $101 : 99$ or of $10201 : 9999$, which is nearly equal to the former.

33. EXAMP. II. A half of the ratio of $101 : 100$ is equal to the ratio of $\overline{101}^{\frac{1}{2}} : \overline{100}^{\frac{1}{2}}$, or of $10.049 : 10$ nearly; and by this proposition it is equal to the ratio of $100^{\frac{1}{2}} : 100$, or of $201 : 200$, or of $10.05 : 10$.

METHODS OF COMPARING RATIOS.

34. DEF. Proportion is an equality of ratios. When the ratios of $A : B$ and $C : D$ are equal, they are said to be proportional, and usually written thus $A : B :: C : D$, or A is to B as C to D .

C

35. Cor.

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35. Cor. If A, B, C, D be proportional, and $A = \frac{1}{4}B$, or $\frac{m}{n} \times B$, C will be equal to $\frac{1}{4}D$, or to $\frac{m}{n} \times D$; or if these magnitudes be incommensurate, and A be greater, or less, than any part or parts of B , C will be greater, or less, than the same part or parts of D . If A be contained between $\frac{m}{n} \times B$ and $\frac{m+1}{n} \times B$, (m and n being any numbers whatever,) C will also be contained between $\frac{m}{n} \times D$ and $\frac{m+1}{n} \times D$. This is evident from the definitions of ratios and proportion; for if A were contained between $\frac{18}{17} \times B$ and $\frac{19}{17} \times B$, or $\frac{m}{n} \times B$ and $\frac{m+1}{n} \times B$, and C were not contained between $\frac{18}{17} \times D$ and $\frac{19}{17} \times D$, or $\frac{m}{n} \times D$ and $\frac{m+1}{n} \times D$, A 's magnitude compared with B 's, would not be equal to C 's magnitude compared with D 's, or the ratios of $A:B$ and $C:D$ would not be equal, and they would not be proportional.

36. PROP. If A cannot be greater than, equal to, or less than, any part or parts of B , but at the same time C is greater than, equal to, or less than, the same part or parts of D , they will be proportional, or $A:B :: C:D$.

DEM. Let $A:B :: E:D$; and if $A = \frac{m}{n} \times B$, $E = \frac{m}{n} \times D$ (35), and, from the hypothesis, $C = \frac{m}{n} \times D$: therefore $E = C$, and $A:B :: C:D$. Let these magnitudes be incommensurate, and if A be contained between $\frac{m}{n} \times B$ and $\frac{m+1}{n} \times B$, E will be contained between $\frac{m}{n} \times D$ and $\frac{m+1}{n} \times D$ (35), and, from the supposition, C will

C will be contained between $\frac{m}{n} \times D$ and $\frac{m+1}{n} \times D$. The difference therefore between C and E is between $\frac{m}{n} \times D$ and $\frac{m+1}{n} \times D$, and consequently is not greater than $\frac{D}{n}$; and because this is true whatever be the magnitude of the number n , which may be unassignably great and $\frac{D}{n} = 0$, C will be equal to E , and $A : B :: C : D$. Q. E. D.

37. Cor. 1. If four magnitudes A, B, C, D be proportional, the products of the extreme, and middle, terms are equal, or $A \times D = B \times C$. For let $A = \frac{m}{n} \times B$, and therefore $C = \frac{m}{n} \times D$ (35): consequently $A \times D = \frac{m}{n} \times B \times D$, and $B \times C = B \times \frac{m}{n} \times D = A \times D$. If they be incommensurate, let A be greater than $\frac{m}{n} \times B$, and less than $\frac{m+1}{n} \times B$, and C will be contained between $\frac{m}{n} \times D$ and $\frac{m+1}{n} \times D$ (35); therefore $A \times D$ is contained between $D \times \frac{m}{n} \times B$ and $D \times \frac{m+1}{n} \times B$, and $C \times B$ is contained between $B \times \frac{m}{n} \times D$, and $B \times \frac{m+1}{n} \times D$. The difference of these products is therefore not greater than $\frac{DB}{n}$, which, because n may be taken unassignably great, $= 0$; and consequently $A \times D = B \times C$.

38. Cor. 2. If two products be equal, their factors are proportional. Let $A \times D = B \times C$; and if $A = \frac{m}{n} \times B$, then $C = \frac{A \times D}{B}$
= (by

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$=$ (by substituting A 's value) $\frac{m}{n} \times D$; or if A be less than $\frac{m+1}{n} \times B$, but greater than $\frac{m}{n} \times B$; then C , being equal to $\frac{A \times D}{B}$, will be less than $\frac{m+1}{n} \times D$, and greater than $\frac{m}{n} \times D$; consequently (36) $A : B :: C : D$. In the same manner it may be proved, that $A : C :: B : D$, if they be similar; or that $A : BC :: 1 : D$, and $B : A :: D : C$, or $B : A \times D :: 1 : C$, or $A : B :: C : \frac{BC}{A}$. Any equation is therefore resolvable into a proportion by so arranging the terms, that the rectangle of the extreme terms may be one side of the equation, and that of the mean terms the other.

39. Cor. 3. Ratios, which are equal to any ratio, are equal to each other. Let $A : B :: C : D$ and $C : D :: E : F$; and if A be equal to $\frac{m}{n} \times B$, or contained between $\frac{m}{n} \times B$ and $\frac{m+1}{n} \times B$, C will be equal to $\frac{m}{n} \times D$, or contained between $\frac{m}{n} \times D$ and $\frac{m+1}{n} \times D$ (35), and also E will be equal to $\frac{m}{n} \times F$, or contained between $\frac{m}{n} \times F$ and $\frac{m+1}{n} \times F$ (35); therefore A and E are either equal to, or contained between, the same parts of B and F respectively, and $A : B :: E : F$ (36).

40. Cor. 4. If $A : B :: C : D$ they will be proportional inversely, or $B : A :: D : C$. Let A be equal to $\frac{m}{n} \times B$ or contained between $\frac{m}{n} \times B$ and $\frac{m+1}{n} \times B$, and C will be equal to $\frac{m}{n} \times D$, or contained between $\frac{m}{n} \times D$ and $\frac{m+1}{n} \times D$ (35); therefore $\frac{n}{m} \times A$ is equal to

or

or greater than, and $\frac{n}{m+1} \times A$ is less, than B , or B is contained between $\frac{n}{m+1} \times A$ and $\frac{n}{m} \times A$, or equal to $\frac{n}{m} \times A$; and for the same reason D is contained between $\frac{n}{m+1} \times D$ and $\frac{n}{m} \times D$, or equal to $\frac{n}{m} \times C$; therefore (36) $B : A :: D : C$.

41. Cor. 5. Magnitudes are proportional to their equimultiples or equal parts. Let $A = \frac{m}{n} \times B$, and $2A$ will be equal to $\frac{m}{n} \times 2B$, and $\frac{1}{2}A = \frac{m}{n} \times \frac{1}{2}B$; therefore (36) $A : B :: 2A : 2B :: \frac{1}{2}A : \frac{1}{2}B$; or let A be contained between $\frac{m}{n} \times B$ and $\frac{m+1}{n} \times B$, and $2A$ will be contained between $\frac{m}{n} \times 2B$ and $\frac{m+1}{n} \times 2B$; and $\frac{1}{2}A$ between $\frac{m}{n} \times \frac{1}{2}B$ and $\frac{m+1}{n} \times \frac{1}{2}B$; therefore (36) $A : B :: 2A : 2B :: \frac{1}{2}A : \frac{1}{2}B$, &c.

42. Cor. 6. If $A : B :: C : D$, $A^m : B^m :: C^m : D^m$ and $A^{\frac{1}{m}} : B^{\frac{1}{m}} :: C^{\frac{1}{m}} : D^{\frac{1}{m}}$. For m times, or an m^{th} part of, the ratio of $A : B$ is equal to m times, or an m^{th} part of, that of $C : D$, because these ratios are equal to each other; therefore (34) $A^m : B^m :: C^m : D^m$; and $A^{\frac{1}{m}} : B^{\frac{1}{m}} :: C^{\frac{1}{m}} : D^{\frac{1}{m}}$.

43. Cor. 7. If $A : B :: C : D$; then $A+B : B :: C+D : D$. Let A be equal to $\frac{m}{n} \times B$, or contained between $\frac{m}{n} \times B$ and $\frac{m+1}{n} \times B$, and (35) C will be equal to $\frac{m}{n} \times D$, or contained between $\frac{m}{n} \times D$ and

and $\frac{m+1}{n} \times D$; therefore adding B to both, $A+B$ will be equal to $\frac{m+n}{n} \times B$, or greater than $\frac{m+n}{n} \times B$ and less than $\frac{m+1+n}{n} \times B$; and $C+D$ will be equal to $\frac{m+n}{n} \times D$, or greater than this, and less than $\frac{m+1+n}{n} \times D$; and consequently (36) $A+B : B :: C+D : D$. In the same manner it may be proved, that $A-B : B :: C-D : D$; and $A : B :: A \pm C : B \pm D$.

44. DEF. *That ratio is said to be greater or less than another, whose antecedent has a greater proportion to its consequent than the antecedent of the other to its consequent.*

Two ratios, whose antecedents, or consequents, are the same, are easily compared: for it is evident that the ratio of 6:3 is greater, and the ratio of 4:3 less, than the ratio of 5:3; and the negative ratio of 3:6 is greater, and of 3:4 less, than the ratio of 3:5.

45. *The first method of comparing unequal ratios is to reduce them to other ratios equal to them with a common antecedent, or consequent.*

EXAMP. To compare the ratios of $A : B$ and of $C : D$. Find a ratio equal to either of them, whose antecedent or consequent is the antecedent or consequent of the other; thus, $A : B :: C : \frac{BC}{A}$ (38), and consequently the ratio of $A : B$ being equal to that of $C : \frac{B \times C}{A}$, is greater or less than the ratio of $C : D$, according as $\frac{BC}{A}$ is less or greater than D .

Let

Let the ratios of $3:5$ and of $6:9$ be compared; and $3:5::6:10$ (10), therefore the ratio of $3:5$, being the same with that of $6:10$, is a greater negative ratio than that of $6:9$.

46. *The second method of comparing ratios, is to divide the antecedents by their respective consequents, and that ratio will be the greatest, the quotient of whose terms is the greatest.*

EXAMP. The ratio of $A:B$, being equal to that of $C:\frac{BC}{A}$ (38), is greater than, equal to, or less than the ratio of $C:D$, according as $\frac{B \times C}{A}$ is less than, equal to, or greater than D , or $\frac{C}{A}$ less than, equal to, or greater than $\frac{D}{B}$, or $\frac{A}{C}$ greater than, equal to, or less than $\frac{B}{D}$. This method is expeditious and true, but let it be remembered that $\frac{A}{B}$ is a number as every quotient is, and does not measure the ratio of $A:B$, nor is the relative magnitude or ratio of the ratios of $A:B$ and $C:D$ equal to the ratio of $\frac{A}{B}:\frac{C}{D}$.

47. *The third method of comparing ratios. If A be contained between $\frac{m}{n} \times B$ and $\frac{m+1}{n} \times B$, and C be contained between $\frac{m}{n} \times D$ and $\frac{m+1}{n} \times D$, the ratios of $A:B$ and $C:D$ are equal: if therefore A be greater than $\frac{m}{n} \times B$ and C be not greater than $\frac{m}{n} \times D$, or if nA be greater than mB , but nC not greater than mD , the ratio of $A:B$ is greater than that of $C:D$.*

EXAMP.

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EXAMP. Let the ratios to be compared be $7:5$ and $4:3$. Multiply the first and third by 3, and the second and fourth by 4, and the resulting numbers are 21, 20, 12, 12; and the first multiple being greater than the second, but the third not greater than the fourth, 7 has to 5 a greater ratio than $4:3$.

48. DEF. Any magnitude A is said to be, or vary, directly as another B , when A is to any new value of A as B to a corresponding new value of B , that is, when $A:a::B:b$, a and b being corresponding new values of A and B .

FIG.IV. EXAMP. I. If the line mn move parallel to itself, and its extremities m and n be always in the lines LP , LQ ; Lm varies as mn . For $LA:La$ (a new value of Lm) $:: mn:AB$ (a corresponding new value of mn).

EXAMP. II. The area Lmn varies directly as the square of Lm or mn ; for $Lmn:LAB$ (a new value of Lmn) $:: mn^2:AB^2$ (the square of a new value of mn).

49. DEF. Any magnitude A is said to be, or vary, inversely as another B , when A is to a new value of A , as a cotemporary new value of B is to B , or when $A:a::b:B$.

EXAMP. I. In the lever the power is inversely as the perpendicular let fall from the center of motion upon its direction; for let F and f be two powers in equilibrio, P and p the perpendiculars let fall from the center of motion upon their directions; and $F:f$ (a new value of the power) $:: p$ (the perpendicular let fall upon f 's direction) $: P$.

EXAMP.

EXAMP. II. If a given space be described in different times T and t , with unequal uniform velocities V and v , the velocities and times are inversely as each other; for $V : v$ (a new value of the velocity) $:: t$ (a new value of the time corresponding to v) $: T$.

EXAMP. III. If two variable right lines AB , AC form the equal FIG.V. rectangles AD , EF , they are inversely as each other; for (Euc. VI. 14.) $AB : CF$ (a new value of AB) $:: CE$ (a new value of AC) $: AC$.

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50. When any quantity A varies directly, or inversely as any other B , an analogy is always implied, and A and B are not necessarily homogeneous.

51. Cor. 1. If A be directly as B , and B directly as C , A will be directly as C ; for (48) $A : a :: B : b$ and $B : b :: C : c$; therefore (39) $A : a :: C : c$, and A is directly as C (48).

52. Cor. 2. If A be as B and B as C ; A will be as $B \pm C$, or as \sqrt{BC} ; for $A : a :: B : b :: C : c$, and (43) $B : b :: B \pm C : b \pm c$; therefore (39) $A : a :: B \pm C : b \pm c$, and A is as $B \pm C$ (48): and because the ratios of $B : b$ and $C : c$ are equal, $B^2 : b^2 :: BC : bc$ and $B : b :: \sqrt{BC} : \sqrt{bc}$. B therefore, and consequently A (51), is as \sqrt{BC} .

53. Cor. 3. A varies as $m \times A$, or as $\frac{A}{m}$; for $A : a :: m \times A : m \times a$, or as $\frac{A}{m} : \frac{a}{m}$ (41); therefore (48) A is as $m \times A$, or as $\frac{A}{m}$.

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54. Cor. 4. If A be as B , A^2 will be as B^2 , and A^m as B^m , and $\frac{A^{\frac{1}{m}}}{A^m}$ as $\frac{B^{\frac{1}{m}}}{B^m}$; for since $A:a::B:b$ (48) $A^m:a^m::B^m:b^m$, or $\frac{A^{\frac{1}{m}}}{A^m}:\frac{B^{\frac{1}{m}}}{B^m}::\frac{B^{\frac{1}{m}}}{B^m}:\frac{b^{\frac{1}{m}}}{b^m}$ (42); therefore A^m is as B^m and $\frac{A^{\frac{1}{m}}}{A^m}$ as $\frac{B^{\frac{1}{m}}}{B^m}$ (48).

55. Cor. 5. If S be as $V \times T$, $\frac{S}{T}$ will be as V , $\frac{S}{V}$ as T , and $\frac{S}{V \times T}$ a given quantity: for (48) $S:s::VT:vt$ and $\frac{S}{V}:\frac{s}{v}::T:t$ and $\frac{S}{T}:\frac{s}{t}::V:v$, and $\frac{V}{VT}:\frac{s}{vt}::1:1$ (41); therefore $\frac{S}{V}$ is as T , $\frac{S}{T}$ as V , and $\frac{S}{VT}$ does not vary (48).

56. Cor. 6. If A be as B , and C as D , $A \times C$ will be as $B \times D$; for $A:a::B:b$ and $C:c::D:d$; therefore the ratios of $A:a$ and $C:c$ are equal to the ratios of $B:b$ and $D:d$, and the sums of these ratios are equal, or the ratio of $AC:ac$ is equal to that of $BD:bd$ (18), and $AC:ac::BD:bd$ (34), and AC is as BD (48).

57. Cor. 7. If A be inverfely as B , then B will be inverfely as A ; for $A:a::B:b$ (49) and $b:B::a:A$; and if B be to b as 4:5, $A:a::5:4$.

58. Cor. 8. If A be directly as B , and B inverfely as C , A will be inverfely as C ; for $A:a::B:b$ (48) and $B:b::c:C$; therefore $A:a::c:C$ (39), and A is inverfely as C (49).

59. Cor. 9. If A be inverfely as B , and B inverfely as C ; A will be directly as C : for $A:a::b:B$ and $b:B::C:c$, therefore $A:a::C:c$ and A is as C .

60. Cor.

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60. Cor. 10. If $A \times B$ be always the same, A is inverfely as B ; for $A \times B : a \times b :: 1 : 1$ and $A : a :: b : B$; or otherwise, since $A \times B = a \times b$, $A : a :: b : B$ (38).

61. Cor. 11. If A be inverfely as B , it will be as $\frac{1}{B}$ and $v.v$; for $A : a :: b : B$ (49) $:: \frac{1}{B} : \frac{1}{b}$ (41), therefore A is as $\frac{1}{B}$ (48). If A be as $\frac{1}{B}$, $A : a :: \frac{1}{B} : \frac{1}{b}$ (48) $:: b : B$ (41); and A is inverfely as B (49).

62. Cor. 12. If A be as $\frac{L}{M}$, A will be inverfely as $\frac{M}{L}$; for $A : a :: \frac{L}{M} : \frac{l}{m}$ (48) $:: \frac{m}{l} : \frac{M}{L}$ (41); and A is inverfely as $\frac{M}{L}$ (49).

63. DEF. If any quantity A be dependent upon feveral others P, Q, R, S, T , all independent of each other, so that none of the quantities P, Q, R can vary, but A varies in the same ratio, nor either S or T can vary, but A varies in a contrary ratio; A is said to be as P and Q and R directly, and S and T inverfely.

Thus, the fraction $\frac{LMN}{XZ}$ varies as L and M and N directly and X and Z inverfely; because none of the factors in the numerator can vary, but the value of the fraction varies in the same ratio; and none of the factors in the denominator can vary, but the value of the fraction varies in a contrary ratio. If any one of the factors in the numerator be multiplied by 2, 3, &c. the value of the fraction is 2, 3, &c. times its former value; and if any of the factors in the denominator be multiplied by 2, 3, &c. the value of the fraction becomes $\frac{1}{2}, \frac{1}{3},$ &c. of its former value.

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64. Cor. 1. If any magnitude A vary as P and Q and R directly, and as X and Z inversely, these magnitudes being independent of each other, and any one may vary without affecting the rest, A will vary as $\frac{P \times Q \times R}{X \times Z}$. For let any new cotemporary values of P, Q, R, X, Z , be respectively equal to p, q, r, x, z , and let A , when changed in the ratios of $P:p, Q:q, R:r, \frac{1}{X}:\frac{1}{x}$, and $\frac{1}{Z}:\frac{1}{z}$, become a ; that is,

Let $P:p::A:B=A$'s value after the change of P .

$Q:q::B:C=A$'s value after the change of P, Q .

$R:r::C:D=A$'s value after the change of P, Q, R .

$\frac{1}{X}:\frac{1}{x}::D:E=A$'s value after the change of P, Q, R, X .

$\frac{1}{Z}:\frac{1}{z}::E:a=A$'s value after the change of P, Q, R, X, Z .

The sum of the ratios of $A:B, B:C, C:D, D:E, E:a$, or the ratio of $A:a$ (16) is equal to the sum of the ratios of $P:p, Q:q, R:r, \frac{1}{X}:\frac{1}{x}, \frac{1}{Z}:\frac{1}{z}$, or to the ratios of $\frac{PQR}{XZ}:\frac{pqr}{xz}$ (18); and A varies as $\frac{PQR}{XZ}$ (48).

65. Cor. 2. If any of the quantities P, Q, R, X, Z be given or remain invariable, they are to be rejected, and A will vary as the rest. Let P be given or $P=p$, and, because $P:p::A:B, A=B$ (38), and the ratios of $B:C, C:D, D:E, E:a$, or the ratio of $B=A:a$, is equal to the ratio of $\frac{QR}{XZ}:\frac{qr}{xz}$ (18), and A is as $\frac{QR}{XZ}$ (48). Let Q be also given or $Q=q$, and $A(=B=C):a::\frac{R}{XZ}:\frac{r}{xz}$, or A is as $\frac{R}{XZ}$ (48). Hence may be understood what is meant by philosophical writers, when they say the velocity (V) varies

varies as the space (S) directly, and the time (T) inversely, or as $\frac{S}{T}$; and, if T be given, V is as S , and, if S be given, V varies as $\frac{1}{T}$.

66. EXAMP. I. The number of feet (S) described in any time (T), with an uniform velocity (V), encreases or decreases directly with V and T , both of which are independent of each other, and therefore varies as $V \times T$. If V be doubled or encreased in the ratio of 2 : 1, and T be encreased in the ratio of 3 : 1, S will be encreased in the ratios of 2 : 1 and 3 : 1, or of 6 : 1.

67. EXAMP. II. The quantity of matter (\mathcal{Q}) in different bodies varies as the magnitude (M) multiplied into the density (D). For if the magnitude vary in the ratio of $M:m$, and the density or closeness of the constituent parts, in the ratio of $D:d$, the quantity of matter will be changed according to both these ratios, or $\mathcal{Q}:q::M:m$ and $D:d::M \times D:m \times d$, or \mathcal{Q} varies as $M \times D$ (48).

68. EXAMP. III. The velocity (V) of a body moving uniformly in the peripheries of different circles, varies as the radius (R) directly, and periodic time (P) inversely. For, if P remain constant whilst the space described or periphery varies, the velocity will encrease or decrease directly as the space, or as R , because the peripheries of circles are as their diameters; and if the periodic time be encreased in the ratio of 2 or 3 : 1, or universally in the ratio of $P:p$, the velocity will be changed in the ratio of 1 : 2 or 3, or of $p:P$. If therefore the radius be changed in the ratio of $R:r$, and periodic time in the ratio of $P:p$, and V become v , $V:v$ as $R:r$, and $\frac{1}{P}:\frac{1}{p}$, or as $\frac{R}{P}:\frac{r}{p}$, and the velocity is as $\frac{R}{P}$ (48).

INTRODUCTION, &c.

69. EXAMP. IV. If V be the velocity communicated, by the action of a force F , to a quantity of matter represented by \mathcal{Q} , and these quantities be supposed to vary, V will be as $\frac{F}{\mathcal{Q}}$; for if the force become $2F$, $3F$, &c. the velocity will become $2V$, $3V$, &c. and if a quantity of matter become $2\mathcal{Q}$, $3\mathcal{Q}$, &c. the velocity communicated will be $\frac{V}{2}$, $\frac{V}{3}$, &c. and V therefore is as $\frac{F}{\mathcal{Q}}$ (64).

A SYSTEM

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S Y S T E M
O F
N A T U R A L P H I L O S O P H Y.

M E C H A N I C S.

M ECHANICAL philosophy acquired the name from its utility in the construction of machines; but it is now, in a more general sense, understood to comprehend two branches of science, cultivated at different periods of time, denominated Statics, or the science of the equilibrium and relation of powers, and Dynamics, or the science of actual motion. The first professes to describe the construction, properties, and mechanical advantages of machines, and their various combinations, calculated to sustain the pressure of heavy bodies and facilitate their motion; and to investigate the equilibrium of powers acting upon them, or their relative magnitudes, when by opposite exertions they destroy each other's effect or remain quiescent. This branch comprehends that part called the mechanical powers, and is sometimes called practical mechanics. The object of dynamics is the nature, genesis, and change of actual motion, or an investigation of the direction,

rection, quantity, and law of variation, of a force or power capable of generating any motion or change of motion; and vice versâ. This comprehends the laws by which all motions are regulated, the motions resulting from collisions, the theory of oscillations, projectiles, and centripetal forces, and is sometimes called rational mechanics. The principles of statics were calculated and established by Archimedes, and have, since that period, been almost exhausted by the labours of succeeding writers. Galileo demonstrated the laws of descent of heavy bodies, and from him originated the science of dynamics, which has since been prosecuted to an amazing extent in Euler's *Mechanica*, and Newton's *Principia*. As the objects of mechanical philosophy are the equilibrium and motions of bodies; a description of the different qualities, or mechanical affections of matter, producing pressure, motion, and other phenomena, ought to be premised. Mechanics therefore might with great propriety be divided into two parts; of which the first would contain the properties of matter, or the existence, intensity at a given distance, and laws of variation at different distances, of those general principles which are the apparent origin of motion, and the continuation of motion in the material world, investigated by analysis: and the second would be an application of these principles, containing synthetic demonstrations of their effects upon machines, commonly called the mechanical powers, when in equilibrio, and their relation to actual motions, rectilinear and curvilinear, uniform and variable. This is the most obvious division of the subject, and every deviation from it must be attributed to the design of this selection, which is solely the utility of the academic student.

MATTER, AND ITS PROPERTIES.

CHAP. I.

OF MATTER.

70. **M**ATTER is the substance employed in the formation of that part of the creation, whose existence is evidenced by the testimony of the senses; and the most characteristic and prominent marks of it seem to be extension and solidity. Indeed solidity is the most discriminating mark of matter, as it distinguishes it from every thing else; but tastes, smells, sounds equally indicate the existence, though whatever is solid and extended is the common and most general description, of matter. Many other qualities invariably adhere to all matter with which we are acquainted, whether hard, soft, or fluid; and from our ignorance of its internal constitution, or that latent principle, by whose influence its qualities are connected, and from which they necessarily derive their origin, these, and many others inaccessible to the senses, may be ingredients in its essence. All our knowledge of matter, as far as it is related to the present subject, is either geometrical, or philosophical: the first considers matter as being of some magnitude, or circumscribing space and having some figure, then called body, and is usually denominated stereometry, or the mensuration of magnitudes of three dimensions, length, breadth, and thickness; and the second comprehends all the properties of matter addressed to the senses, which may be stiled physical or philosophical, because all the phenomena of nature are conceived to result immediately

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diately from them; as extension, solidity, inertia, and those apparently more active qualities, gravity, magnetism, electricity, cohesion, elasticity. These last are called mechanical affections of matter, or mechanical causes, because they are conceived to reside in matter, and all mechanical phenomena, or changes of motion observable in the material world, appear to be directed by their necessary and uniform agency. The philosophical properties of matter are distinguishable into two kinds, general, and specific: the first are such as universally adhere to every species of matter, and of which no art hath been able to divest them, as *extension, solidity, mobility, quiescibility, inertia, figure, attractions and repulsions*, and probably many more which are too subtle for the observation of sense, or have yet eluded the inquisition of the philosopher; and the second are such as are attached only to particular species of matter, as opacity, transparency, hardness, fluidity, colour, magnetism, elasticity, &c. Another discrimination of these qualities is, that some are incapable of intension or remission, as extension, solidity, inertia, mobility, figure, which are inseparable from a body and its component parts; and others are relative and capable of encrease and decrease, as attractions and repulsions, whose intensity depends upon the magnitude and distance of the attracting or repelling bodies, and as that magnitude and distance may vary without limit, the influence of these powers may, by remoteness, decrease without limit, and become evanescent. A knowledge of the existence, and discrimination of these qualities, is entirely derived from experiment and observation, and constitutes the sum of all that is known concerning matter. The internal constitution of it is unknown, and its effects cannot be investigated without the assistance of experience; and whether any more, besides those eight general properties enumerated, belong to it, and what, or whether any, connection subsist between those already discovered, and to what particular constitution specific qualities are owing, must be derived from the same source, experiment, and if ever discovered, will probably result from the labour of the chymist.

CHAP. II.

OF EXTENSION.

71.* **T**HERE are three kinds of extension, lineal, as the line AB ; superficial, having length and breadth, as $ABCD$; and solid, having length, breadth and thickness, as AH . Ideas of these three kinds of extension are undefinable, and only to be acquired by the senses of seeing and feeling. The perceptions introduced to the mind by seeing or feeling two distant parts of the line AB , a surface $ABCD$, and material body AH , convey respectively the meaning of lineal, superficial and solid extension; which are dissimilar magnitudes, and incapable of comparison with each other. A line, having no breadth, cannot, however repeated, constitute a surface; and a surface, having no thickness, cannot constitute a solid. The parts of a line, though divided without limit, are still lineal extension, and the parts of a surface, or solid, though divided into a number of parts unassignably great, are still superficial and solid extension. This quality is so far essential to matter that it cannot be divested of it, or conceived to exist without it.

FIG.
VIII.

72. DEF. *A magnitude is said to be finite when an equal to it can be assigned, or when its encrease and decrease are limited within assignable bounds.*

73. DEF. *A magnitude is said to be infinitely great or small, when no finite ratio obtains between it and a finite magnitude, or when its encrease or decrease is not limited by any assignable boundary.*

74. PROP.

* Keil's Physics. Muschenbroek.

FIG. VI. 74. PROP. *A finite right line EX is divisible in infinitum.*

DEM. A mathematical point A may be taken between X and E , and another B between A and E , and this process may be continued without limit, because the intermediate points $A, B, \&c.$ having no length, can never coincide with X or E , or with each other. Q. E. D.

FIG. VII. Otherwise: Through E and X draw EP, XC parallel to each other, and, because they never meet, a number of points $a, b, c, d, \&c.$ greater than any assignable number may be taken in XC ; and if right lines be drawn from them to any point P on the other side of EX , they will divide it into a number of parts equal to the number of points $a, b, c, \&c.$: for if they did not, two lines must pass through the same part, that is, either intersect each other in EX , or coincide till they arrive at it, and then diverge, both of which are impossible. There is therefore no assignable limit to the divisibility of EX . Q. E. D.

FIG. VIII. 75. Cor. 1. The finite surface $ABCD$ is infinitely divisible; for it is equally divided with AB , by drawing mathematical right lines $pf, qg, rh, \&c.$ from the points $p, q, r, \&c.$ parallel to AD , which having no breadth cannot coincide. If $ABCD$ be a boundary of the solid AH , and mathematical planes be drawn through $pf, qg, rh, \&c.$ parallel to AF , they cannot coincide, having no thickness, and consequently will divide the solid into the same number of parts with the surface $ABCD$, or line AB .

FIG. IX. 76. Cor. 2. Every finite line, surface and solid, is therefore composed of an unlimited number of parts; and each of these parts is composed of an unlimited number, $\&c.$ For let LM be infinitely greater than Ln , and take $LM : Ln :: Ln : Lo :: Lo : Lp, \&c.$; and drawing any finite line MA , making an angle with LM ,
and

and nB , oC , pD , &c. parallel to MA and terminated by LA ; MA is infinitely greater than nB ; nB than oC ; oC than pD , &c. But MA is infinitely divisible, and right lines, drawn from L to the points of division, will cut nB , oC , &c. into the same number of parts with it.

77. Cor. 3. Matter is therefore infinitely divisible, because extended. And this proposition and corollaries are applicable to magnitudes of every kind, as velocities, forces, times, &c.

S C H O L I U M.

78. The terms infinitely great and small are relative, and imply a comparison with an assignable magnitude, and compared with it, all infinitely great or small magnitudes are to each other in a ratio of equality; but, compared with each other, they admit of the same inequality of ratios with finite magnitudes. The line AM is infinitely divisible, and if the points of division be equidistant, any one part xy , multiplied into their number, is equal to AM ; and $xy : AM ::$ unity : a number unassignably great, or xy is an infinitesimal of the first order. The infinitesimal xy is infinitely greater than wt , and wt than rs ; for $xy : wt :: wt : xs :: LM : Ln$. But infinitesimals of the same order may be to each other in any assignable ratio; for, let LM be to LF as 3, 4, 5, or or $m : 1$, and FG , being drawn parallel to MA , will be equally divided with it, and the infinitesimal $vz : xy :: LM : LF :: 1 : 3, 4, 5$; or any number m . This is called the mathematical divisibility of matter, and is equally applicable to space and all other magnitudes, which may be represented by extension; but that an actual divisibility, or actual separation of the parts of matter, is limited by certain inviolate bounds, is inferred from the identity of natural substances.

PHENOM. I. Salt dissolved in a menstruum, becomes the same salt when the menstruum is supposed to dry gradually. Salt converted

verted, by a chymical process, to an acid spirit, is, by reversing the process, regenerated into the same salt. Metals liquidated by fire, are restored to their pristine state by the attraction of cohesion. But were the decomposed particles of the salt, or metal, attenuated by the action of the menstruum or fire, and again to coalesce by their cohesive force, substances of a different texture and appearance, and of different specific gravities, would result. The most powerful natural agent, with which we are acquainted, is fire; which, whether artificial or collected in the focus of a burning glass, only attenuates bodies to a limited degree by converting them into a thick smoke, glass, &c. Every species of matter fluctuates and decays by the separation of its component parts, and is renovated by their accession; but were the nutritious particles, administering to the increase of natural substances, capable of divisibility or diminution by attrition, &c. new species of substances would result with new properties and characters. Water and earth, composed of old particles, and fragments of particles produced by attrition, would have a different specific gravity, from water and earth composed of entire and unbroken particles; and the nature and textures of these and all other substances would by repeated attritions, be perpetually changing, which is contrary to experience. Bodies therefore break, not in the midst of solid particles, but where those particles cohere in a few points; and the divisibility of matter is only a separation of its constituent parts, effected by a dissolution of that cohesive force which unites them, and is limited by ultimate elementary particles which are a perfect repletion of space, without pores and indivisible.

C H A P. III.

O F S O L I D I T Y.

79. **T**HE second philosophical quality of matter is that by which it occupies space to the exclusion of all other matter, and is called its solidity; but the meaning of it cannot be conveyed by words, being undefinable and only to be acquired by the sense of touching. The perception introduced to the mind by the insuperable resistance felt in a body is an idea of solidity. This resistance, and total exclusion of matter from space already occupied, is inseparable from all matter with which we are acquainted, whether hard, soft, or fluid; and the little resistance experienced in some bodies, is not owing to a want of solidity, but to their fluidity and softness, by which qualities their component parts are easily displaced. That hard and soft bodies are solid, and cannot occupy the same part of space, appears from uncontroverted experience; for the opposite sides of the substance compressing them are never found to meet, but by removing the intervening parts. And the same is discovered to obtain in all fluids within the reach of observation; for no portion of water, air, mercury, &c. can occupy the same place, because they afford an insuperable resistance preventing the coincidence of the opposite sides of the substance compressing them; and their dimensions are never found to be diminished by pressure without an adequate cause, viz. compressibility, and transmission through the vacuities of the vessel containing them. Numberless facts demonstrate electric and magnetic effluvia, whose parts are immeasurably minute, to be capable of impulse and resistance like other fluids; and it is inferred from analogy, that these fluids, circumscribed and compressed by plane surfaces, would invariably oppose their junction, were the vacuities in those surfaces less than the particles of the

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the fluids. Solidity therefore is an universal attribute of matter, but its cause is unknown and probably undiscoverable by human faculties. The approach of one body to another is apparently limited by the actual contact of their nearest parts, which absolutely fill the spaces occupied by them, and consequently a nearer access and greater surfaces of contact, only result from their dislocation; but this is not allowed without controversion. Compressibility, contraction by cold, elasticity, &c. prove that the minute parts of some bodies are not in mathematical contact, and their resistance is an indication of repulsion; and the nearer access therefore of two bodies is imagined to be impeded by a strong repulsive power, which being overcome, a mutual penetration of parts ensues, and the bodies occupy the same part of space. But this doctrine is still only hypothetical; and though the minute parts of some bodies do exert an influence at a distance, and a repulsive power be confessed to obtain between them, it cannot be concluded, generally, to be the only cause preventing their nearer approach, nor admitted as a general principle in nature, contrary to the common apprehension of mankind, till established by satisfactory and uncontroverted experiments. The transmission of one body through another, and apparent penetration of parts, seem, and may be conceived, to result from empty unoccupied space between those parts; and the quantity of these vacuities is collected from the following phenomena.

80. PHENOM. Many vacuities or pores are actually visible, through a microscope, in every species of animals, vegetables and fossils. The bottom of the sea is visible at a greater depth than sixty feet. A man's finger, placed before the aperture of a dark chamber, is transparent by the passage of light through its pores: and light is transmitted in almost every direction through glass and water, and when condensed three thousand times in the focus of a burning glass, it seems to be admitted into water and glass without obstruction. Electric effluvia pass through gold with a velocity unassignably great, and the magnetic power is transmitted through every species of matter, except iron, without diminution.

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The volatile spirit of sulphur tinges, with a brown colour, silver surrounded with repeated coverings of cloth or paper; the scents of musk, civet, &c. pass through wood. Air and water imbibe each other; oils penetrate the vacuities of sulphur and some stones; mercury penetrates the pores of gold, brass, and is transmitted through human skin, leather, &c.; water is transmitted through the membranes of animals, and the fine tubes of vegetables, and may, by compression, be forced through gold, silver, &c.* From numberless chymical experiments it appears, that all animals, vegetables, and fossils, yield water plentifully by the force of heat; and all bodies, whether hard or fluid, admit the particles of fire into their pores, through which it passes and is dissipated.

81. DEF. *The magnitude of a body is the magnitude of solid extension, that is, length, breadth and thickness, or the number of cubical inches contained in it.*

82. DEF. *The quantity of matter in a body is the number of equal particles contained in it. If the matter composing different bodies be reduced to equal particles without pores, the quantity of matter in each will be equal to one particle multiplied respectively into their number.*

83. DEF. *Density of a body is the contiguity, or close adhesion, of its particles; but this term usually implies the ratio obtaining between the number of equal particles, or quantities of matter, in bodies of the same magnitude.*

84. DEF. *Homogeneous bodies are those which have the same density in every part; heterogeneous are those which have not.*

85. DEF.

* In the celebrated Florentine experiment, a quantity of water was enclosed in a hollow sphere of silver, and then forcibly compressed by screws, till the fluid was seen to ooze through the pores of the metal and cover its surface like a dew.

85. DEF. *The porosity of a body is the remoteness of its elementary particles, or the ratio of the quantities of vacuity in different bodies of equal bulk.*

86. PROP. *If Q, D, B represent respectively the quantity of matter, density, and magnitude, of a body, and be supposed to vary; Q will vary as $D \times B$.*

DEM. If D , or B , be encreased or diminished in the ratio of 2, 3, &c. to 1, Q will evidently be encreased or diminished in the same ratio, and D and B are unconnected; therefore (64) Q is as $D \times B$. Q. E. D.

87. Cor. 1. Since Q is as $D \times B$, if q, d, b be cotemporary new values of Q, D, B ; $Q : q :: DB : db$, and $\frac{Q}{D} : \frac{q}{d} :: B : b$, & $\frac{Q}{B} : \frac{q}{b} :: D : d$ (41); or if D be given, Q is as B ; and if B be given, Q is as D ; and if Q be given, D is $\frac{1}{B}$ and B as $\frac{1}{D}$. This cor. follows also from (65).

88. Cor. 2. If numbers, or lines, whose ratio is the same with that of D and B , be substituted for them, Q is properly expressed by the product of these numbers, or a rectangle whose sides are these lines.

89. Cor. 3. If the density of a body be encreased in the ratio of 2, 3, &c. to 1, the porosity (P) will be diminished in the ratio of 1 to 2, 3, &c.; and consequently P is as $\frac{1}{D}$, or (87) as $\frac{B}{Q}$. If B be given, P is as $\frac{1}{Q}$; and if Q be given, P is as B ; and P being given, B is as Q .

EXAMP.

S O L I D I T Y.

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EXAMP. I. The density of gold is to that of water as $19\frac{1}{2} : 1$, and consequently the relative quantity of pores in gold and water is as $1 : 19\frac{1}{2}$. The relative densities of gold and cork are as $81\frac{1}{2} : 1$, and their quantities of pores therefore as $1 : 81\frac{1}{2}$. If one half of the magnitude of gold were vacuous, the relative quantities of pore and solid parts in the water and cork, would be respectively as $39 : 1$ and $163 : 1$.

EXAMP. II. If a body whose magnitude is A were constructed of particles cohering in such a manner as to have $\frac{1}{2}$ of its magnitude vacuous, and these particles were similarly constructed of other less particles having $\frac{1}{2}$ of their magnitude vacuous, and the third order of particles were elementary, and a perfect repletion of space, the quantity of vacuity would be equal to $A \times \frac{1}{2} + A \times \frac{1}{2} \times \frac{1}{2} + A \times \frac{1}{2} \times \frac{1}{2} \times \frac{1}{2} = A \times \frac{1}{2} + \frac{1}{4} + \frac{1}{8} = A \times \frac{7}{8}$. And the magnitude of vacuity is to the magnitude of matter as $7 : 1$.

S C H O L I U M.

90. What the figure and magnitude of the elementary particles of matter are, cannot be known from the senses, which, with every microscopical assistance, are unable to discern them. An elementary particle of matter hath probably never yet been seen. A number of elementary particles, uniting by the power of cohesion, form greater particles; and these, uniting again by the same power, form greater still; and, this process being made repeatedly, a corpuscle is at length formed of a sensible bulk. All bodies seem to be composed of these derivative corpuscles, which, formed of more or fewer repeated unions, compose bodies more or less dense. These derivative corpuscles seem sometimes to be similar: if a beam of light be separated by a prism into small coloured rays, and any slender ray of the same colour be minutely examined, its component parts seem to be similar, because they affect the sight exactly in the same manner. Pure mercury squeezed through the pores of leather, or raised into fume, and received upon clean glass,

exhibit globules similar and undistinguishable. This is also observable of the vapours of pure water. And, in other species of matter, the derivative particles are combined from others exactly similar to them; a red globe of blood is observed, through a microscope, to be composed of six yellowish ferous globes, and every one of these is composed of six lymphatic globes; but farther the microscope does not enable us to proceed. Every species of matter, however different in density, may be conceived to be formed of equal and similar elementary particles. Thus, the particles

FIG. X. *A, B, C*, composed each of six equal elementary particles, are different; and these particles may, by a different combination, form different still; and these repeated coalitions may, by changing the original and succeeding numbers and their positions, form masses of matter differing in endless variety.

U N I F O R M M O T I O N.

91. From the solidity of matter, whether owing to repulsion or actual contact, results its capacity of impulse or of being protruded, and consequent mobility of body, as far as the cause of mobility seems to be known. If the body *A*, in motion, strikes another body *B*, not retained in its place by any force, *B* will be protruded and move; because, from their solidity, they cannot penetrate each other's dimensions; and the communication of motion in this case is the necessary consequence of solidity. Motion supposes the successive existence of the body moving in different parts of space, and therefore cannot be understood without presupposing a knowledge of space and time; which therefore must be premised.

92. DEF. *Space is that which contains the whole sensible creation, is unlimited, and its parts are homogeneous, inseparable, co-existing and unresisting.*

93. Cor. Since space is unlimited in every direction, and its parts similar and undistinguishable, any portion of it can only be ascer-

ascertained by its relation to some assumed sensible mark or object. Hence of place, which is a part of space, there are two kinds, absolute and relative.

94. DEF. *The absolute place of a body is that portion of this unlimited space, which is occupied by it when fixed and immoveable; and the relative place of a body is its situation with respect to some assumed mark or object, which itself may be moveable.*

95. Cor. The absolute and relative place of a body coincide, when it and the assumed object are immoveable, or remain in the same part of fixed space. But we cannot perceive the absolute place of any object.

96. DEF. *Duration is that which flows uniformly, is unbounded, continuous, whose parts are similar, and no two exist together.* An idea of duration is obtained by observing the interval between our ideas, or between the successive appearances of any external object; and, as it is, strictly, only measurable by a portion of duration, and no two parts exist together, consequently cannot be compared by juxta position, time, which is a part of duration, is of two kinds, absolute and relative.

97. DEF. *Absolute time is a portion of duration, whose quantity is only known by a comparison with another portion; and consequently the relation between any two parts of absolute time is undiscoverable. Relative time is a part of duration which elapses during any motion of body, or any succession of external appearances.*

98. DEF. *An instant is the boundary between any two contiguous portions of time, as a point is the boundary of any contiguous lines; and a moment is any small portion of time.*

99. DEF-

99. DEF. *Absolute motion is a change of absolute place, and absolute rest is a permanence in the same absolute place.*

100. DEF. *The direction of motion is the position of the line upon which it is made. When the motion is rectilineal, its direction is this right line; when curvilineal, the direction is the tangent to that point of the curve where the moving point is.*

101. DEF. *Velocity is the quickness or slowness of motion, or the rate at which a body moves.*

102. DEF. *A body is said to move with an uniform, accelerated, or retarded velocity, when it continues the same, encreases, or decreases. When the encrease, or decrease, of velocity is the same in any equal times, the acceleration, or retardation, is said to be uniform; and when this encrease or decrease of velocity, encreases or decreases in any equal times, the acceleration, or retardation, encreases or decreases in the same ratio.*

103. DEF. *The line joining all the successive places, through which a moving body passes, is called the space described.*

104. PROP. *If S represent the space uniformly described with the velocity V in the time T, and these magnitudes be supposed to vary, S will vary as $V \times T$.**

DEM.

* This proposition is proved, not inelegantly, by the following process: Let S and s be respectively described uniformly, with the velocities V and v, in the times T and t; and $T:t::S:\frac{S \times t}{T}$ (38) = the space described in the time t with an uniform velocity equal to V; and $\frac{S \times t}{T}:s::V:v$ and $S \times t \times v = s \times V \times T$ (37), and (38) $S:s::V \times T:v \times t$. Q. E. D.

DEM. If either V or T be encreased or diminished in the ratio of 2, 3, 4 $n : 1$, the space will evidently be encreased or diminished in the same ratio; and because V and T are unconnected, or either of them may be changed without affecting the other, S will vary as $V \times T$ (64). Q. E. D.

105. Cor. 1. Because S is as $V \times T$, V will be as $\frac{S}{T}$, T as $\frac{S}{V}$, and, if S be given, T as $\frac{1}{V}$ and V as $\frac{1}{T}$ (65).

106. Cor. 2. If V be given, S is as T , or the spaces described, with the same uniform velocity, are true measures of time, and may be substituted for it. And, T being given, S is as V , or the spaces described in the same time uniformly, are measures of the velocity: if V , or any other symbol, be called velocity, it denotes the length or number of feet described uniformly in one second, or any other time. Let any values of the velocity be to each other as $A : B$, or as 3 : 2, and any cotemporary values of the time be as $C : D$, or as 5 : 4, and the spaces described will be to each other as $X : Z$, or as $3 \times 5 : 2 \times 4$. PLATE
I.
FIG. III.

107. Cor. 3. If lines therefore be substituted for V and T , the space S will be as the rectangle, whose sides are these lines; and, if numbers be substituted for them, the space will be equal to the product of these numbers, the time equal to the quotient of the space divided by the velocity, and the velocity equal to the quotient of the space by the time. Let V be equal to n feet in one second, and $T = p$ seconds; and S is actually equal to $n \times p$ feet, T actually equal to $\frac{S}{V}$ seconds, and $V = \frac{S}{T}$ feet in one second.

VARIABLE MOTION.

108. PROP. *In variable finite velocities, the velocity during an infinitely small time, is uniform.*

DEM. The encrease, or decrease of velocity produced in any finite time, is finite, and, if the whole encrease or decrease be divided into any infinite number of increments or decrements, each will be infinitely small and vanish, compared with the whole velocity, which therefore is uniform. Q. E. D.

109. Cor. If therefore S' , V' , T' , represent corresponding increments of S , V , T , respectively; V' will be evanescent compared with V ; S' will vary as $V \times T'$ (104), V as $\frac{S'}{T'}$, and T' as $\frac{S'}{V}$ (105).

PLATE I.
FIG. XI. 110. PROP. *If the abscissa AS of any curve DEF represent the whole time, and the ordinates AD, BE, CF, &c. be as the velocities at the instants A, B, C, &c.. the spaces described in the times AB, BC, will vary as the areas AE, BF, &c.*

DEM. For let Am , mn , np , &c. be moments of time, and the spaces described in these moments, are as the rectangular parallelograms Aq , mr , nv , (107) &c.; and the whole spaces, described in the sums of these moments, i.e. in the times AB , BC , are as the sums of these parallelograms, or as the areas AE , BF . Q. E. D.

111. Cor. 1. If the relation between V and T , or the ordinate and abscissa, be given, the relation between the spaces, or the areas AE , BF may be found.

112. Cor.

112. Cor. 2. If the area AE be always as the space described (S) in the time AB or T , the velocity is as the ordinate BE ; for according to this supposition, S' is as $AD \times Am$, or $AD \times T'$, and (109) as $V \times T'$; therefore V is as AD . PLATE I. FIG. XI.

113. Cor. 3. If AS represent the space described, and the ordinates $AD, BE, CF, \&c.$ be always inversely as the velocities at those points; the times of describing $AB, BC, \&c.$ will be as the corresponding areas $AE, BF, \&c.$ For the time of describing any small space Am is as $\frac{S'}{v}$, or as $Am \times AD$; and consequently the sum of the moments, or whole time of describing AB , is as AE .

114. PROP. *The acceleration and retardation of velocity, vary as the change uniformly produced directly, and the moment of time in which they are produced inversely.*

If the changes of velocity, uniformly produced in the same time, be as $2:1$, the acceleration is as $2:1$ (102); and if the times, in which the same change of velocity is effected, be as $1:3$, the law of acceleration is as $3:1$, or inversely as the times; because if the times were as $3:3$, the changes of velocity uniformly produced in these equal times, which measure the ratio of acceleration, would evidently be as $3:1$, therefore the acceleration is as $\frac{V'}{T'}$, and in the same manner the retardation is as $-\frac{V'}{T'}$. Q. E. D.

115. Cor. If the velocity vary as the time, the acceleration and retardation are constant; for $V:v::T:t$, and, supposing $n =$ the number of moments of time, or changes of velocity, $\frac{V}{n} (V') :$
 $\frac{v}{n} (v') :: \frac{T}{n} (T') : \frac{t}{n} (t')$; or V' is as T' and $\frac{V'}{T'}$ is given.

G

NOTE.

N O T E.

PLATE I. *116. PROP. In variable motions, if the velocity (V) in any point L , be as any power of the space described, the time (T) of describing that space may be found.
FIG. XII.

Let the body begin to move from A , and $AL = z$, and V be as z^n , and T the time of describing AL . \dot{T} is as $\frac{\dot{z}}{V}$ (109) as $\frac{\dot{z}}{z^n}$, and $T = \frac{z^{1-n}}{1-n} + C$ (correction). Q. E. D.

117. Cor. 1. If V be constant, or $n = 0$; T will be as (z) the space described.

118. Cor. 2. If V be as the space described, or $n = 1$; $T = \frac{1}{0}$ or is infinite, and the body will never move from the point A . There are therefore no motions in nature, whose velocities are not in a less ratio than that of the spaces described, in the beginning of motion.

119. Cor. 3. If V be as the power of the space described, whose exponent is $\frac{1}{2}$, $-\frac{1}{2}$, -1 , -2 , &c. the time will be as that power of the space described, whose exponent is $\frac{1}{2}$, $\frac{3}{2}$, 2 , 3 , &c.

FIG. XIII.

120. Cor. 4. Let the body move in AL , and the velocity (V) be as the ordinates of the curve line AM , which meets LA in A , and consequently at A , $V = 0$. It is evident, that if the time of describing AL be finite, or (118) V be as some power of AL , whose exponent is less than unity, or fractional, the tangent at A will be perpendicular to AL . Let V be as AL^n and \dot{V} as $n \times AL^{n-1} \times \dot{AL}$, and $\dot{V} : \dot{AL} :: n \times AL^{n-1} : 1$; and, if \dot{V} be as any power of AL , whose exponent is less than unity, viz. $\frac{1}{2}$, $\dot{V} : \dot{AL} :: \frac{1}{2} \times AL^{-\frac{1}{2}} : 1 :: \frac{1}{2AL^{\frac{1}{2}}} : 1 :: \frac{1}{0} : 1$, when AL vanishes, and consequently the fluxion of the ordinate LM is infinitely greater than that of the absciss AL .

121. EXAMP. If AMB be a semicircle, $AB = 2a$, $AL = s$, and the velocity at any point L , as LM , or such as would describe $n \times LM$ feet in 1"; $T =$ the fluent of $\frac{s}{v}$, or of $\frac{s}{n \times LM}$, or $\frac{s}{n \times \sqrt{2as - s^2}}$, or $\frac{as}{na \times \sqrt{2as - s^2}}$, or $\frac{AM}{n \times AC}$ seconds. Therefore the whole time of describing $AB = \frac{AMB}{n \times AC}$ seconds, which is a given quantity, and whatever be the length of AB , it will be described in the same time.

PLATE II.
FIG. XIV.

122. PROP. The times of describing AL and al (T and t) with velocities which are always to each other as the ordinates LM , lm of similar curves AM , am , are equal.

* Euler's Mechanica.

DEM.

DEM. Let AL and al be homologous spaces, whose ratio is that of $p : q$, and if $AL = s$, and $LM = v$; $al = \frac{AL \times q}{p}$, and $lm = LM \times \frac{q}{p}$. The time of describing AL or $(T) = \text{flu. } \frac{\dot{s}}{v}$, and $t = \text{flu. } \frac{al}{lm} = \text{flu. } \frac{AL}{LM} = \text{flu. } \frac{\dot{s}}{v} = T$. Q. E. D.

123. DEF. *The scale of velocity is a line whose ordinates are as the velocities at those points from whence they are drawn. If the body move in AL , as in the last proposition, Ec. AM is the scale of velocity.*

124. DEF. *The scale of time is a line whose ordinates are as the time. If the ordinate MQ be always as the time of describing AM , the line AQ is called the scale of time.*

PLATE
II.
FIG. XV.

125. PROP. *If the scale of velocity (V) be given, the scale of the time (T) may be found.*

For \dot{T} is as $\frac{\dot{s}}{V}$ and $T = \text{flu. of } \frac{\dot{s}}{V} + C$, and if the relation of the ordinate and absciss, or V and S be known, the fluent of $\frac{\dot{s}}{V}$ may be found. Q. E. D.

126. Cor. If the scale of velocity AN be a circular arc, the time is that arc; for \dot{T} is as $\frac{\dot{s}}{V}$ or $\frac{\dot{s}}{MN}$, or AN , and T is as AN .

127. PROP. *The scale of time AQ being given to construct the scale of velocity.*

\dot{T} is as $\frac{\dot{s}}{V}$, and V as $\frac{\dot{s}}{\dot{T}}$ or $\frac{MQ}{MO}$, supposing QO to be perpendicular to the curve. Take therefore $MN = \frac{MQ}{MO}$, which is $= \frac{\dot{s}}{\dot{T}}$ or as V , and AN is the scale of velocity.

128. Cor. 1. If AQ be a right line, and T as $m \times S$; \dot{T} is as $m \times \dot{s}$, and V is as $\frac{\dot{s}}{\dot{T}}$ as $\frac{\dot{s}}{m \times \dot{s}}$ as $\frac{1}{m}$, and AN is a right line.

129. Cor. 2. If T be as S^m and \dot{T} as $m \times S^{m-1} \dot{s}$; V will be as $\frac{\dot{s}}{m \times S^{m-1} \dot{s}}$, or as $\frac{1}{m \times S^{m-1}}$. If AQ be the common parabola, or $m = \frac{1}{2}$; $\frac{1}{m \times S^{m-1}}$ will be equal to $2s^{\frac{1}{2}}$, and AN is also a parabola; and the relation of $T(MQ)$ and the absciss AM or S being known, the relation of $V(MN)$ to S may always be found.

R E L A T I V E M O T I O N.

130. DEF. *Relative motion is a change of relative place, and relative rest is a permanence in the same relative place.* Relative and apparent motion are sometimes distinguished from each other, the first being defined to be that which is attributed to a moving object by an observer in motion, and the second that attributed to an object really quiescent by an observer in motion.

131. DEF. *The relative velocity of two bodies is the velocity with which they accede to, or recede from, each other.*

132. Cor. 1. Relative motion and rest of a body coincide with absolute, when the assumed object, by which its situation is determined, remains in the same part of fixed space. If the earth be quiescent, every ship which moves, or is quiescent, with respect to a fixed object upon the shore, moves also absolutely, or is absolutely quiescent.

PLATE
II.
FIG.
XVI.

133. Cor. 2. If the object, by which the situation of a body is determined, move, absolute and relative motion and rest do not coincide. If the body B , fixed at the point B in the line AB , be the assumed object with which the situation of A , placed in the same line and partaking of its motion, be compared, its absolute velocity will be the sum or difference of its own velocity and that of the line, according as they move in the same, or opposite, directions. If the line AB move in the direction YX , with an uniform velocity of 100 feet in 1", and A also move in the same direction, with an uniform velocity of 50 feet in 1"; A 's whole absolute velocity in the direction YX is 150 feet in 1", and its relative velocity, or uniform recess from B , is 50 feet in 1". If the line AB move uniformly in the direction YX with a velocity of 100 feet in 1", and A move in the opposite direction XY , with a velocity of 100 feet in 1", it will be absolutely quiescent, and its relative

relative velocity, or B 's access to it, will be 100 feet in 1" uniformly: if A 's velocity, in the last supposition, be 50 feet in 1" uniformly, its relative velocity, or approach to B will be 50 feet in 1", and its absolute velocity will also be 50 feet in 1" in an opposite direction.

134. Cor. 3. A spectator therefore, placed upon the moving line AB at the point B , only perceives A 's relative motion, which may be always different from the true.

*135. PROP. *If two bodies A and B move at the same time from A in the directions AM and AL, with uniform velocities a and b respectively, their relative velocity will be uniform, and to a as the sine of $\angle BAD : \angle BDA$, and to b as $\sin. \angle BAD : \sin. \angle DBA$ or $\angle BAF$, supposing B and D to be cotemporary positions of A and B, and AF to be parallel to BD.*

PLATE
II.
FIG.
XVII.

DEM. Let $B, D; C, E; M, L;$ be cotemporary positions of A and B ; and (106) $a : b :: AB : AD :: AC : AE :: AM : AL$; and consequently BD, CE, LM , which measure the relative velocity, are parallel to each other, and encrease uniformly because AM does. But the relative velocity is to $a :: LM : AM :: \sin. \angle MAL : \sin. \angle ALM$, and it is to $b :: LM : AL :: \sin. \angle LAM : \sin. \angle LMA$ or $\angle MAH$. Q. E. D.

136. Cor. 1. The only sensible motion of B to a spectator placed at A , and supposing himself to be quiescent, will be along Ab parallel and equal to ML ; and when A hath really described AM , B will appear to have described $Ab = ML$ uniformly.

137. Cor. 2. If A and B move uniformly with any equal velocities in parallel directions AM, BL , their relative situation is not changed, being always equal, and parallel, to AB ; and if any space.

FIG.
XVIII.

space $AMBL$, containing any number of bodies, move, their relative situations are not affected by it. To a spectator therefore placed in A , in this moving space, and supposing himself to be quiescent, any body B will appear to be quiescent.

138. Cor. 3. It is evident that the absolute motion of a body may be changed into innumerable relative motions uniform and rectilinear, if its absolute motion and that of the body, by which its situation is determined, be so.

PLATE
II.
FIG.
XIX.

139. PROP. *If two bodies A and B move, at the same time, from the point A, in the directions AM, AL, with uniform velocities, as AM, AL respectively, and B be considered as quiescent; A's apparent motion will be uniform and its direction and velocity as the diagonal of a parallelogram whose sides are AM, and AD equal and opposite to AL.*

DEM. Let m, l be cotemporary positions of A and B , or let $Am : Al :: AM : AL$ and lm, Lm are parallel. But when B is at l and L , A 's distance is equal to lm, LM ; or, supposing B to be quiescent at the point A , and AE, AD be respectively equal and opposite to Al, AL , and En, DN be drawn parallel to AM , A will appear at n and N . But An varies as Am , and therefore increases uniformly, and AN is the diagonal of a parallelogram, whose sides are $AD = AL$, and AM . Q. E. D.

140. Cor. 1. The apparent and real velocities of A are to each other as $AN : AM :: \sin. \angle MAD : \sin. \angle NAD$.

141. Cor. 2. Supposing A and B to move as before, and A to be quiescent, the distances and directions of A from B will be Ab , and AH , supposing them to be respectively equal and parallel to ml, ML . Whilst B therefore really describes AL , it appears to describe $AH = AN$, and in the same right line with it; and the
apparent

apparent motion of A and B , seen from each other, are equal and opposite.

142. Cor. 3. An uniform relative motion along the diagonal of any parallelogram $AMND$ may be considered as resulting from, and equivalent to, two uniform absolute motions in the sides AM and $AL = AD$; and these motions are in the same plane, and their velocities are to each other as the diagonal and sides.

FIG.
XIX.

143. Cor. 4. If the motions of A and B be opposite, their relative motion is equal to the sum of their real motions, AN being, in that supposition, equal to $AM + AD$. If the motions of A and B conspire, their relative motion is equal to the difference of their real motions, or of AM and AD . And if the real velocities of A and B be variable, according to the same law, their relative velocity will vary according to that law.

144. Cor. 5. The real motions in AM and AL are said to be equivalent to a motion in AN , because they produce the same effect in the same time, as if B were quiescent at A , and A were to describe AN uniformly in that time. And in the same manner two motions in AN and AD are equivalent to a motion in AM , or its opposite and equal AP , according as the body describing AD , or AN , is supposed to be quiescent.

145. Cor. 6. If the real motion of A and B , or AM and AL , and their inclination be known, the relative motion of either of them may be found; for AM , MN and $\angle MAD$, and consequently its supplement AMN , being known, the base AN , and the $\angle MAN$, or the inclination of A 's apparent to its real path, may be found.

146. Cor.

146. Cor. 7. The relative and real motion of A , or AM , AN , and the angle MAN , being known, the real motion of B , which is supposed to be quiescent, or $AL = AD$ may be found.

PLATE
II.
FIG. XX.

147. PROP. *If two bodies A and B move at the same time uniformly in the directions AM, BL, with velocities equal to a and b, the relative motion of B will be uniform and rectilinear upon the line BN, supposing AN to be always equal, and parallel to right lines joining the contemporary positions of A and B.*

DEM. Let $m, l; M, L$, be cotemporary positions of A and B , or let $a : b :: Am : Bl :: AM : BL$, and ml, ML will be the distances of B from A in the points m and M . And if A be considered as quiescent at A , taking An, AN , respectively parallel and equal to ml, ML , B will appear at n and N in the line BnN , which is therefore the apparent path of B . But $a : b :: Am (nl) : Bl :: AM (NL) : BL$; therefore BnN is a right line, and increases uniformly, because it varies as BL . Q. E. D.

148. Cor. If A and B move as in this proposition, and B be considered as quiescent, A 's apparent path will be in the right line ADE , parallel and equal to BnN , as is evident by taking BD and BE respectively parallel and equal to lm and LM . A 's velocity will therefore be uniform and equal to B 's velocity, when A was considered as quiescent.

149. The absolute and relative motions of B , and their directions being known, the direction and velocity of A 's motion, which occasioned B 's relative motion, may be found; for, let A and B be cotemporary positions of A and B , and let Bl, Bn , be the spaces described in the same time by B 's absolute and apparent motion, and if ln be joined, AM drawn parallel to it, will be A 's path, and A 's velocity : B 's velocity :: $nl : Bl$.

150. Cor.

S O L I D I T Y.

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FIG.
XX.

150. Cor. 1. An absolute uniform and rectilinear motion BL may be changed into any infinite number of uniform relative rectilinear motions, by changing the direction and velocity of A 's motion; for BN may be drawn of any length, and in any direction, and the direction and velocity of A , moving uniformly so as to produce this relative motion, may be found by drawing AM parallel and equal to NL .

151. Cor. 2. When two bodies move uniformly in right lines, and one of them is considered as quiescent, the relative motion of the other is therefore uniform and rectilinear; and the spaces, described by it relatively, vary as the velocity multiplied into the time (104), or S varies as $V \times T$, T as $\frac{S}{V}$, and V as $\frac{S}{T}$.

152. Cor. 3. If the relative motion of B be variable, the absolute motion of either A or B is variable.

153. PROP. *If a body A move in any curve line AmM, and another body B move in any other curve BIL, B's apparent motion may be found.*

PLATE
II.
FIG.
XXI.

Let m, l, M, L , be cotemporary positions of A and B ; and B 's distance and the right lines in which it appears at m and M , are ml and ML . Or, if A be considered as quiescent, and An, AN , be taken respectively parallel and equal to ml, ML , B will appear at n and N , and its path will appear to be BnN . Q. E. I.

154. Cor. It is evident that the curve BnN , the relative path of B , may be described by the motion of A in different curves; but all the curves described by A will be equal and similar, and have their corresponding parts parallel, because they are always subtended by right lines equal and parallel to nl, NL , &c.

H

155. PROP.

FIG.
XXII.

155. PROP. *If B be absolutely quiescent, and A move in any curve line AmM uniformly, B will appear to move uniformly in a curve equal and similar to AmM, and these curves are similarly situated with respect to the points A and B.*

Let m, M be any contiguous positions of A , which being considered as quiescent at the point A , the cotemporary positions of B are n and N , supposing An, AN to be respectively parallel and equal to Bm, BM . And because $An = Bm$ and $AN = BM$, and the angle $NAn = \text{angle } mBM$, the small arc $Nn = mM$, and they make equal angles with the distances AN, BM , and An, Bm . Q. E. D.

156. Cor. 1. *If B move in the curve BnN, and A be quiescent at the point A, it will appear to describe the curve AmM similar and equal to BnN.*

PLATE
III.
FIG.
XXIII.

157. Cor. 2. *If A and B, placed in the same right line ABC, describe the circles AmM, BlL in the same time uniformly, and A be considered as quiescent, B's apparent path is a circle whose center is A and radius AB. For let m, l, and M, L, be cotemporary positions of A and B, and B will appear at l and L; or, since A is considered as quiescent, if An, AN, be drawn parallel and equal to ml, ML, it will appear at n and N, and its apparent distance from A always = AB, and the direction of B's apparent motion is the same with that of A and B.*

158. Cor. 3. *If B be considered as quiescent, and they move as in the last corollary, A's apparent path will be a circle whose center is B and radius BA, and the direction of its motion is the same with that of B in that corollary. For taking any cotemporary positions m, l, and M, L; and, drawing Bn, BN respectively parallel and equal to ml, mL, A will appear at n and N, and its apparent distance from B always = AB.*

159. Cor.

S O L I D I T Y.

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FIG.
XXIV.

159. Cor. 4. The relative motions of A and B seen from each other, are equal and in similar curves. For let m and M be any contiguous positions of A , and l, L corresponding positions of B ; and joining ml, ML , B 's relative motion, whilst it really describes lL , will make it describe nN , supposing A to be quiescent, and An, AN to be respectively equal and parallel to lm, LM . Supposing B to be quiescent, and Bb, BH to be equal and parallel to ml, ML , or An, AN , A will appear to describe bH , which is equal and similar to Nn .

160. Cor. 5. If B be immoveable and not placed in the plane of A 's motion, it will appear to describe a line equal and similar to the line described by A , and they will be in parallel planes; for taking any two contiguous portions of A 's path, the right lines joining the real and apparent places of A and B are equal and parallel, and consequently the small right lines joining them are equal and parallel: and this and the preceding corollaries are true wherever the eye's imaginary place is.

S C H O L I U M.

161. To distinguish real and absolute, from relative motion, is an important, but very difficult problem; because the same apparent motions may result from real motions combined in endless variety. Every apparent motion of a body results from, and may be explained by, its real motion and that of the spectator and v, v ; and therefore its real motion, and the real motion of the spectator, or this, and the apparent motion of the body, must be presupposed, and, with the assistance of mathematics, the other may then be detected. A spectator ignorant of the earth's annual and diurnal motion, and supposing himself to be quiescent, must draw erroneous conclusions, in all his reasonings, concerning the absolute motion of any body, because these motions of the earth will communicate an apparent motion to a body absolutely quiescent, and affect the absolute motion of a body moving (134). The relative motions of bodies must therefore be ascertained from observation, and their real motions, when

discoverable, are deduced from these, the properties of real motion, and the quantity of mechanical powers, or qualities in matter, that are found to generate motion. If the properties, invariably attached to all real motions, obtain in any observed motion; or if any mechanical cause be proved to exert an influence, which is exactly competent for the production of any observed motion, it may safely be inferred that this motion is not imaginary, but real.

EXAMP. I. The variation of distance between a ship and a remote object is owing to the motion of the ship, when the wind and current act and are competent for its production.

FIG.
XXV.

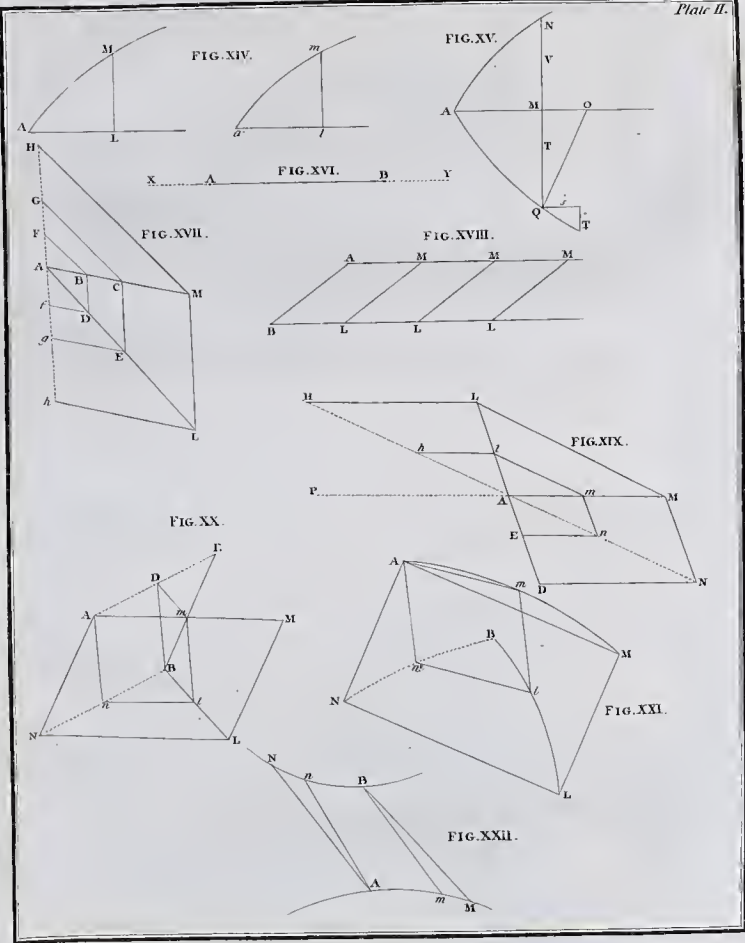
EXAMP. II. A ship sailing from south to north at the rate of five miles in an hour, observes another ship which sailed from the same point *E* going to the south-west at the rate of $7\frac{1}{10}$ miles in an hour nearly, and from hence it appears, that the last ship sailed from east to west at the rate of five miles in an hour nearly (139).

EXAMP. III. The magnitude of the earth's attractive power, or any other mechanical cause, is incompetent for the production of the diurnal revolution of the heavenly bodies, which is therefore probably apparent.

EXAMP. IV. The aberrations of the fixed stars are proved geometrically to result necessarily from the progressive motions of light and of the spectator combined, and are consequently only apparent.

EXAMP. V. A real circular motion is inseparably attended with a centrifugal force, and distinguishable therefore by this effect. Thus the diminution of gravity, which is observed in acceding towards the equator, affords a proof of the diurnal rotation of the earth.

C H A P.



C H A P. IV.

INERTIA OF MATTER.

162. DEF. *POWER* and force are used synonymously to signify any action upon a body, which produces motion, or a tendency to motion; as animal exertions and the influence of physical causes, gravity, elasticity, magnetism, &c. When a force acts always with the same intensity, or produces the same effect in a given time, it is said to be uniform or constant, and variable when it does not.

163. DEF. Any momentary impression upon a body producing motion, or a tendency to motion, is called an impulse, or percussion, of a power or force.

164. DEF. Whatever resists the action of a force is called an obstacle.

165. DEF. *Inertia*, or *vis inertiae*, of matter, is that quality by whose influence a change, in the direction or quantity of its motion, cannot be produced without the application of a force. The cause of a body's continuing in a state of rest, or of uniform rectilinear motion, is not any external force, but the nature and constitution of matter; and this internal cause or principle is called its inertia.

166. This property of matter, or inertia, is essential in the constitution of a system, whose preservation is necessarily dependent upon the regular observance of prescribed, determinate, laws; but
its.

its existence can only be collected from experience. *First*, Every change in the situation of a body, from rest to motion, and from motion to rest and to an encrease or decrease of motion, indicates an inertia, which is found to pervade every species of matter accessible to observation, the particular facts, adduced in proof of it, being innumerable and concurring to establish its universal existence. A quiescent body is never discovered to move, and motion, or change of motion, is never induced, without the actual impression of a cause able to produce these effects; for matter not only continues in a state of rest, by an incapacity to give motion to itself, but never ceases to move, or changes the quantity and direction of its motion, without the application of some philosophic influence adequate to their production, as magnetism, elasticity, gravity, &c. or animal exertion, as percussion, protrusion, &c. The time of the motion of a wheel upon its axis, or brass topp upon a polished surface, is encreased with the diminution of friction upon the axis, or surface; and the motion of a body, placed upon the deck of a ship, and partaking of its motion, continues in the same direction after the ship has ceased to move; and these motions are evidently discontinued by the operation of a cause, i. e. friction, competent to destroy them. New motions are observed without any sensible material impulse, resulting apparently from an innate tendency to motion; thus, a body not supported descends towards the earth, and, projected in a direction not perpendicular to the earth's surface, deviates from the line of projection with a velocity perpetually variable; and new motions arise amongst the minute particles of bodies; but these are the necessary result of established natural powers, gravity, elasticity, &c. *Secondly*, This quality is sometimes called the *vis insita*, or *vis inertiae*, of matter, from the similitude of its effects to those of animal powers exercised upon a body, both producing a change of motion. If a body *A*, placed upon a polished table, or suspended by a rope, be quiescent, it will continue in that state till urged by some external cause; and if another body *B* be projected with any velocity, and impinge upon *A*, it will communicate motion to it, and be itself retarded. When the effect of *B*'s impact

impact is considered in relation to the change produced in A , B is said to act upon A ; and, when considered in relation to the change of motion produced in itself, it is said to be resisted by A . Action is usually ascribed to a moving, and resistance to a quiescent, body; but they may both be considered as actions, because they produce similar effects, that is, a change of motion. A , by its vis inertiae, destroys a part of B 's motion, and B , by endeavouring to retain its present state, that is, by its vis inertiae, protrudes, and communicates motion to, A . *Thirdly*, Every change in the situation of a body is therefore universally allowed to indicate some philosophical influence, or animal exertion; and whether a body be quiescent, or moving, the quantity of its inertia is invariably the same; for it is discovered from experiments, the only source of information, that a force, communicating a velocity equal to a to the quiescent body A , will communicate an increase of velocity, equal to a , to that body moving. The force of gravity, which communicates a velocity nearly equal to thirty-two feet in one second of time, to a body falling from a state of rest, communicates an increase or decrease of velocity equal to this, in one second, to a body already moving, according as it conspires with, or is opposite to, the direction of the body's motion. And if the body B , moving with a velocity equal to b , impinge upon A at rest, and communicate a velocity to it equal to a , and lose a velocity equal to l , it appears from experiments, that B moving with the velocities $b + m$, $b + 2m$, $b + 3m$, &c. will communicate an increase of velocity equal to a , and lose a velocity equal to l , by impinging on A moving in the same direction with velocities equal to m , $2m$, $3m$, &c. Consequently the vis inertiae of A , estimated by the velocity, and increase or decrease, communicated to it, is the same in a state of motion and rest. *Fourthly*, As the whole vis inertiae of a body is composed of that of all its parts, and as we cannot conceive the vis inertiae of the same or equal particles to be increased or diminished, the whole vis inertiae of different bodies will vary as the number of equal particles, or quantities of matter, contained in them. The invariable result of experiment is, that bodies equal to

to A , $2A$, $3A$, &c. and heterogeneous bodies, whose weights are as 1, 2, 3, &c. receive the same velocity by the action of forces, whose magnitudes are as 1, 2, 3, &c.

167. DEF.* *The moment, quantity of motion, vis insita, of a body, are used synonymously to denote the impetus, or force, with which it moves, or its capacity to communicate motion or pressure.*

168. PROP. *If M , Q , V , represent moment, quantity of matter, and velocity, respectively, and be supposed variable, M will vary as $Q \times V$.*

DEM. Presuming that equal particles of matter have an equal inertia, it is evident that the pressure of one particle upon an immoveable obstacle, or the force requisite to generate or destroy its velocity, will vary as this velocity, or as V ; and because the impetus of any body is composed of that of every particle or Q , and Q and V are independent, M will vary as $V \times Q$ (64). Q. E. D.

Another demonstration.

The moments of different bodies vary as the uniform forces capable of generating or destroying them, in equal times, being their whole cotemporary effects; but forces equal to F , $2F$, $3F$, &c. will evidently communicate, in the same time, the same velocity to 1, 2, 3, &c. equal particles endued with an equal inertia, and velocities as 1, 2, 3, &c. to one particle. Therefore F , which is as M , is increased in the same ratio, with the velocity or V , and number of equal particles or Q , and, these being independent, vary as $Q \times V$ (64). Q. E. D.

* Keil's Physics, Lect. IX.

SCHOLIUM.

169. These, or similar demonstrations, are usually adduced in support of this proposition, which is perhaps more properly and with more conviction, demonstrated by experiments. The moments of different bodies are collected by measuring the magnitudes of the forces required to produce them, or the magnitudes of their cotemporary effects similarly produced, in the simplest and most intelligible instances, which may be deemed the common measure of their moments.

EXP. I. If two unequal spherical bodies, moving in opposite directions, meet and after impact be quiescent, their moments must be equal; but their quantities of matter are always found to be inversely as their velocities. Or, if they move in the same direction, and one overtakes the other, the velocity gained by the struck body, and lost by the striking body, are always found to be inversely as their quantities of matter.

EXP. II. If a body A be placed on the same side of the fulcrum of a straight lever, at the distance of 1, 2, 3, &c. feet from A , it will move with velocities as 1, 2, 3; and be restored to an equilibrium by the bodies A , $2A$, or $3A$, &c. placed, at the distance of one foot, on the other side. And, in general, a body, whose weight is $n \times A$ pounds, at the distance of one foot from the fulcrum, is found to equilibrate with bodies, whose weights are $\frac{n \times A}{2}$, $\frac{n \times A}{3}$, $\frac{n \times A}{4}$, &c. at the distance of 2, 3, 4, &c. feet from the fulcrum respectively. The pressures or moments, are justly inferred from these experiments, to be as their weights multiplied into their distances, or as the quantities of matter multiplied into their velocities. And the same conclusion results from experiments upon every other, however complicated, machine, when allowance is made for friction, and the velocity is properly estimated.

170. Cor. 1. Since M varies as $\mathcal{Q} \times V$, V will vary as $\frac{M}{\mathcal{Q}}$, and \mathcal{Q} as $\frac{M}{V}$ (55). If lines be taken in the ratio of \mathcal{Q} to V , M will vary as the area of a rectangle, whose sides are these lines; and if numbers be taken in that ratio, M will vary as their product.

171. Cor. 2. If S represent the space described, in the time T , with the velocity V , M , varying as $\mathcal{Q} \times V$, will vary as $\frac{\mathcal{Q} \times S}{T}$ (105), S , varying as $V \times T$, will vary as $\frac{M \times T}{\mathcal{Q}}$, and T will vary as $\frac{S \times \mathcal{Q}}{M}$.

172. Cor. 3. If two bodies therefore A and B , moving in opposite directions with velocities respectively equal to a and b , meet, and after impact be quiescent; or if their effects to produce motion upon any machine be opposite and equal, and their velocities, when properly estimated, be a and b ; $A \times a = B \times b$, and $A : B :: b : a$ (38), or A and B are to each other inversely as their velocities. The converse of this is true, and if the bodies be inversely as their velocities, or $A : B :: b : a$, their moments are equal, or they are in equilibrio upon any machine, and $A \times a = B \times b$ (37).

SCHOLIUM.

173. Forces are distinguished by some foreign writers into two kinds; 1st, of bodies at rest, and 2dly, of bodies in motion.

First, *The force of a quiescent body, such as is conceived to reside in one lying upon a table, suspended by a string, sustained by springs, &c. is called its pressure, tension, force, solicitation, vis mortua, &c. and is estimated by the weight with which the table is pressed, the string stretched, spring bent, &c. To this kind are referred centripetal and centrifugal forces, as their effects are similar to those of weights, pressures, tensions.*

Secondly,

Secondly, *The force of a moving body, arising from its inertia, by which it communicates motion, or a change of motion, to another body by impact, overcomes gravity, friction, and other pressures, and is only destroyed by an extinction of its motion, is called the vis motrix of that body, or its vis viva, to distinguish it from the vis mortua.*

Concerning the measure of the first of these forces no doubts are entertained, the pressure of one particle, upon the arm of a lever, being universally allowed to be as its velocity, and of a number of particles, as that number and velocity conjointly; but the measure of the vis viva hath long been a subject of warm contention between the adherents of Newton and Leibnitz, the former maintaining it to be the product of the quantity of matter and velocity, and the latter the product of the quantity of matter and square of the velocity. The term force being defined to be that, which, acting upon a body, communicates motion, or a tendency to motion, the magnitude of the force of a moving body, at any instant of time, may be best, and indeed can only be clearly, conceived by measuring the motion communicated by it; and in this mensuration three things are to be considered, viz. the magnitude of the body or quantity of matter to be moved, the velocity communicated, and the time in which it is communicated. Some of these circumstances may be, as they have been, omitted, and different measures of the same force have resulted. The followers of Newton conceive forces to act, for the same time, with their respective intensities continued uniformly, and estimate their magnitudes by their cotemporary effects. If two forces F and f be supposed to act, at the same instant, upon two bodies A and B , with uniform intensities, and communicate velocities to them, equal to V and v respectively, in the same time, the magnitudes of these forces are to each other as $A \times V : B \times v$. And in the communication of motion by the impact of one body upon another, this law is found to obtain, wherever the effects are distinctly understood, and can be ascertained, the quantity of motion, if measured by the product of the quantity of matter and velocity, remaining invariable, when estimated in the same direction. But the followers of Leibnitz, adopting a different definition of force, derive a different conclusion: they do not suppose the force to

act with uniform intensity, as it may decrease gradually to evanescence, which happens in the collisions of bodies, and actions of springs; and the time of action, and sometimes the direction, is disregarded, and not deemed to affect the result. Though it be true in some particular cases of the communication of motion from one body to another, that the force, according to their acceptance of the term, varies as the product of the quantity of matter and square of the velocity, or that this product is the same before and after impact, this conclusion cannot be affirmed to obtain generally. The doctrine advanced by Newton is universally true, according to his meaning of the term force; but whether the opinion of Leibnitz be true or not, is best known from experiments, the result of which is generally repugnant to it: and it would therefore, perhaps, create less confusion to adhere to the old definition of force, include the time of action, and suppose the intensity of the force, during that time, to be the same. The distinction between these two kinds of forces hath often been urged to be superfluous, for the following reasons. *First*, The force of a moving body is a vague and undefinable term; for there is no force in a body considered absolutely, except its inertia, which is always the same, whether the body be quiescent or moving; but if a moving body impinge upon another body either moving or quiescent, its inertia exerts itself as a force, whose magnitude is relative and depends upon its velocity, and the magnitude and velocity of the struck body. If the body A , moving with the same velocity a , impinge upon the quiescent bodies B , $2B$, $3B$, &c. or upon the same body B quiescent, and moving with different velocities equal to b , $2b$, $3b$, &c. it will in every case produce a different effect, and the change of state, both in the impinging and struck body, will be different. The force of springs, and animal exertions, is also relative, and it seems therefore improper to talk of the absolute force of a moving body, springs, &c. because it is relative. *Secondly*, No idea can be formed of an instantaneous communication of motion; for, in the collision of bodies, the parts which come into contact first, are displaced and move, and some time elapses before motion be communicated to the whole body. This interval of time between the first contact of the nearest parts of the impinging

ing bodies, and their motion, is the time of collision, in which the bodies exert a mutual, and perpetually variable pressure, by which their state is gradually changed. This is most observable in soft bodies, or those whose parts yield to percussion; and, as the parts of all bodies do yield, obtains universally, in a greater or less degree, according to the different contexture of their parts. The forces of percussion are therefore pressures, exerted with variable intensity for a short time; and, to comprehend their magnitude, this time ought to be ascertained, the quantity of pressure, for every instant of this time, computed, and the whole collected into one sum. This discrimination of forces seems therefore to be unnecessary, because the action of collision is nothing more than the operation of a continued pressure, and pressures and percussions are similar, and properly expressed by the same term, force. The arguments adduced, in support of their opinion, by the advocates of Leibnitz, are derived from experiments upon, the collision of bodies, the action of elastic springs, the composition and resolution of motion, the indentings or cavities formed in clay, tallow, &c. by falling bodies, and the velocity of fluids in hydraulics. I will only produce a few experiments, which may convey some useful knowledge, are most strongly urged in support of this hypothesis, and appear most to militate against the old opinion.

EXP. I. If one spring be requisite to destroy the motion of a body, whose velocity is 1, four springs, equal to it, are requisite to destroy its motion when its velocity is 2, nine springs when its velocity is 3; and generally the number springs varies as the square of the velocity of the moving body.

EXP. II. The number of men, horses, &c. of equal strength requisite to communicate velocities, as 1, 2, 3, &c. to any body, are as 1, 4, 9, &c. or as the square of the velocity.

These similar experiments may be true, and indeed seem to be established.

established by Gravesand, Desaguliers, &c.; but they prove nothing against the common opinion of the force of bodies, which is not measured by effects produced in any unequal times. And besides, the effect of any number of agents, as springs or animal exertions, cannot properly be said to measure their force; because their whole force is not exerted, that of each, considered individually, being greater than when acting in conjunction with the rest, and will be greater or less according as the number of the rest is less or greater.

EXP. III. If the velocity of water, striking upon the float-board, or ladle-board, of a mill, be as 1, 2, 3, &c. the effect of the wheel is found to be as 1, 4, 9, &c. or as the square of the velocity of the stream; and the force of the water, being measured by its effect, varies therefore as the square of the velocity. In this experiment the number of particles of water striking the board in a given time, encreases as the velocity, and, the force of each being encreased in the same ratio, the effect ought to be as the square of the velocity. Theory is very seldom confirmed by practice without allowances for the operation of collateral causes, which, perhaps, cannot be made with demonstrative accuracy. Such experiments as this are too imperfectly understood to justify any general conclusions against experimental proofs clearly comprehended, or theoretic demonstrations, resulting from data, which, perhaps, do not take place in the experiment.

EXP. IV. Mr. Smeaton has proved, by some very ingenious experiments (Phil. Transf. Vol. LXVI. pag. 469.), that the mechanic power to be expended, varies as the square of the velocity to be generated, and the velocity as the impelling power multiplied into the time of action. But the mechanic power, according to Mr. Smeaton, is measured by the weight multiplied into the space descended, and from this definition, his conclusion is inevitable: but this power is different from what is usually meant by moment, force, &c. In N^o. 2, 3, the weight *S* descends through
spaces

spaces equal to 10, and $2\frac{1}{2}$, turns, in $28\frac{1}{4}$ " and $14\frac{1}{4}$ " respectively, and communicates velocities nearly as 2 : 1 ; and in this experiment the weight S acts with the same intensity for times, whose ratio is 2 : 1, but its capacity to communicate motion, if measured by the space, must be as 4 : 1.*

EXP. V. Equal cavities are formed in clay, tallow, &c. by equal bodies falling through spaces inversely as their quantities of matter; but the cavities being equal, the causes, or forces of the bodies, are equal, or as $\frac{S}{S}$ (supposing S to represent the space) or as $S \times Q$, from the supposition, or as $Q \times V^2$.

EXP. VI. A ball shot with velocities as 1, 2, 3, &c. is found to form cavities in wood, clay, &c. whose depths are as 1, 4, 9, &c. or as the squares of the velocities.

The two last experiments prove nothing against the common doctrine; and though they are of practical utility, yet admitting the propriety of measuring forces by their whole effects produced in any unequal times, they do not vary, in these experiments, as the square of the velocity: because the cavities formed are not the whole effect, a greater or less motion being communicated to the parts contiguous to the cavity in the struck body, according to the magnitude and velocity of the striking body. A small body, moving with a great velocity, may by impact overcome the cohesive force of the particles of the struck body, and effect a separation between them; and from its great velocity, or short time of action, little motion will be communicated to the contiguous parts; but

* These experiments are immediately deducible from mechanical principles. Let $F =$ the radius of the barrel M ; $V =$ velocity generated in K in the time T ; X and $Z =$ the respective spaces passed through by S and K in the time T ; and $1 = K$'s distance from the axis Be . And from established principles, the force impelling K is .8 oz. $\times F$, or as F , and, because F is constant, $F \times T$ is as V . $F \times T^2$ is also as Z ; but $X : Z :: F : 1$, therefore X (or the space descended by S , which is as the mechanic power) $= Z \times F$, and is as $F^2 \times T^2$, or as V^2 . (Phil. Transf. Vol. LXVI. p. 469.)

but a great body, moving slowly, will form no cavity in the obstacle; because the moment of those parts of its surface in contact, is less than the cohesive force of the parts of the obstacle, and the parts struck will, by their unknown connection with the parts adjoining, communicate motion to them, and these to the next, till the whole obstacle move. Two bodies therefore moving with the same moment, may produce quite dissimilar effects; and as lines, surfaces, &c. are measured by the application of a common measure, the magnitudes of forces are to be measured by their effects, which ought to be similar in every respect, and this similitude is most observable and best understood in experiments upon the lever and other simple machines.*

L A W S O F M O T I O N.†

174. The laws of motion are general rules, or consequences, resulting from the nature and constitution of matter, and its relation to mechanical forces, the causes of motion, by which it is governed, and from which it cannot deviate without suffering a total change of its nature. Their truth is established by a number of concurring and uncontroverted observations, all information upon natural powers and motions derived from them being only deducible from thence; and, when thus established, they are esteemed to be uniform characteristic marks obtaining in all material motions.

FIRST

* Let s represent the effect, or space through which any obstacle is impelled by a moving body, v be the velocity of the body: and if s be as v^2 , \dot{s} is as $2v\dot{v}$, or as $v\dot{v}$, or as $\frac{\dot{v} \times s}{t}$ (if t be the time), and \dot{v} therefore is as \dot{s} . The decrement of velocity is therefore as the time in which it is produced, and the resistance must be constant; which seems to be a necessary property of the vis viva; but as this seldom happens, and the vis viva is really a pressure, and may be measured in the same manner with the vis mortua, the term force may express them both.

See Defagulier, Vol. II. p. 51. Maclaurin's Account of Sir I. Newton's Discoveries, p. 178. Discours sur les Loix de la comm. du Movement, Oper. Tom. III. and Dissertat. de vera notione virium vivarum. Act. Petropol. Tom. I. p. 131. Muschen. Int. ad Phil. Nat. Vol. I. p. 83, &c. Phil. Transf. No. 371, 375, 376, 396, 400, 401, 459.

† Maclaurin, Book II. Chap. II. Newt. Prin. p. 13. Keil's Physics, Lect. XI, XII. Helsham, Lect. III. IV.

LAWS OF MOTION.

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FIRST LAW OF MOTION.

175. *Every body perseveres in a state of rest, or of uniform rectilinear motion, unless affected by some external cause, as animal exertion or physical power.*

The truth of this law is collected by observation. It is nearly synonymous to the inertia, or that quality, of matter, by which it continues in every new state; for, this being established, it is a necessary result that it cannot begin to move from a state of rest, nor, if in motion, can it accelerate, retard, or change the direction of that motion.

176. Cor. Every body, moving in a curve line, or in a right line with a perpetually varying velocity, is acted upon without intermission during its motion, by some external force; and every body, moving in a curve line, has a perpetual tendency to move in a right line, and to leave the curve in the direction of a tangent to that point of the curve where the body is.

SECOND LAW OF MOTION.

177. *Motion, and change of motion, are proportional to, and in the direction of, the force impressed.*

The motion communicated to a quiescent body is proportional to the cause producing it; and when a body is accelerated or retarded, during its motion, the acceleration and retardation are proportional to the force producing them. If a body move upon an inclined plane AB , or in any curve line LM , by the action of a force tending to any point S , the change of motion is not proportional to the whole force directed to that point, but to that part which acts in the direction AB , or NT touching the curve where the body is; for this part, only, accelerates or retards the

K

body's

FIG.
XXVI.

LAW S O F M O T I O N.

body's motion. When water or air act upon the vanes of a mill, or sails of a ship, the change of motion is not proportional to the whole force of the water or air, but to that part which is actually impressed upon them; for, if the velocity of the water or air be equal to that of the vanes or sails, they will move on together without any acceleration; and, if the direction of the air or water be oblique to the motion of the vanes or ship, some force will be ineffectual, and the change of motion will be proportional to the remainder, acting in the direction of their motion.

The truth of this law is demonstrated by numberless experiments upon the motion, and change of direction and motion, communicated by material impulse. Every experiment upon the effects of conspiring and opposite forces, and the effects of forces, acting obliquely, or resulting from the composition and resolution of forces, demonstrated experimentally, exhibit concurring proofs of its truth; for were motion, and change of motion, not proportional to, and in the direction of, the force impressed, it may be proved geometrically that these effects would not take place. The conclusion is therefore inevitable, and, being thus established by uniform experience, may be deemed a fixed principle in nature, and applied not only to the explication of those particular experiments from whence it was deduced, but of all other natural phenomena. This law is also a proof of the inertia of matter, or, that being established, a corollary from it; for if matter can produce no change in itself, it must move in the direction of the force impressed; and, since the *vis inertiae* of a body is the same, whether moving or quiescent, it is evident that the genesis, encrease, and decrease, of velocity equal to V in any body, require exactly the same force. If the force of collision, or impulse of any force such as gravity, be capable of generating a velocity, equal to V , in the body A , it will in the same time generate an encrease, or decrease, of velocity equal to V in that body, according as it conspires with, or opposes its motion.

178. Cor. 1. The velocity generated by one single impulse equal to nF , n being a number, is equivalent to that generated by n successive

cessive conspiring impulses, each of which is equal to F ; and consequently the velocity communicated to a body by any number of forces acting in the same direction is the same, whether they act together or separately; because the genesis, and encrease, of a velocity $= V$, require exactly the same force.

179. Cor. 2. The velocity generated, in any body, by any number of unequal conspiring impulses $X, Y, Z, \&c.$ is as their sum; for let $x, y, z, \&c.$ be the velocities respectively generated by them in any body A , and (177) $X:Y::x:y$, and $X:X+Y::x:x+y$, and $X:X+Y+Z::x:x+y+z, \&c.$ And the velocity, communicated to any body by a force $= F$, in a given time, is as the magnitude of that force; for supposing this force to act by impulses, whose magnitudes are $X, Y, Z, \&c., F = X+Y+Z, \&c.$

180. Cor. 3. The magnitudes of any forces are therefore as the spaces uniformly described in a given time, by the velocity which they communicate to the same body (106).

181. Cor. 4. The velocity generated in a body in the same time by two forces X, Y , acting in opposite directions, is as their difference; for (177) $X:Y::x:y$ and $X:X-Y::x:x-y$.

T H I R D L A W O F M O T I O N.

182. *Action and reaction are always equal and opposite; or the mutual actions of two bodies are always equal, and in opposite directions.*

This is another rule observable in all the motions of nature, resulting, like the two first, from the inertia of matter. Were matter divested of this property, motion would be communicable without resistance, and consequently without effecting any change in the force communicating it. The action of all forces, whether operating by the visible impact of one body upon another, or the

invisible agency of gravity, magnetism, &c. consists in producing pressure and motion; and reaction in supporting this pressure, or resisting the production of motion. And these this law asserts to be equal, when estimated in opposite directions; which, being proved experimentally, like the two first, in all communications of motion with which we are acquainted, may be esteemed a general principle pervading the whole material system.

In the communication of pressure upon any immovable plane, whether arising from the protrusion, gravity, or force of impact, of a body; the meaning of the law is, that the resistance of the plane, and an opposite force equal to that producing the pressure, have precisely the same effect, as they only destroy the force of protrusion, weight, or impact. In the communication of motion by visible impact, the meaning of the law is, that action is mutual, equal and opposite, or that the quantities of motion lost and gained, which measure this action and reaction, are equal when estimated in opposite directions. It is always supposed, that the quantity of motion is estimated by the product of the quantity of matter and velocity. No idea can be formed of the loss of motion, except by communication; and that the quantity, lost by the impinging body, is gained by the struck body, appears from innumerable experiments upon the collision of bodies. If a perfectly hard, or a soft body, A , moving with the velocity V , impinge upon an equal unelastic body B , they will move together, after impact, with a velocity equal to $\frac{V}{2}$; and if A , moving with the velocity $V + u$, impinge upon B , moving in an opposite direction with the velocity u , they will move together, after impact, with a velocity equal to $\frac{V}{2}$. And whatever be the magnitudes of A and B , it is the invariable result of experiments, that the quantity of matter in A , multiplied into the velocity lost by it, is equal to the quantity of matter in B , multiplied into the velocity which it gains by impact; or if l and g represent the velocities lost by A and gained by B respectively, $A \times l = B \times g$. When A and B are perfectly elastic, the velocities lost and gained are doubled, and the law still obtains, $A \times 2l$ being equal to $B \times 2g$. If motion be communicated to any body A , by
any

any force X protruding, pulling, &c. the meaning of the law is, that the reaction, or resistance, of A destroys such a part of X 's force, as would be destroyed by a similar force, capable of generating that motion in A , acting in opposition to X . In the communication of motion by unknown means, as by gravity, magnetism, repulsion, &c. the law asserts that the body, attracting or repelling, moves in an opposite direction to that of the body attracted or repelled, and with an equal quantity of motion. The attraction between the earth, and any body upon its surface, is mutual and equal; for, were it not, a rectilinear motion would ensue from the stronger attraction, which is contrary to experience. And, since gravity is an innate principle, it, and its effects will remain the same when that body is detached from the earth, and consequently their attractions continue to be mutual and equal, and they will meet with equal moments. If two magnets A and B , whose weights are unequal, be placed upon two pieces of wood, floating in water within the reach of attraction, they will meet with velocities inversely as their quantities of matter; and, if a reed be inserted between them to prevent their junction, they will be quiescent, which they would not be, were their attractions unequal. When the weight of A is equal to 2, 3, 4, &c. times that of B , its velocity, in an opposite direction, is equal to $\frac{1}{2}$, $\frac{1}{3}$, &c. of B 's, the products of their weights and velocities being always found to be equal. And if A is repelled, B is also found to be repelled in an opposite direction, and their velocities are always inversely as their weights. This law, being found to obtain in all actions of bodies within the reach of experiments, is inferred to obtain universally through the material system.

183. Cor. In the impact of bodies therefore the quantity of motion, estimated in the same direction, is the same before and after impact. If A , moving with a velocity equal to a , impinge upon B , moving in the same, or an opposite, direction, with a velocity equal to b , the sum of their moments after impact, is equal to $Aa + Bb$, or $Aa - Bb$, according as they move in the same, or opposite directions; for, if they move in the same direction, the en-crease

crease of moment communicated to B is destroyed in A , and their sum continues the same; and, if they move in opposite directions, the least moment, and a part of the other body's moment equal to it, are destroyed, so that their sum after impact continues to be equal to $Aa - Bb$, the same as before impact.

COMPOSITION AND RESOLUTION OF FORCES.

PLATE
III.
FIG.
XXVII.

184. * The same motion may be communicated to a body B by a single force Z , and any number of conspiring forces, into which it is said to be resolvable, and they are said to be compounded into Z ; thus the body B may describe the line BP , in the same time, by a force Z acting in that direction, or by two, three, &c. forces X , Y , &c. inclined to it. For it is evident, that B , acted upon at the same time, by two forces X , Y , in directions making an angle, which cease to act after the body has left the point B , will move in some intermediate rectilinear direction, and this, by changing the inclination and magnitude of X and Y , which are variable without limit, may be any intermediate line whatever BP . The force Z is said to be resolvable into the two, X and Y , producing the same effect with it; and, vice versâ, X and Y are said to be compounded into one, Z , which produces the same effect with them.

FIG.
XXVIII.

185. PROP. *A body B urged at the same time, by two forces x , y , whose action ceases when the body has left the point B , and whose magnitudes and directions are as two right lines BM , BN , making any angle, will move as if it were impelled by one force z , whose magnitude and direction are as BP , the diagonal of a parallelogram, whose sides are BM and BN .†*

DEM.

* Newt. Princip. p. 14. Muschen. Ch. X. Hellsam, Lect. IV. Maclaurin's Newt. Ch. II. Graves. Lect. I. Ch. XIII. Hermannus, Lect. I. Ch. II. Varignon. Tom. I. Sect. I.

† The magnitudes of forces, whether animal exertions as percussions, protrusions, &c. or philosophic powers, as gravity, magnetism, &c. can only be compared by measuring their effects, which are supposed to be produced uniformly in the same time. And, in this supposition, it is evident, that the conclusions will be the same, whether the forces X , Y be impulses, or the same number of impulses always equal to themselves.

DEM. The access of B to, or recess from, MP , is not affected by the action of y ; nor its access to, or recess from, NP , by the action of x (second law of motion): therefore they will carry B to MP and NP respectively, in the same time, whether they act together or separately. But if they act together, B , being carried to both MP and NP , must be found at P their intersection, and describe BP uniformly (first law of motion); which is the magnitude and direction of z (180). Q. E. D.

Otherwise,

Let BM , moving parallel to NP uniformly, arrive at NP in the same time that B describes BM uniformly, which consequently at the end of that time will arrive at P . And if npm , be any new position of BM , and the body at p ; $Bn : np ::$ velocity of $BM : \text{velocity of the body} :: BN : BM$ (106); therefore p is always in the diagonal BP ; Bp encreases uniformly because Bn does; and BP is the magnitude and direction of z (180). Q. E. D.*

186. Cor. 1. A force z , producing the same effect with two x and y , which are as BM, BN , is as twice Bs , supposing s to be the bisection of BP .

187. Cor. 2. If any two forces act upon a body, in directions which make an angle, it cannot be quiescent.

188. Cor. 3. A force represented by a side of a triangle BP , may be substituted for two represented by the remaining sides BM, MP ; and it is said to be equivalent to them, because they produce

* This proposition, which is deduced from the second law of motion, may be demonstrated not inelegantly from the third. The forces x and y , urging the body B obliquely, partly oppose each other, and since, from the third law of motion, the opposite parts are equal, the path of B will be equally distant from any cotemporary positions of it, M and N , when they act separately; which therefore is Bs , s being the bisection of BP . And because each of the forces x, y carry B through Bs , they will both carry it through $2Bs$ or BP , the diagonal of a parallelogram, whose sides are BM and BN .

duce the same effect, or make a body describe BP in the same time. And these three forces BM, BP, BN are in the same plane (Euc. B. XI. P. II.)

FIG.
XXVIII.

189. Cor. 4. A force BM is equivalent to BA and AM ; and BN to BC and CN ; and therefore these forces BA, AM, BC, CN , are equivalent to, and may be substituted for, the force BP . And any number of forces, whose directions are in the same plane, may be reduced to one in that plane producing the same effect with them.

FIG.
XXIX.

190. Cor. 5. A force AH is equivalent to two represented by AB and AG ; and AD to AH and AF ; therefore any number of forces $AG, AB, AF, \&c.$ in different planes, may be reduced to one AD producing the same effect with them.

FIG.
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191. Cor. 6. The converse of this is true, and one force BP , or AD , may be resolved into any number, either in the same, or different planes, producing the same effect with it.

FIG.
XXX.

192. Cor. 7. The velocity generated by z : velocity generated by $x + y :: BP : BM + BN$; and the quantity of motion is therefore diminished by composition, and increased by resolution; but, when estimated in any given direction BL , it remains invariable. For let BP, BM, BN , be each resolved into two, one in the direction BL , and the other perpendicular to it; and from similar triangles

$$Py : Dy :: yL : yn,$$

$$\text{and comp. } PD : nL :: Dy : yn :: BE : Bm,$$

$$\text{therefore } BL = Bm + Bn.$$

193. Cor. 8. The same force z , and consequently the same velocity, may be generated by an infinite number of pairs of forces, because

because the same right line may be the diagonal of an infinite number of parallelograms.

194. Cor. 9. The parts, of the two forces, directly opposed to each other, are to the parts, acting in the diagonal, as $Mm + Nn$ to $Bm \pm Bn$, according as the angle MBN is not greater, or greater than a right angle. If x and y are given, and the angle MBN be supposed variable, and become equal to two right angles, then N will come to N' , P to P' , s to s' (s always bisecting BP and MN) and $2Bs'$, or the conspiring parts of x and y , $= BP' = BM - BN$; if this angle vanish, then N will come to N'' , P to P'' , s to s'' , and $2Bs''$, or the conspiring parts of x and y , $= BP'' = BM + BN$.

FIG.
XXXI.

FIG.
XXXII.

195. Cor. 10. If BP and the angle BMP be given, the curve line passing through M , in different positions of M , is a circular arc subtended by BP ; if BP and $BM + MP$ be given, this curve line is an ellipse, whose focal distance is BP ; and if BP and the difference of BM and MP be given, the curve is an hyperbola and BP its focal distance.

PLATE
IV.
FIG.
XXXIII.

196. PROP. Any three forces x, y, z , whose directions are Bx, By, BN in the same plane, acting upon a body B without producing motion, are to each other as the three sides of a triangle respectively parallel to their directions.

FIG.
XXXIV.

DEM. Let BN be the magnitude of z , and be resolved into two, in the directions yB, xB ; viz. BM, BP ; which must be respectively equal to x and y , because in equilibrio with them, and acting in the same directions; therefore x, y, z , are to each other as $BM, BP (MN), BN$ respectively. Q. E. D.

197. Cor. 1. x, y, z are to each other as the three sides of a triangle respectively perpendicular, or equally inclined to their directions; for this is similar to the former.

198. Cor. 2. Any two of these forces are to each other inversely as the sines of the angles, which their directions make with the direction of the third. For,

$$\begin{aligned} z : y &:: BN : BP :: \sin. \angle BPN \text{ or } PBM : \sin. \angle BNP \text{ or } NBM, \\ z : x &:: BN : BM :: \sin. \angle BMN \text{ or } MBP : \sin. \angle MNB \text{ or } NBP, \\ \text{and } x : y &:: BM : MN :: \sin. \angle BNM \text{ or } NBP : \sin. \angle MBN. \end{aligned}$$

199. Cor. 3. The angle $\angle xBy$, and magnitude of x and y , being given, the magnitude and direction of z may be found; for BM , MN and the angle BMN , the supplement to the angle $\angle xBy$, being known, BN , and the angles MBN and MNB or NBP , may be found by trigonometry.

FIG.
XXXV.

200. Cor. 4. Any number of forces in the same plane, acting upon a body B which remains quiescent, may be reduced to two equal and opposite: if BM , BN , BD , act upon B , they are equivalent to BP and BC in opposite directions, which must be equal, because B is quiescent.

201. Cor. 5. Any number of forces in different planes, urging a body without producing motion, may be reduced to two in the same plane, equal and opposite. If BD and DC be elevated above the plane BNM , and PBC be the common intersection of the planes BMN and BDC ; BP , which is equivalent to BM and BN , must be equal to BC , which is equivalent to BD and DC , because they are opposite and destroy each other.

202. Cor. 6. If more than three forces BM , BN , and BD , each of which is combined from two, act in different planes, upon a body B without producing motion, and are estimated in any direction BL by letting fall parallel lines upon it, the forces resulting will produce no motion in that direction; for the parts of BM , BN , resulting from this resolution, which are in the direction BL , are equal and opposite to the part of BD reduced to the same direction, and the sums of the opposite parts, on each side of BL , must also be equal, because B is supposed to be quiescent.

203. Cor.

203. Cor. 7. If the forces BM, BN, BC , which act upon a body without producing motion, be reduced to any plane FH , by resolving each into two, one perpendicular to it, and the other parallel to the lines joining the points, where the perpendiculars meet it, the forces resulting will produce no motion in that plane. For BP , being taken equal and opposite to BC , will be equivalent to BN and BM , and BC, BP will be projected into equal straight lines bc, bp ; but BN, MP , being equal and parallel, are projected into equal and parallel lines bn, mp , and the figure $bmpn$, is a parallelogram. And bm, bn , being equivalent to bp , will be equivalent to its equal and opposite bc .

FIG.
XXXVI.

204. PROP. *If a body be acted upon by two similar variable forces, in directions parallel to BP, BQ , making any angle, which act when the body hath left the point B , it will describe a right line.*

FIG.
XXXVII.

DEM. Let the forces act by impulses at the beginning of equal particles of time, and let BD, DE, EF , and BG, GH, HI be the relative magnitudes of corresponding impulses. By the action of the two first impulses BD, BG , the body will move in the direction BK (185); by the two next DE, GH , it will describe, in the same direction, KL , &c.; and, because the forces are similar, $BD : DE :: BG : GH$, and componendo $BD : BG :: BE : BH$, and consequently BKL is a right line. Q. E. D.

205. Cor. 1. Since this demonstration does not depend upon the magnitude of the particles of time, into which the whole time is supposed to be divided, it will obtain when these particles are evanescent and the forces act incessantly.

206. Cor. 2. If the forces be constant, the velocity in BM is uniformly accelerated; for, in this supposition, BD, DE , &c. are equal, and $BD : DE :: BK : KL$; but BK, KL which are therefore equal, &c. measure the increments of velocity communicated at B and K (106).

207. Cor. 3. Every body moving in a curve line, must be acted upon by, at least, two dissimilar forces.

208. Cor. 4. The whole force in the direction BP is the sum of the impulses $BD, DE, EF, \&c.$ and the force in the direction BM , is the sum of the impulses $BK, KL, LM, \&c.$ and these forces are similar; for $BD, DE, \&c.$ measure the intensity of the force in the direction BP at the points B and D ; and $BK, KL, \&c.$ measure the intensity of the force in the direction BM , at B and K , and $BD : DE :: BK : KL$.

FIG.
XXXVII.

209. PROP. 86. *If a body be urged at the same time by two constant, or similar variable, forces x, y , whose magnitudes and directions are the two sides of a parallelogram BF, BI , it will move in the same manner as if it were urged by a constant or similar variable force z , represented by the diagonal.*

DEM. The impulse BK is equivalent to BD and BG ; KL to DE and GH ; LM to EF and $HI, \&c.$ consequently the sum of these impulses BM is equivalent to BF and BI . Q. E. D.

210. Cor. 1. A force represented by one side of a triangle BM , may be substituted for two similar forces represented by the remaining sides BF, FM .

211. Cor. 2. A force represented by BM may be reduced into any number of forces BR, RS, SM , similar to it, and vice versâ.

212. Cor. 3. Three similar variable forces, urging a body without producing motion, are to one another as the three sides of a triangle, parallel or perpendicular to their directions; and they are to each other inversely as the sines of the angles, which their directions

directions make with the direction of the third (198); and any number of similar forces in the same or different planes, in equilibrio, may be reduced to two, in the same plane, equal and opposite.

S C H O L I U M.

213. Prop. (204) may be demonstrated per saltum, though perhaps less intelligibly, by the following process. Let a body B be urged, at the same time, by any forces x and y , in the directions BM, BN , which carry it respectively through Bm, Bn and BM, BN , in equal times; and it is evident that, at the end of those times, B will be in d and D . If x and y act only at the point B , or if $Bm : BM :: Bn : BN$, the path of B is the diagonal; if they act at B, m, n, M, N , &c. and Bm be to BM as the times of description, and Bn be to BN as the squares, cubes, &c. of those times; $Bn : BN :: Bm^2 : BM^2$ or $Bm^3 : BM^3$, &c.; and the path of B is a parabola, &c.

FIG.
XXXVIII.

214. PROP. 87. *A body B impelled by two forces tending to two equidistant points, Q, R, which at equal distances are equal, and at unequal distances are variable according to any law whatever, will describe a right line.*

FIG.
XXXIX.

DEM. Let these forces act by impulses at the beginning of equal particles of time, and BE, BF being the magnitude of the two first, will make B describe BD , which produced bisects the base QR ; and DL, DM , the two next, will make it describe DN , bisecting the base QR : therefore BN is a right line; and the demonstration obtains when the particles of time are evanescent, and the action of the forces incessant. Q. E. D.

215. Cor. The direction of the combined attraction of all the particles, composing a sphere, passes through the center: for let it be divided into thin laminæ parallel and contiguous to each other,
and

and let Q, R , be any corpuscles in a lamina, equidistant from PB passing through the center, and what has been proved of these two may be proved of every two equidistant from P , which compose the whole lamina; and the same may be proved of every lamina, and consequently of the whole sphere.

S C H O L I U M.

PLATE
V.
FIG. XL.

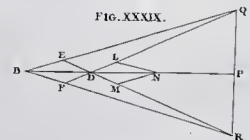
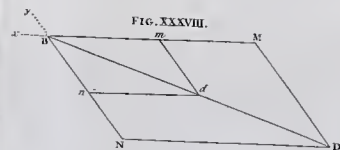
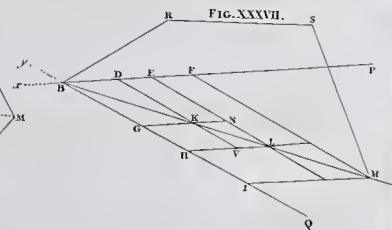
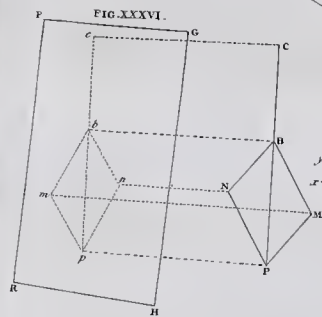
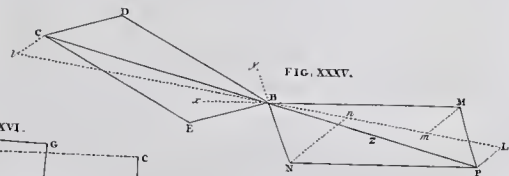
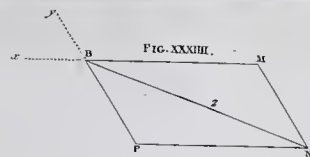
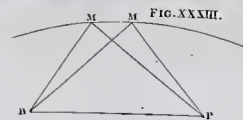
FIG.
XLI.

FIG.
XLII.

216. Three forces, acting upon a body without producing motion, are to each other as the three sides of a triangle parallel to their direction, which consequently meet in the same point of the body, and are in the same plane (EUC. B. XI. P. II.). This may be evinced otherwise; for since a body, acted upon by any number of forces, cannot be at rest, unless they oppose each other, and the opposite parts be equal, and in the same direction; it is clear that a plane AB may pass through P , acted upon by x, y, z in different planes, so that they shall be on the same side of AB , and therefore, as they partly conspire, they must communicate motion to the point P , and to the whole body. And if any two of these forces, x, y , act against the same point p of the body AB , and x acts against q ; x and y will communicate motion to p (187) and z to q , and AB cannot therefore be quiescent. If the directions of x and y meet in a point p without the body, and z act at the point q , AB cannot be at rest; because x and y have the same effect with, and are equivalent to, some force in an intermediate direction pd , which is not opposite to the direction of z , and cannot therefore destroy it. It is supposed, that the directions of these forces are inclined to each other; for, if they be parallel, BA may be quiescent, when they act upon different points.

FIG.
XLIII.

217. PROP. *If a body B be acted upon by any number of forces, at the same time, whose magnitudes and directions are BC, BA, BF, BH, &c. and from the bisection, m, of the diagonal BD of a parallelogram, whose sides are BC, BA, a right line be drawn to the extremity F, of BF, cutting the diagonal BE, of a parallelogram whose sides are BD, BF, in n, and from n a right line be drawn to the extremity, H, of BH,*
a cutting



cutting the diagonal BG of a parallelogram, whose sides are BE , BH , in p , $Bp : BG :: \text{unity} : \text{number of forces}$.

DEM. The triangles Bpn and GpH , Bnm and nEF , are similar; and consequently $Bp : pG :: Bn : GH(Bn + nE) :: Bm : Bm + EF(Bm + BD \text{ or } 3Bm)$; therefore $Bp : BG :: Bn : 4Bm :: 1 : 4$. The process is similar when there are more forces. Q. E. D.

218. Cor. 1. It is evident, from the construction of the figure, and similar triangles, that $CA : Cm :: 2 : 1$,
 $mF : mn :: 3 : 1$,
 $nH : np :: 4 : 1$; and if the number of forces were equal to any number, m , nH would be to $np :: m : 1$.

219. Cor. 2. The force BG , resulting from the action of any number of forces BC , BA , BF , BH , &c. either in the same, or a different, plane, is to their sum as $BG : BC + BA + BF + BH$, &c. and if the magnitudes and directions of these forces be given, the magnitude and direction of BG may be found; for CB , BA , and the $\angle CBA$ being given, BD , and the $\angle s CBD$, DBA , are known; and DB , BF , and the $\angle DBF$ (which may be found) being given, BE is also given, and in the same manner any other diagonals may be investigated.

220. Cor. 3. If a body therefore be acted upon by any number of forces BA , BC , BD , BE , BF , BG , the directions of the forces resulting from the action of two, three, four, &c. are found by joining AC , bisecting it in m , and taking $mD : mn :: 3 : 1$, $nE : nP :: 4 : 1$, $pF : pq :: 5 : 1$, $qG : qr :: 6 : 1$, and the directions are Bm , Bn , Bp , Bq , Br . Their magnitudes are $2Bm$, $3Bn$, $4Bp$, $5Bq$, $6Br$, &c.

FIG.
XLIV.

FIG.
XLV.

221. Cor. 4. The same force BP may be generated by the action of 1, 2, 3, &c. forces, either in the same, or different, planes, whose directions and quantities are variable in infinitum: for let it result from the combined action of four forces, and taking $BP : Bq :: 4 : 1$, through q draw any line whatever qA , and take $qA : qP :: 3 : 1$; through p draw any line pC , and let pC be to $pn :: 2 : 1$; and through n draw any line nD , and taking nE equal to it; the forces BA , BC , BD , BE , are equivalent to BP . The lines qA , pC , nD , &c. may be drawn in any directions, in any different planes, and of any different lengths.

FIG.
XLVI.

222. PROP. *The directions of three forces x , y , z , acting upon a body B without producing motion, will meet in the center of gravity of a triangle, whose distances from the angular points, are as the magnitudes of the forces.*

DEM. Produce AB , and taking $BE = AB = z$, complete the parallelogram $BDEC$ by drawing from E lines parallel to BC , BD ; BE is equivalent to BC and BD , and equal to z , which is equivalent to x and y , which must therefore be equal to BC , and BD , respectively; but $BG = \frac{BE}{2} = \frac{AB}{2}$, and B is therefore the center of gravity of the triangle. Q. E. D.

Otherwise:

Let BA , BC , BD , be the magnitudes of z , x , y , respectively, and joining DC and bisecting it in G , and joining GA , $AB = 2BG$ (217). Q. E. D.

FIG.
XLVII.

223. PROP. *If a body B , be acted upon by four forces, in different planes, x , y , z , w , or BA , BC , BE , BF , and remain at rest, they are to each other as the three sides and diagonal of a parallelopiped parallel to their directions respectively.*

DEM.

DEM. The forces BA, BC are equivalent to BI , and BI, BE to BD , which must be equal and opposite to w or BF , because B is quiescent. Q. E. D.

224. Cor. An infinite number of parallelopipeds may be described, whose sides and diagonal are the same: for let x, y, z, w , be invariable, and taking BC, BA , equal respectively to y and x , making any angle whatever ABC , and BD, ID in a different plane, respectively equal to w and z , draw BE parallel to ID , and DE to BI . Bisect BI in H , and join EH , intersecting DB in G ; and because the triangles BGH, EGD are similar, $EG : GH :: ED (2BH) : BH :: 2 : 1$; therefore these two forces combined with BE produce a force passing through G (217); and because $DG : GB :: EG : GH :: 2 : 1$ and $DB : GB :: 3 : 1$, the force w is equal to DB . But the angle CBA is infinitely variable, and BD, FD may be drawn in any planes whatever.

FIG.
XLVII.

225. PROP. If four forces, x, y, z, w , which are as BA, BD, BE, BF respectively, act in different planes, upon the body B without producing motion, B is the center of gravity of a triangular pyramid whose base is EAD and vertex F .

FIG.
XLVIII.

DEM. Join AD , and bisect it in m , and, taking mn to $nE :: 1 : 2$, n will be the center of gravity of the triangle EAD , and nB is the direction of the force resulting from x, y, z combined, and its magnitude is equal to $3 \times Bn$, which must be equal and opposite to w or BF , because B is quiescent: BF is therefore equal to $3Bn$, and B consequently the center of gravity of a pyramid, whose base is EAD and vertex F . Q. E. D.

226. Cor. 1. It is evident that $BF = 3Bn$, because the force resulting from $x + y + z$ is to those forces as $3Bn : BA + BD + BE$ (220) and $w : x + y + z :: BF : BA + BD + BE$ (hypoth.) therefore $w = BF = 3Bn$.

M

Cor.

227. Cor. 2. The force resulting from the combined action of BA , BD , BE passes through the center of gravity of a triangle, which is formed by joining A , E , D ; and the direction of a force, equivalent to them, passes through the same point.

FIG.
XLIX.

228. Cor. 3. An indefinite number of pyramids may be described, having the distance of their angular points from the center of gravity always the same; for let BA , BD , BE , BF , be invariable, and making the angle ABD of any magnitude whatever, draw CD , BD , in any plane whatever, equal respectively to BE , BF , and complete the parallelogram $BCDE$, and the process is the same as that in (224).

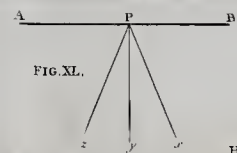


FIG. XL.

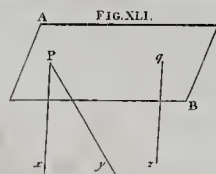


FIG. XLI.

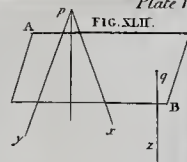


FIG. XLII.

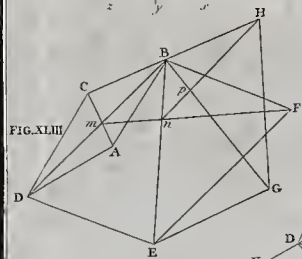


FIG. XLIII.

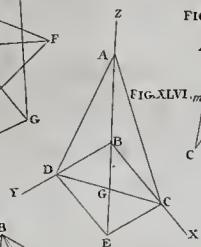


FIG. XLIV.

FIG. XLVI.

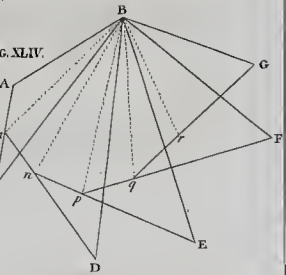


FIG. XLV.

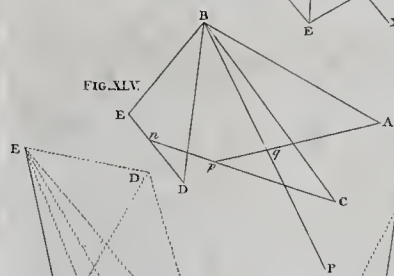


FIG. XLIX.

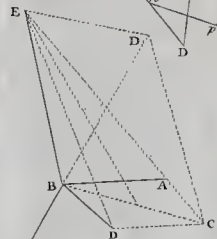


FIG. XLVII.

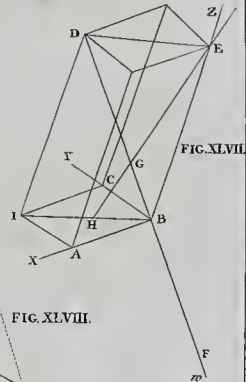
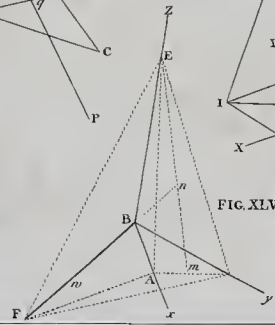


FIG. XLVIII.



C H A P. V.

O F A T T R A C T I O N.

229. DEF. ***T**HAT power or principle in nature, by whose influence bodies, or the constituent parts of bodies, accede, or have a tendency to accede, to each other, without any sensible material impulse, is called attraction.*

230. DEF. *That natural power, by whose influence different bodies, or different parts of the same body, recede, or have a tendency to recede, from each other, without any sensible material effect, is repulsion.*

When the changes of motion in the body *A* are uniformly observed to depend upon the situation and distance of another body *B*, a connexion is understood to obtain between them, expressing some quality or mechanical affection of matter, such as gravity, cohesion, elasticity, magnetism, electricity, which appears to reside in matter, and by whose agency this change of motion is conceived to be produced. If the direction of *A*'s motion, or tendency to motion, be towards *B*, it is said to be attracted by *B*; and if it be from *B*, it is said to be repelled by *B*.

231. That there are in nature motions and tendencies to motion, *conatus accedendi & recedendi*, both in aggregate bodies, and in their minute constituent particles, without any sensible cause, and indeed inexplicable by any known properties of matter, is unquestionably certain. These demonstrate a source of motion distinct from, and repugnant to, any material impulse with which

A T T R A C T I O N O F G R A V I T Y.

we are acquainted, which is denominated attraction, a term indifferently applied to the mechanical affection and its effects, as all inexplicable powers are only measurable by their effects. These effects being found by observation to be different at different distances, the intensity of the power producing them is inferred to be variable, and the law of variation is discovered by an actual mensuration of these effects at different known distances. By this process, these material affections, though intelligible only in their effects, are considered as quantities capable of mathematical comparison, and their ratios are compared with those of lines and numbers. Attraction is usually divided into two kinds: *first*, that which operates at sensible distances, as gravity, electricity, magnetism; and, *secondly*, that, whose effects are limited to almost insensible distances, which is called the attraction of cohesion.

A T T R A C T I O N O F G R A V I T Y.

232. The existence of this species of attraction is evidenced by uniform experience. Matter when supported is heavy, and, the support being removed, accedes towards the earth; when projected obliquely, it deviates from the line of projection, and vibrates when suspended by a string inclined to the earth's surface. And the direction of this motion of matter, and tendency to motion, is invariably the same, towards the center of the earth nearly. A mechanical affection is justly inferred from these phenomena (230), which is conceived to produce this motion and tendency to motion, and is promiscuously called gravity, gravitation, or the attraction of gravity. But more correctly the tendency of a body towards the earth, which is measured by the velocity acquired in a given time, is called the accelerating force of gravity; the weight or moment of a body, which is measured by the product of the quantity of matter and the accelerating force, is called its vis motrix; and the power conceived to be in the earth, which produces this weight and tendency towards it, is called the absolute force or attraction of gravity; and a body influenced by it, is said to be attracted by, or gravitate towards, the earth. All hard and soft
bodies

bodies are subject to this attractive power, and have weight; air, water, and other fluids, and the vapours arising from them, and all odorous substances, and exhalations from all terrestrial bodies, gravitate towards the earth, as they may be collected in a balance and actually weighed; and the weight of these last substances is also rightly inferred from the decrement of weight sustained by the evaporable matter. Even fire and light, and the most volatile vapours, seem to be under the controul of this universal principle, and no matter in the vicinity of the earth, accessible to experiments, hath yet been discovered to be uninfluenced by it, and consequently all matter may be considered as gravitating towards the earth, till future satisfactory experiments induce a different opinion. This attractive power is not confined to bodies in the vicinity of the earth, but extends also to the moon; for the moon describes equal areas in equal times about the earth's center, and is therefore urged by a force directed to that center (NEWTON. Sect. II. P. II.) and coincident in direction with gravity: and it appears, from astronomical observations, that it is equal in quantity to the force of gravity at that distance, and they are consequently the same power. And because the revolutions of their satellites round jupiter and saturn, and of the primary planets round the sun, are phenomena similar to that of the moon round the earth, both being acted upon by forces directed to their respective centers, which vary according to the same law; the satellites therefore gravitate towards their primaries, and these towards the sun (7). The primary planets gravitate towards each other; for jupiter and saturn, in conjunction or at their least distance, are discovered to produce very sensible irregularities in each other's motions; and the motions of their satellites are said also to be subject to irregularities, which are sensible at their least distance where their action is greatest. All the great bodies, composing the solar system, are therefore impressed with this principle of attraction; which, being attached to the whole of any body, must pervade every constituent portion; and the minutest particles of matter gravitate, though perhaps insensibly, towards each other. The operation of this principle between two small bodies, at the surface of the earth, is only rendered insensible from the predominant influence of the earth;

ATTRACTION OF GRAVITY.

earth; for a pendulous body was observed to be considerably deflected from its vertical situation by the attraction of the mountain Chimborazo in Peru; and, by some ingenious experiments made with great precision, Dr. Maskelyne ascertained the quantity of attraction of the mountain Schehallien in Scotland.* Every mountain therefore, and every less portion of the earth, possesses this quality; which might have been inferred from (215). For if a particle of matter be attracted by every part of a sphere of matter, equally at equal distances, the direction of the combined attraction will pass through its center and $v.v$; but it appears, from experience, that all bodies descend in directions perpendicular to the earth's surface, and, because the earth is nearly spherical, the whole force, producing this descent, is directed nearly to the center, and is consequently combined of the force of every particle.

233. PROP. *The accelerating force of gravity at equal distances from the earth's center, is the same in all bodies, whether quiescent or moving, and whatever be their magnitude, figure, density.*

DEM. This proposition is only demonstrable from experiments. Two equal wooden boxes, suspended by threads of eleven feet in length, one of which was filled with wood, and an equal weight of gold fixed in the center of oscillation of the other, were discovered by Sir I. Newton, to perform all equal vibrations in the same time; and numberless experiments demonstrate that all bodies,

* If a mountain do possess an attractive power, a body suspended near it will be deflected from its vertical position, and point to a false zenith, and the arc, measuring the distance of this from the true, is the measure of the mountain's attraction. In an observation on the north side of Schehallien, the plumb line, being deflected towards the mountain, pointed too much to the south, and gave the distance of a star from the zenith too much to the north; and from an observation made on the north side, the distance was too much to the south. The difference of the latitude of the two stations, collected by these observations, must be greater than it really is, by the sum of the arcs measuring the deflections from the two zeniths. From observations of ten stars near the zenith, Dr. Maskelyne found the difference of the latitude of the two stations to be $54''6$; and, from a mensuration of their distance, it was only $43''$: the difference of these is $11''6$, and the half of it $5''8$, measures the attraction of the mountain. This experiment is most convincing, and decisive in support of the universality of gravitation.

dies, however different their magnitude, figure, and density, descend in an exhausted receiver, whatever be its height, exactly in the same time; and consequently the tendency of bodies towards the earth, or their accelerating force, is the same whether quiescent or moving, &c.* Q. E. D.

234. PROP. *The weight of bodies, at the surface of the earth, are proportional to their quantities of matter.*

DEM. The weight of a body is evidently as the number of equal particles or quantity of matter contained in it, multiplied into the tendency towards the earth, or accelerating force of each, and this being given (233) varies as the quantity of matter. Q. E. D.

Otherwise :

The weight of any two bodies A and B , or the forces producing them, are clearly proportional to their vis inertiae; for if the vis inertiae of A were to that of B , as 1, 2, 3, &c. : 1, it is evident that the weights or forces, producing the same velocities in A and B in the same time, will be as 1, 2, 3, &c. : 1; and consequently the weight, being as the vis inertiae, will be as the quantity of matter (166). Q. E. D.

235. Cor. If the accelerating force of gravity were encreased in any ratio, the weight of a given body would be encreased in the same ratio. Substituting therefore W , \mathcal{Q} , F , for the weight, quantity of matter, and force of gravity, respectively, and supposing them to be variable; W will be as $\mathcal{Q} \times F$; F as $\frac{W}{\mathcal{Q}}$; \mathcal{Q} as $\frac{W}{F}$, and, if W be given, F and \mathcal{Q} are inversely as each other.

S C H O-

* All bodies, whether great or small, dense or rare, acquire a velocity in falling 1'', which would carry them uniformly through 32.2 feet in 1'', and an encrease of velocity, equal to this, is found to be added in every succeeding second of time. e

S C H O L I U M.*

236. In the same place, and at the same distance from the æquator, bodies are found to descend through the same space in the same time always; but, the force of gravity being diminished unequally by the diurnal rotation of the earth, they do not, and ought not to, descend through equal spaces, or perform equal vibrations, in the same time, at different distances from the æquator. However, this unequal diminution of gravity, arising from the different centrifugal forces, does not afford a solution for the different spaces fallen through; and there must be some other causes of inequality, which are probably the difference of distance from the center, and different species of matter near the places of observation.

237. PROP. *The force of gravity varies as the square of the distance from the center of the earth inversely.*

PLATE
VI.
FIG.
LI.

DEM. Gravity acts in right lines, and, whatever be the cause, is equally diffused over every point of the surface $ABDC$, equidistant from the center of the earth S ; and that same influence is also equally diffused over the surface $abcd$, similar to the former and similarly situated. Taking $almn$ equal and similar to $ABDC$, the influence of gravity upon $almn$: influence upon $abcd$ or $ABDC$:: $almn (ABDC) : abdc :: SA^2 : Sa^2$. Q. E. D.

Another

PLATE
VI.
FIG.
L.

* PROP. *Supposing the earth to be spherical, the diminution of gravity, arising from the centrifugal force, varies as the square of the cosine of latitude.*

DEM. Let Pp be the earth's axis, and $ÆQ$ its æquator, and the centrifugal force at Q : centrifugal force at A :: $QC : AD$; the centrifugal force at A : that part opposite to gravity :: $AN : AL$ (supposing NL to be perpendicular to AC) :: $AC : AD$; and consequently the centrifugal force, or diminution of gravity at Q : the diminution of gravity at A :: $QC^2 : AD^2$. This diminution of gravity varies therefore as AD^2 , or as the square of the cosine of the place's distance from Q . Q. E. D.

Another demonstration :

The magnitude of the earth's attractive force, at different distances, varies as the space through which it impels a body in equal times; but a body at the earth's surface falls through 16.1 feet in 1", and the moon, at her mean distance, nearly 60 semidiameters of the earth, falls through $\frac{16.1}{60 \times 60}$ feet in 1" (from observation): presuming therefore that the moon is retained in her orbit by the earth's attractive power, the force of gravity at the earth's surface, is to the force at the moon, as $16.1 : \frac{16.1}{60 \times 60} :: 60 \times 60 : 1$.

Q. E. D.

238. Cor. 1. The first of these demonstrations is applicable to all forces diffused by rectilineal effluvia of matter from a center, as heat, smells, &c.

239. Cor. 2. Putting d for the semidiameter of the earth, and W for the weight of a body at the earth's surface, its weight at any other distance nd (n being a number) will be equal to $\frac{W}{n^2}$. A weight, of 3600 lb. at the earth's surface, is equal to $\frac{3600}{60 \times 60}$ or 1 lb. at the distance of the moon.

240. Cor. 3. If BC and bc be homogeneous and similar, their weights at unequal distances are equal; for they are to each other as $\frac{BC}{SA^2} : \frac{bc}{Sa^2}$ (237) as 1 : 1. If they be not homogeneous, their weights will be as their quantities of matter, that is, as their weights at equal distances, divided by the squares of their distances.

S C H O L I U M I.

241. The variation of the force of gravity was collected from observation in Peru by Mess. Condamine and Bouguer. Condamine found that the same pendulum upon mount Pichinca, at Quito, and upon the banks of the river of the Amazons, performed respectively 98720, 98740 and 98770 vibrations, in twenty-four hours. And Bouguer observed upon the summit of Pichinca, and upon the sea-shore, being $\frac{1}{1349}$ th part nearer to the center of the earth, that the lengths of isocronous pendulums were to each other as 438.71 to 439.7. These observations demonstrate a diminution of the force of gravity in receding from the earth's surface; but the law of variation cannot be collected from them, or any other similar experiments, with precision; because the different states of the atmosphere, of heat and cold, and different densities of the earth contiguous to the different places of observation, would proportionably affect the experiments, and these cannot be ascertained with sufficient accuracy. From the observations of Bouguer, the force of gravity upon the mountain and upon the sea-shore, is as 99575 : 99802; and, supposing it to vary inversely as the square of the distance, it is as 99575 : 99722; which differs less from the true law than could be expected.

S C H O L I U M II.

242. An investigation of the cause of attractions and their mode of operation, and particularly of this mechanical affection of matter, its gravity, so universally prevalent, hath long been the object of the philosopher's researches, but without effect; the subject is still involved in impenetrable darkness, and his wishes are unsatisfied. Dr. Halley refers it to the immediate agency of the Creator; Mr. Cotes deems it essential to matter, like extension and mobility, &c. and Sir I. Newton seems to entertain a different opinion, and attributes it to some undiscovered and invisible mechanical affection of matter. Electrical and magnetic experiments prove the existence

ence of a subtle fluid pervading the pores of the densest bodies, and the gravity of bodies may arise from the action of a fluid, whose particles are so small as to pass through all bodies, and does not therefore resist their motion; and, if endued with a strong repulsive power, it will by expanding itself press upon gross bodies. This æther, Newton supposes, is, from its repulsive force, much rarer within the dense bodies of the sun and planets, than in empty space, and in receding from them becomes perpetually denser and more elastic, and therefore occasions their gravity, every body being impelled from the denser towards the rarer parts of the fluid. This however is conjectural, and its truth or falsehood can only be established by future experiments. All that is certainly known about this mysterious cause, is, that it cannot arise from the pressure of a fluid similar to those with which we are acquainted, whether soft or elastic, quiet or agitated, for the following reasons. *First*, Because a fluid, subtle enough to penetrate the inmost recesses of the densest bodies, and so rare as not to impede their motion, cannot be conceived capable of communicating motion to them. *Secondly*, Because the attraction of gravity penetrates the inmost recesses of bodies, and their weights are proportional to their quantities of matter, not magnitude of surface; but the pressure of every fluid with which we are acquainted, varies as the surface opposed to it. *Thirdly*, Because the attraction of gravity acts with equal intensity upon bodies, whether quiescent or moving; but all known fluids act upon quiescent and moving bodies with different forces. But, however unsearchable the efficient cause of gravity may be, its final cause is most conspicuous, being the preservation of the earth and other planets, and of their periodical revolutions, which are only continued by an uninterrupted exertion of this mechanical affection of matter.

C H A P. VI.

ATTRACTION OF COHESION.

* **A** SECOND species of attraction, distinguished from that of gravity by the intensity of its action, and the small distances to which its influence extends, is that observable between the minute particles of bodies, whether hard, soft, or fluid. Hard and soft bodies have an appearance of endless composition; but they are composed of insensibly minute particles which cohere together and form greater, and these, being successively compounded by the influence of the same principle, form successively greater, till, by repeated coalitions, they become objects of sense; and it appears, from many phenomena, that the particles of all known fluids have also a tendency to cohere with others. This mutual tendency to accede to each other, of the component parts of hard, soft, and fluid bodies, not arising from any sensible material impulse, indicates that philosophical relation denominated attraction (230); which, being great at the point of contact, and evanescent at the least sensible distance, coincides, in these marks of similitude, with innumerable phenomena, and the material power, from whose agency they are conceived to result, is called by the same name, the attraction of cohesion, because they are all effected by the mutual action of particles, whose force is great when contiguous, and, when removed to any sensible distance, evanescent. But whether the phenomena, which have these marks of coincidence and similitude, result from a common cause, varying at different distances, like the attraction of gravity, according to the same law of the distance, can only be ascertained by experiments, which are yet too undecisive to allow of accurate discrimination. The laws and limits of the action of these minute particles are still unknown,

• Helsham, Lect. I, II. Muschenbroek, Ch. XVIII. Newt. Opt. p. 380. Hamilton's Tracts, Lect. II.

known, and, from their almost insensibly minute sphere of action, perhaps undiscoverable; but its intensity seems to vary according to some higher inverse law of the distance than the duplicate, and therefore soon evanescent, that it may not interfere with the attraction of gravity.

243. PROP. *If the particles composing two spherical surfaces AFC, BDE, attract with forces, equal at equal distances, and at unequal distances, varying as the n^{th} power of the distance inversely, a particle P will be attracted by them with a force varying inversely as that power of the distance, whose exponent is $n - 2$.* FIG. LII.

DEM. The attraction of P towards these surfaces, is as the number of particles and power of each, or as $\frac{AFC}{PA^n} : \frac{BDE}{PB^n}$, or as $\frac{PA^2}{PA^n} : \frac{PB^2}{PB^n}$, or as $\frac{1}{PA^{n-2}} : \frac{1}{PB^{n-2}}$. Q. E. D.

244. Cor. 1. If $n = 2$, the attractions of these laminæ are as $1 : 1$, or equal to each other; and the attraction of the solid $BECA$ is equal to that of one lamina multiplied into their number. The attraction therefore of this solid is to that of $PAC :: BA : PA$; and consequently the attraction of a portion of matter, in contact with P , is not much greater than when removed to a small distance.

245. Cor. 2. If $n = 3$, the attractive powers of any laminæ BDE, AFC , are to each other as $\frac{1}{PB} : \frac{1}{PA}$, or as the ordinates BI to AL of an equilateral hyperbola; and the whole attractive power of the body $BECA$ is to that of $PAC :: \text{area } BL : \text{area } AR :: \text{a finite quantity} : \text{a quantity infinitely greater}$. According to this supposition, the attraction of P is not much affected by the addition, or deduction, of new matter at a small distance; for the encrease,

encrease, or decrease, of force resulting from it, would be finite, and consequently infinitely less than that resulting from contact. But if the surface of *P* were enlarged, and in contact with new matter, attracting according to this law, its cohesive force would be proportionably encreased.

246. Cor. 3. If this attraction, at a finite distance, had a finite ratio to that of gravity, at the point of contact it would be infinitely superior to that of gravity; and, if at the point of contact, the ratio between this attraction and gravity were finite, at any assignable distance it would be evanescent.

247. Cor. 4. The attraction of cohesion, being great in contact and insensible at a small distance, varies in a higher than the inverse duplicate ratio of the distance.

FIG.
LII.
LIV.

248. Cor. 5. Since the cohesive force of any particle *P*, in contact with the body *PAC*, varies as the quantity of surface, it will be greatest when the surfaces of contact are plane. Two species of matter constructed of dissimilar particles, *a, b, c, d, e, f, g, &c.* *l, m, n, o, p, q, &c.* have different degrees of hardness; for the surfaces of contact, and cohesive forces of *a, b, c, d, &c.* are greater than those of *l, m, n, &c.*

249. Cor. 6. Of unequal similar particles, the smallest are capable of the firmest cohesion, because they have the greatest surfaces of contact compared with their magnitudes, the surfaces being diminished only as the squares, and the magnitudes as the cubes, of any lines similarly placed in them.

250. Cor. 7. Hard bodies therefore may be conceived to be constructed of very small particles, or such as are terminated by plane surfaces; for collections of these particles cemented by this principle

principle of cohesion, with which they are proved to be impressed, would constitute hard bodies. And soft bodies may be conceived to be combined from large spherical particles, or such as have many angular points, and do not allow of great surfaces of contact.

SCHOLIUM.

251. The existence of this natural power is collected, and a general idea of it formed, from the following familiar observations and experiments. *First, An attraction subsists between the component parts of hard and soft bodies.* The force required to separate the contiguous parts of hard bodies, which is equal to the force of cohesion, is much superior to their gravity, though they be full of pores and consist of portions collected into one mass, and touching only a few points; and it must be much greater between the smallest particles, whose surfaces are in contact, without any great intermediate vacuities. This attraction is very strong in all tenacious and viscid bodies, as pitch, rosin, &c. which adhere to every thing in contact with them; and if a body very soft and viscid be extended, it will again contract its dimensions by the tendency of the parts to each other. Large ships, floating near each other in a calm sea, have a great tendency to each other, and are with difficulty kept from coming together. If the surfaces of the segments of two leaden bullets, whose diameters are not greater than $\frac{1}{4}$ th of an inch, be polished, and compressed together by a gentle twist, it is said that a weight of 100 lb. is frequently required to separate them. And if two polished plates of glass, marble, brass, or any metalline substances, be compressed together and suspended in an exhausted receiver, the weight of the lower plate will not dissolve their cohesion; and if any fine thread be wrapped round one of them several times, at considerable intervals, after compression, their cohesion is still sensible, though they be removed from actual contact by the thickness of the thread. The cohesive forces of these substances is augmented by moistening their surfaces with water, oil, grease, &c. which expels the air from their pores, or, if very tenacious, keeps it confined, and prevents the
action

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action of its repulsive power. The cohesive force of two polished plane surfaces of metal, whose diameters were two inches, heated in boiling water, and besmeared with grease, oil, &c. was overcome by the following weights :

	Cold grease.	Hot grease.
Planes of glass by	130 lb.	by 300 lb.
brass	150	800
copper	200	850
marble	225	600
silver	150	250
iron	300	950.

When these surfaces were moistened with water, oil, &c. and compressed together, their cohesion was overcome by the weights in the following table :

With water	by 12 oz.
oil	18
Venice turpentine	24
rosin	850 lb.
tallow-candle	800
pitch	1400.

Secondly, An attraction subsists between the particles of any portions of water, oil, mercury, and all other fluids, except air, fire, and light. Small portions of these fluids form themselves into globular drops both in the open air and in vacuo; and if a drop of mercury be poured upon clean paper, or a drop of water upon the leaf of a plant, their spherical figure is not changed, the gravity of their parts being unable to dissolve their cohesion; and two drops of any fluid, not much attracted by the surface on which they are placed, will coalesce into one when in contact, as is often observable in drops of water lodged upon the leaves of plants. If mercury, well purged of air, be poured into a clean glass tube 70 inches long, so that its parts be contiguous to each other and to the glass, after inversion the whole column will remain suspended; but the pressure of the atmosphere sustains only 29 or 30 inches, and the remainder must be supported by some other agent, which is chiefly the mutual adhesion of the parts; for if they be discontinued

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105

tinued by a bubble of air intervening, or by shaking the glass, the column subsides to the height of 29 or 30 inches.

Thirdly, This species of attraction subsists between the particles of different fluids, and between them and other substances. If a piece of fir wood, whose surface is one square inch, be soaked in water and float upon its surface, a weight of 50 grains, besides an equivalent to its own weight, is required to detach it from the fluid; and when its surface is enlarged, a proportionably greater weight is required. Water rises near the sides of the vessel containing it, round a glass bubble floating upon its surface, up capillary tubes of glass, plants, &c. A remarkable instance of this attraction occurs in the new invented water-pump. If two wheels, or pulleys, *B* and *D* in the same plane, be made to revolve with great velocity, and *D* be immersed in water, a column of water will ascend with the rope; and, if there be two or more grooves in *B* and *D*, and as many ropes pass over them, at the distance of about an inch, the columns raised by each rope will cohere, and the quantity of water be much augmented.* Air is incorporated with most hard bodies, possesses their interstices, and probably serves as a bond of union to their constituent parts. It appears from many experiments, that the quantity of air, detached from some hard bodies in which it was consolidated, by the action of fire, or some particular fermentation, is very considerable; and Dr. Hale discovered that $\frac{2}{3}$ of a calculus humanus were air. It is attracted by water, and perhaps all fluids, into whose pores it insinuates itself, is intimately mixed

FIG.
LV.

* There are two of these machines, at Windsor, of different dimensions. The depth of the well at the round tower is 178 feet: *D* is the wheel in the water, its diameter = 12 inches, the thickness of the rope $DB = \frac{1}{2}$ inch nearly; diameter of the upper wheel *B* = 13 inches; diameter of *C* = 11 inches; diameter of *E* = 4 feet 6 inches; a power, applied at *F* turns the wheel *E* round, and that, by means of the string *EC*, communicates motion to *G*, which has the same axis with *B*.

FIG.
LV.

In the other machine the depth of the well is 95 feet, and, in this, the quantity of water, raised by the utmost efforts of a man, was at the rate of nine gallons in a minute. At Paris 500 pounds of water were elevated through an altitude equal to 240 feet, in ten minutes, when the diameter of the rope, surrounding the pulleys, did not exceed six French lines. In the beginning of the motion, the column, adhering to the rope, is always less than when it has been worked for some time, and continues to encrease, till the surrounding air partake of its motion.

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mixed with them, and coheres too strongly to be separated without the agency of fire, or some more attractive substance. If sal ammoniac, or corrosive mercury sublimite, be dissolved in water, bubbles of air will disengage themselves from the water, and, adhering very tenaciously to the thin particles of the salt, ascend with them, and burst upon the surface. When water, beer, &c. are poured into a glass, or other vessel, many bubbles of air are detached from the fluid by the stronger attraction of the vessel, and adhere to its sides and bottom; and, if these be rough, or have more points of contact, the number of bubbles will be increased. Light is attracted by glass, water, and all transparent mediums, as appears by its refraction and reflection; and also by all other substances, as is observable in an experiment of Sir I. Newton.* A beam of light admitted through an aperture into a dark chamber, and passing by the edge of any substance, will be deflected from a rectilinear course, and by repeated attractions and repulsions describe an undulating line in its passage.

OF HARDNESS, SOFTNESS, AND ELASTICITY.

252. DEF. *A hard body is that whose parts are not easily moved from their places; as wood, metal, stones, &c.*

Bodies constituted in such a manner that their component parts do not give way to compression, and whose cohesive force is insuperable, may be called perfectly hard. The hardest bodies with which we are acquainted, as adamant, flint, gems, tempered steel, &c. are full of vacuities, and contiguous portions of them, not being in contact all round, are penetrable, and separable by the action of a sufficient force. Fire insinuates itself amongst the vacuities of the densest and hardest bodies, and produces fusion or expansion, which cannot be effected without an entire separation, and change of place, of their component parts. Perfect hardness therefore seems to be confined to the elementary particles, or least parts into which bodies can be divided, whose figure and dimensions cannot be separated by any known power, and are impenetrable.

253. DEF.

* Newton's Optics, p. 317.

253. DEF. *A soft body is that whose parts change their position by the action of a small force, and retain it when the force is removed; as butter, snow, wax, &c.*

Bodies constructed in such a manner that their vacuities are not replete with any fluid, and whose parts do not repel each other, and are only retained in their places by their inertia, may be called perfectly soft; but no species of matter is perfectly soft, for the particles of all bodies, not possessed of a repelling power, require a much greater force to separate them, than what is equivalent to their inertia. The degrees of hardness and softness are infinitely variable, and the limit where hardness ends, and softness begins, cannot be defined. In the congress of hard and soft bodies, there is no cause of separation after they come into contact, and they will therefore either move on together, or be quiescent, after impact; and, the effects of collision being ascertained when there is no repulsive power, allowance must afterwards be made for their different degrees of repulsion.

254. DEF. *An elastic body is that which changes its figure, or the position of its parts, by the action of a force, and recovers, or has a tendency to recover, its figure.*

Elasticity is said to be perfect, when the parts of the body return to their first situation with a force equal to the force displacing them, and imperfect, when they do not.

255. PROP. *If a perfectly hard body A impinge upon a perfectly elastic immovable body B, it will be reflected with a velocity equal to that of impact.*

DEM. The particles, composing the surface of B, are continually removed from their places by the compression of A, till it and their resistance become equal; and then, returning to their first situation with a force equal to that of compression, and

in a direction opposite to A 's motion, A must evidently be repelled with a velocity equal to that of impact. Q. E. D.

256. Cor. 1. If A and B be both perfectly elastic, the particles, composing their surfaces, recede equally, though the recess is less than when A was perfectly hard; but the effect is the same, the whole force of restitution being equal to that of compression, and conspiring to repel A .

257. Cor. 2. If B be moveable, the velocity lost by A is double of that lost by impact only; for the parts of B 's surface, contiguous to the point of impact, restoring themselves with a force equal to that which displaced them and acting against A , it will be equally retarded by the forces of impact and restitution; and, for the same reason, the velocity communicated to B is double of what it would be, were both bodies perfectly hard.

258. Cor. 3. The relative velocity of A and B is the same before and after impact, or the velocity with which they accede to each other before impact, is equal to the velocity with which they recede from each other after impact; for the forces of impact and restitution, being equal and opposite, produce the same effects in opposite directions, and the relative velocity, which is destroyed by impact, must be restored by the force of elasticity.

S C H O L I U M.

259. The whole time of contact of A and B , may be divided into two equal periods, the former between their first contact and the cessation of A 's compression, and the latter between this point and their last contact, when they are separated by the restoration of their relative velocity. And, if the time of A 's compression be divided into very small equal parts, A 's decrements of velocity produced in them will perpetually encrease, from the first contact

contact to the end of compression, where its force of protrusion vanishes, and it is stationary for a moment, if B be fixed, or, if B be moveable, its velocity is equal, for a moment, to that of B . The particles then returning, by their elasticity, to their first situation, will communicate, in equal times, perpetually decreasing increments of velocity to A , equal to the corresponding decrements during its compression, which vanish when the surface has regained its natural state, when A leaves B , and is reflected with a velocity equal to that lost by impact. When A and B are both perfectly elastic, the particles, contiguous to the point of impact, give way equally; and when A is perfectly hard, and B imperfectly elastic, or when they are both elastic in different degrees, the particles yield to impact unequally, and a pit or cavity must be formed in one of them.

260. Cor. 4. If the times of compression and restitution be divided into the same number of equal moments, and the decrements and increments of velocity, in corresponding moments of compression and dilatation, be as $n : 1$, the velocity lost by impact is to that communicated in an opposite direction, in the same ratio of $n : 1$ (EUC. B.V. P.XII.). And, if this ratio between these velocities always obtain, their corresponding increments and decrements will always be to each other in that ratio.

S C H O L I U M.

261. *The existence of elasticity is demonstrated by numberless experiments. When two bodies impinge, they must coalesce if the surfaces of contact be immoveable, or, after receding, remain immoveable; and their separation is a proof of elasticity. Metals, semimetals, stones, gems, fossils, cartilages, most fluids, as air and even water, exert an influence opposite to the direction of the force compressing them, and discover a tendency to return to their natural state, which is, in all of them, imperfect and less than the force impressed; but most perfect in glass, ivory, hardened steel and cartilages.

* Muschenbroek, Ch. XVI. Defaguliers, Lect. VI.

lages. Elasticity is encreased by encreasing the density of a body; for metals are rendered more elastic by being beaten with a hammer, and their elasticity, which was not perceptible before, becomes after this very sensible. Steel is more elastic when tempered, and its density is then encreased in the ratio of 7809 : 7738. It is also sometimes encreased by cold, as the range of a cannon ball is greater when the cannon is cold, than when heated, and the string of a violin, or a steel lamina, is inflected, and recovers its situation with less force in hot than in cold weather. The sphere of action of the component parts of an elastic body, and the moments of time in which they lose, and recover, their situation, are too small to allow of decisive experiments for ascertaining their intensity, and laws of operation, at different distances from their natural state; and the result of experiments is only the relative magnitude of the velocity lost by impact, or the whole effect of compression, and that communicated by the whole aggregate force of restitution. Metal fibres, and thin steel laminæ, exhibit no elasticity unless stretched to a certain degree, and inflected by a certain force, as appears from lax chords, which, if a little stretched and removed from their natural state, discover no tendency to return to it; and, when the inflection of a fibre is very great, the influence of elasticity seems to be annihilated, as appears from fibres of wood, which, if inflected beyond a certain limit, remain quiescent and have no tendency to recover their situation. This is also observable in elastic bodies, for their elasticity is only discovered by impact, and the force of impact may be so small as to excite no sensible motion of the constituent parts, or so great as to destroy their elasticity; but the limits where it begins and terminates are unknown. A general idea of elasticity may be formed by considering the most simple cases of the vibrations of fibres, or thin steel laminæ, and conceiving elastic bodies to be composed of them.

FIG.
LVI.

EXP. I. If any fibre, metalline chord, or thin lamina of steel, whose length is AB , be stretched and fixed to two immoveable points A and B , and inflected into the position ACB by a power which ceases to act at C , it will return by its elasticity to its natural state AB ,

AB, and, proceeding with the velocity acquired, continue to perform nearly equal vibrations on each side of *AB* till its motion be destroyed by friction, and the resistance of the air.

EXP. II. If the distance of this fibre from a table, to which it is parallel, be equal to *D*, and a spherical ball *P*, whose diameter is $2D$, be rolled against it, the fibre will be inflected, and this inflection will encrease till *P* be quiescent, and, then returning to its first situation, *P* will be repelled and detached from it, when arrived at *AB*, with a velocity nearly equal to that of impact.

EXP. III. Very small inflections *PD*, *PC*, of the same fibre *AB*, are found to be nearly as the inflecting forces; but when the inflections are considerable, they vary as some power of the force, whose exponent is less than unity or fractional. If *AB* be a small brass wire tended by a weight *W* of 3 lb. and inflected at *P* by weights equal to $\frac{1}{2}$ oz. and 1 oz. successively, $PD : PC :: \frac{2}{5} : \frac{4}{5} :: \frac{1}{2} : 1$; or the inflections are as the inflecting forces.

EXP. IV. If the lengths of the wire be multiplied by 2, 3, &c. and the small inflections *PD*, *PC* be always to each other as $\frac{1}{2} : 1$; the inflecting forces are found to be as 1 : 2 and 1 : 3, &c. or inversely as the lengths of the wire.

EXP. V. If the lengths of the wire be the same as in Exp. III, *W* be equal to 3 lb. and 6 lb. and $PD : PC :: 1 : 2$, the same as before; the inflecting forces are as 1 : 2, or they vary directly as *W*.

262. Cor. 1. If *F* represent the inflecting force, *L* the length of the wire, *I* a small inflection *PD* or *PC*; *I* is as *F* (Exp. III.); *F* is as $\frac{1}{L}$ (Exp. IV.); and *F* is as *W* (Exp. v.). And consequently

if

if I, F, W be supposed to vary, the inflecting force F will vary as $\frac{I \times W}{L}$.

263. Cor. 2. Fibres of unequal thicknesses may be conceived to be composed of a greater or less number of finer fibres of the same thickness, and if W, L, I , be given, it is evident that F will be as the number of smaller fibres, or as the area of a section of the fibre composed of them, or as the square of its diameter (D^2); and consequently F will vary as $\frac{I \times W \times D^2}{L}$.

264. Cor. 3. It is collected from these and other similar experiments, that the elasticity of a stretched fibre appears as soon as it is inflected by the impact of the ball, and continues to encrease to the limit of inflection, where, the moment of the ball and resistance of the fibre becoming equal and opposite, the protrusion ceases; and, because the velocities of impact and resiliion are always nearly equal whatever be the velocity of impact, the velocities of P are equal, at equal distances from the limit of inflection, both in its progress and regress.

FIG.
LVII.

265. Cor. 4. Elastic bodies may be conceived to be formed of elastic fibres or strata, such as AB ; for let the sphere DBE be imagined to be composed of such strata, and stricken at D by a body perfectly hard, and the parts nearest to D , receding by the force of impact, will communicate motion to the contiguous parts, and these to the next, till the different strata be inflected as is represented by the dotted lines in the figure. When the motion of the impinging body is extinguished, and the particles, composing the several strata, are no longer protruded, they will return, by their elasticity, to their first situation; and, if their velocity be perfect, with the same force in acceding to, and receding from, D ; and consequently the impinging body will be reflected with a force equal

equal to the force of impact. And that this is not merely hypothetical, but that motion is diffused from the point of impact to the remote parts of elastic bodies, is presumed from the following experiments.

EXP. VI. If a spherical ball of ivory *A* be pressed against another *B*, whose surface is fresh painted with any colour, it will receive a small point of that colour upon its surface; but if *A* impinge upon *B* with any velocity, the breadth of the spot will be magnified, and become still greater as the velocity of impact is increased. And, because the ball retains its spherical figure after impact, the parts of its surface must have lost, and recovered, their first situation.

EXP. VII. If two glass balls impinge with a proper degree of velocity, the interior parts of the balls will be broken, though the exterior, contiguous to the point of impact, be unbroken.

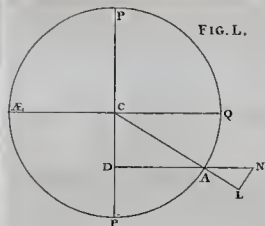
EXP. VIII. If two ivory balls *A* and *B* be suspended from the same point by two strings of the same length, and the less ball *A* impinge upon *B* at rest with a given velocity, *A* will be reflected always to the same height, and *B* will be impelled to the same height, upon the graduated periphery of a circle whose radius is the length of the string. But if either *A* or *B* be hollowed and lead inserted in the center, or nearer to the posterior surface, neither ball, though their weight be the same, will ascend as high as before the insertion of lead.

266. Cor. 1. Motion is therefore communicated from the point of impact to the contiguous parts (EXP. VI.), and diffused from thence to the remote parts of every elastic body (EXP. VII. & VIII.)

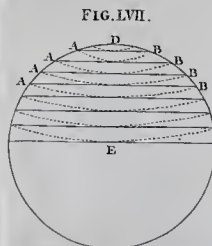
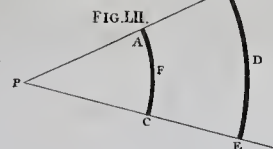
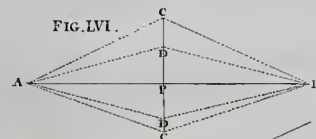
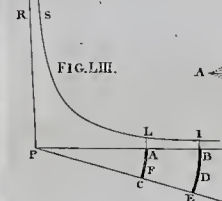
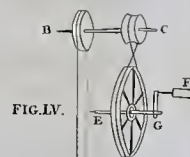
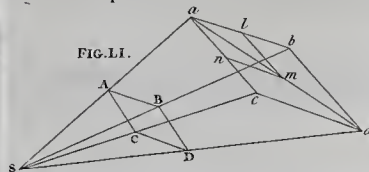
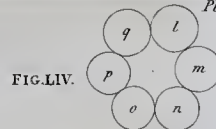
267. Cor. 2. It is evident, from EXP. VIII., that the progressive motion of the parts, from the point of impact, is stopped by the insertion of lead, and consequently that the force of restitution, and change of figure, is less, than before it was inserted.

EXP. IX. A stroke or friction upon the edge of a glass, filled with water, communicates a tremulous motion to the parts of the glass, which is visibly communicated to the water. A reed or stick, placed across the bottom of a large glass bell, will fall when the bell is struck, the stroke producing a change of figure; and, if a piece of metal be fixed near the brim or lip of the bell without touching it, and the bell be stricken by any hard body, it is seen to touch the metal, and a successive of sounds, perpetually decaying, may be heard. If the edge of the bell be pinched, and the fingers suddenly withdrawn, the same sound is heard without producing any sensible motion towards the metal, or displacing the reed across it.

268. Cor. The motion diffused from the point of impact, to the remote parts of an elastic body, is continued for some time, and diminished gradually till it vanishes. And there seem to be two kinds of vibrations of the parts of an elastic body, one of which is quick, and called a tremor of its minute parts, and the other slower and longer, by which its figure is changed, and an impinging body repelled.



a	b	c	d
e	f	g	h





C H A P. VII.

MECHANICAL POWERS.

* **T**HE existence, and intensity of operation, of the mechanical affections, gravity, cohesion and elasticity, and the nature of the other qualities of matter, being ascertained experimentally, they are assumed as established principles, and their efficacy in the production of pressure, motion, and other phenomena, is the next object of mechanical philosophy. There are six simple machines, commonly stiled mechanical powers, from the effects produced, with their intervention, by the action of gravity and animal exertions, viz. the lever, wheel and axis, pulley, wedge, inclined plane and screw. They are all calculated to communicate motion to bodies, and sustain their pressure, for which, the power unassisted by them, is incompetent; and the artifice in all consists in distributing the weight amongst such a number of agents, that the part sustained by the power may bear a small ratio to the whole. Thus, a power incapable of communicating motion to, or supporting the pressure of, a body, without mechanical assistance, may effect its purpose by transferring a part of the weight upon a fulcrum, distributing it amongst a number of pulleys, or placing it upon an inclined plane or screw; and, by this artifice, a power P may keep a weight suspended which exceeds it in any assigned ratio, though without any acquisition of moment in a given direction; for motion is only communicable according to the established natural relations subsisting between matter and motion, and the magnitudes of two powers, in equilibrio, are always inversely as their velocities.

* Keil's Physics, Lect. X. Helsham, Lect. VI. Emerson, Prop. 18, &c. Graves, L. I. C. X. Muschenbroek, Ch. VIII. Varignon, pag. 305. Maclaurin's Newton, B. II. Ch. III. Hamilton's Essay on the Principles of Mechanics. Morgan's Notes to Rohault.

L E V E R.

269. DEF. *A lever is a bar of wood or metal, and is usually represented by an inflexible line, without gravity, revolving about a fixed point, called the fulcrum, by the action of a power upon its arms.*

PLATE
VII.
FIG.
LVIII.

The points W , P , where the weight and power act, are the points of suspension, and the immoveable point F , about which every point of the lever revolves, is called promiscuously the fulcrum, hypo-mochlion, and center of motion. There are three kinds of levers: 1. the fulcrum is between the power and weight, as in the common balance, scissars, snuffers, &c. 2. The weight is between the power and centre of motion, as the oars and rudder of a boat, cutting knives fixed at one end, doors, &c. 3. The power is between the weight and fulcrum, as a ladder raised against a wall, a weight raised by the arm, where the center of motion is at the shoulder, &c.

FIG.
LIX.

270. Cor. 1. Forces whose magnitudes are to each other as PA , PB , PC , &c. acting at P and terminated by a line AC , parallel to PF , have the same effect: for each may be resolved into two forces, one perpendicular, and the other parallel, to PF ; of which, the perpendicular parts are equal and entirely employed in producing a rotation of the lever round F , or in supporting a body W placed on the other side of F , and acting perpendicularly to FW , and the parallel parts only produce a motion in the direction PF , and do not produce any rotation, or contribute to the support of W .

FIG.
LX.

271. Cor. 2. If a given power, represented by PB , act at the same point P in any direction PD , its efficacy to turn the lever round F , or support any body W , is as the chord PC of the circle whose diameter is PB ; for PD being taken equal to PB and resolved

resolved into two forces, one DE perpendicular, and the other PE parallel, to PF , it appears from similar triangles, that DE , the only effective part of the force, is equal to PC .

272. PROP. Two powers W, P , acting upon a lever PW , whose center of motion is F , at the points P and W in the directions WM, PL in the same plane, and in equilibrio, are to each other inversely as the perpendiculars let fall from F upon their directions. PLATE VII. FIG. LXI.

DEM. From F as a center, with the longer perpendicular FL , describe a circular arc cutting the direction of W in D , and, because the efficacy of these forces is the same to whatever points of their directions they are applied, let them be applied at L and D . Let DE represent the magnitude of W and be resolved into two forces, one DG in the direction FD , and the other EG perpendicular to it; and, because DG has no effect in making DF revolve, nor consequently the lever which makes an invariable angle with it, and EG and P act at equal distances from F , in directions perpendicular to those distances, and are in equilibrio, P is equal to EG , and $W : P :: DE : EG :: DF (FL) : FM$, from similar triangles. Q. E. D.

273. Cor. 1. In any lever FP or FHK , those parts of P and W , which are opposite to each other, are inversely as their rectilinear distances from F ; for let WA and PD , be the respective magnitudes of W and P , which act at W and P , and resolving each into two forces, WB, PE , coincident with the lever, or arm FP which makes an invariable angle with it, and AB, DE , parallel and opposite, and drawing the perpendiculars AC, DG , to FP , and FM, FL , to the directions; it appears, from similar triangles, that $DE : AB :: DG : AC :: \frac{DP \times FL}{FP} :: \frac{WA \times FM}{FW} (\because \frac{1}{FP} : \frac{1}{FW} \text{ because } DP \times FL = WA \times FM \text{ from this proposition.})$ FIG. LXII.

274. Cor. 2. If the directions of P and W , acting upon the arms of a straight lever, and in equilibrio, be parallel, they are therefore inversely as their distances, or inversely as the portions of any line drawn through F and terminated by their directions (273). This follows also from the proposition, because the distances, or the segments of any right line on each side of F and terminated by their directions, are, in this supposition, as the perpendiculars let fall from F upon their directions.

276. Cor. 3. The velocity of any point either in a straight, or curved lever, varies as its rectilineal distance from the center of motion F ; for all points describe similar circular arcs, having their centers in F . If P and W act upon the same right line, and their directions be perpendicular, or inclined in the same angle, to their rectilineal distances, their velocities will be as those distances, being measured by the bases of similar triangles which are described in the same time: their velocities are therefore to each other inversely as the opposite parts of P and W (273).

FIG.
LXIII.

277. Cor. 4. If a lever be moveable about an axis AB , or fixed to an axis which is moveable about two centers A and B , and perpendiculars PF , WF , be drawn to this axis from P and W , to which their directions are perpendicular and in the same plane, an equilibrium obtains when $P : W$ inversely as their perpendicular distances from the axis.

FIG.
LXII.

278. Cor. 5. There will be an equilibrium upon the lever when $P \times FP \times \sin. \angle FPL$ (or the angle which P 's direction makes with FP) is equal to $W \times FW \times \sin. \angle FWM$ (or the angle which W 's direction makes with FW ; for $P : W :: FM : FL :: \frac{FW \times \sin. \angle FWM}{\text{rad.}}$
 $:\frac{FP \times \sin. \angle FPL}{\text{rad.}}$; and $F \times FP \times \sin. \angle FPL = W \times FW \times \sin. \angle FWM$, supposing the radius to be given.

279. Cor.

279. Cor. 6. The intensities or moments of any powers $A, B, C, \&c.$ whose directions are parallel, vary as their magnitudes multiplied into their distances from the center of motion; for let $Q, R, S, \&c.$ acting at the same point E , in directions parallel to those of $A, B, C, \&c.$ be in equilibrio with them respectively, or (Cor. 2.) let $A \times AF = Q \times EF, B \times BF = R \times EF, C \times CF = S \times EF, \&c.$; and the efficacy of $A, B, C, \&c.$ to make the lever revolve, or support any power acting against them at E , is evidently as $Q + R + S, \&c.$ or as $Q + R + S \times EF$, or, substituting their equals, as $A \times AF + B \times BF + C \times CF, \&c.$

280. Cor. 7. The intensities or moments of $A, B, C, \&c.$ the sines of whose directions with their rectilineal distances to the same radius are $a, b, c, \&c.$ respectively, will be as $A \times a \times AF + B \times BF \times b + C \times CF \times c, \&c.$; for if Q, R, S act at the same point E , in the same direction, the sine of whose inclination to $EF = z$, and be in equilibrio with them, $Q + R + S \times EF \times z = A \times AF \times a + B \times BF \times b + C \times CF \times c, \&c.$ (278).

281. Cor. 8. If more than two powers act upon a lever, there will be an equilibrium when the sum of the products arising from multiplying each into the perpendicular distance of its direction from the center of motion; or, if their directions be parallel and the lever straight, into its distance from F , on one side, is equal to the sum of the products on the other side. Whatever be the form of the lever, the value of the perpendicular may be substituted for it, and an equilibrium obtains when the sums of the products on each side of F are equal.

282. Cor. 9. Because the efficacy of P and W is the same to whatever part of their direction they are applied, a bended or curved lever may be reduced to a straight one, making an invariable angle with it, and P and W may therefore be always supposed to act in the same right line.

PLATE
VIII.
FIG.
LXV.

283. PROP. *When any number of levers WC, CE, EP, &c. are combined together in the same direction, the ratio of P to W, acting in the same plane at their extremities, in parallel directions, and in equilibrio, is that of $WB \times CD \times EF : BC \times DE \times FP$.*

DEM. Let the forces Q, R, P , acting at the points C, E, P , in directions parallel to those of W and P , be respectively in equilibrio with W, Q, R , and consequently (274) $W : Q :: BC : WB$

$$Q : R :: DE : CD$$

$$R : P :: FP : EF;$$

and, by composition of ratios, $W : P :: BC \times DE \times FP : WB \times CD \times EF$. Q. E. D.

284. Cor. 1. If P and W act in different directions, and Q and R act also in any other different directions, and Q, R, P be respectively in equilibrio with W, Q, R , perpendiculars, let fall from the centers of motion, B, D, F , upon their directions, are to be substituted for the distances.

285. Cor. 2. If any of the forces, in this and the preceding proposition and corollaries, act in different planes, they are to be reduced to the same plane by resolving each into two forces, one in that plane and the other perpendicular to it, of which the former only are effective, and are to be used in the several analogies.

FIG.
LXVI.

286. PROP. *If lines be drawn from F parallel to the directions of P and W which meet in A; P, W, and pressure upon F (Pr) are to each other as the sides and diagonal AB, AC, AF, of the parallelogram CABF, respectively.*

DEM. From F let fall the perpendiculars FL, FM , upon the directions of P and W , and the triangles FCM, FLB (having a right angle in each, and the angles FCM, FBL , either equal to or the

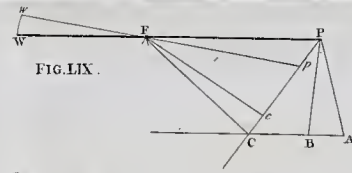


FIG. LIX.

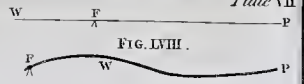


FIG. LVIII.

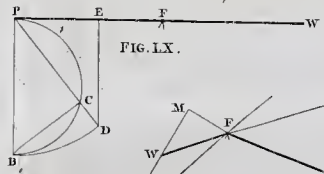


FIG. LX.

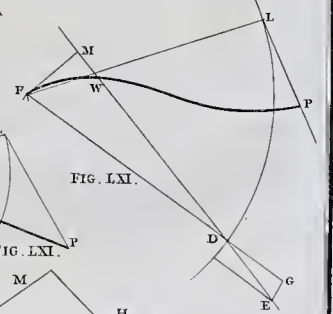


FIG. LXI.

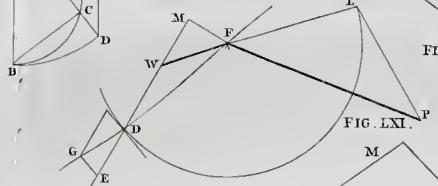


FIG. LXI.

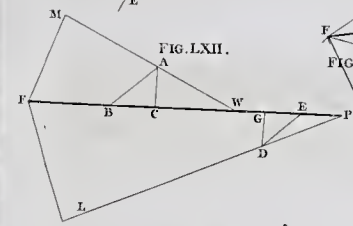


FIG. LXII.

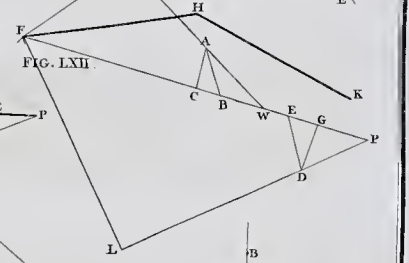


FIG. LXII.

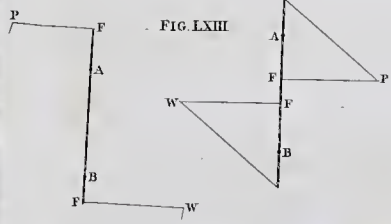


FIG. LXIII.

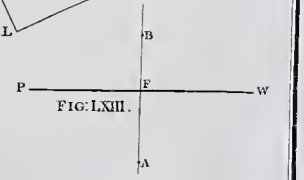


FIG. LXIII.

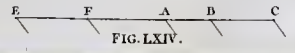


FIG. LXIV.

the supplements of CAB) are similar; therefore $P : W :: FM : FL :: FC : FB$, and a force, whose quantity and direction are AF , is equivalent to P and W (185). Q. E. D.

287. Cor. 1. A power therefore acting at F , whose magnitude is to $P + W$ as AF to $AB + AC$, and whose directions are FA , AB and AC , respectively, will prevent all motion.

288. Cor. 2. The magnitude of P , W or Pr , is as the sine of the angle formed by the directions of the other two; for

$$Pr : P :: FA : AB :: \sin. \text{ of } \angle FBA \text{ or } BAC : \sin. \text{ of } \angle AFB \text{ or } FAC;$$

$$Pr : W :: FA : AC :: \sin. \text{ of } \angle FCA \text{ or } \angle CAB : \sin. \text{ of } \angle CFA \text{ or } FAB;$$

$$\text{and } W : P :: AC : FC :: \sin. \text{ of } \angle CFA \text{ or } FAB : \sin. \text{ of } \angle CAF.$$

289. Cor. 3. Any two of these forces, pressure upon F , P , and W , are to each other inversely as the perpendiculars let fall upon their directions from any point F in the direction of the third; for $Pr : P :: \sin. \text{ of } \angle FCA \text{ or } WAH : \sin. \text{ of } \angle CAF :: WH : WI$ (supposing WH and WI to be perpendicular to the directions of P and Pr , and AW to be the radius); and $Pr : W :: \sin. \text{ of } \angle FBA \text{ or } PAK : \sin. \text{ of } \angle BAF \text{ or } BAN :: PK : PN$, supposing PK and PN to be perpendicular, respectively, to the directions of W and Pr , and AP to be the radius, &c. And if the lever be straight and the directions of P and W parallel, the magnitudes of P , W and Pr , are as the distances of the other two, the perpendiculars then becoming the distances.

290. Cor. 4. If the extremities of the perpendiculars FL , FM , be joined by a right line LM , Pr , P and W are to each other respectively as LM , FM and FL , for this triangle is similar to FAB , as easily appears by describing a circle upon FA as a diameter.

Q

FIG.
LXVII.

291. PROP. *The distance from F, quantity and direction, of any two forces Q and R acting upon a straight lever in the same plane, being given, to find the magnitude, direction and distance of a force equivalent to them.*

Let the directions of Q and R meet in A , and taking $AB : AD :: Q : R$, and completing the parallelogram, the diagonal AE is the magnitude and direction of a force equivalent to them (185). But QR and the angles AQR , ARQ being given, AQ , AR , and the angle QAR may be found; and AB , AD , and the angle BAD being known, AE , the magnitude of the combined force, and the angle BAE may be found; and QA , and the angles AQN , QAN being found, the angle QNA , or inclination of AN , to FN and QN , are known. Q. E. I.

FIG.
LXVIII.

292. Cor. 1. If any number of forces Q, R, S in the same plane, whose magnitudes and directions are AB, AC, ER , act upon the lever FW , and be in equilibrio with any other forces T, V, W , whose quantities and directions are HG, HI, LN , they may be reduced to two which are in equilibrio; for AB and AC are equivalent to AD , and taking EF equal to AD , EF and ER are equivalent to EQ ; and in the same manner HG, HI, LN are equivalent to LP , which is in equilibrio with EQ , because the lever is at rest. And the directions, and distances from F , of LP and EQ , are found as in the proposition.

FIG.
LXIX.

293. Cor. 2. Any number of forces in the same plane, whose magnitudes and directions are AB, AC, EX, HG, HI, LN , may be reduced to one, which is equal to the pressure upon the fulcrum; for, supposing the forces EQ , and LP , resulting from the other forces combined, to act at their intersection S , a third force equivalent to them, and consequently to the pressure upon F , must pass through S (185), and, taking SU, SV , respectively equal to EQ, LP , and completing the parallelogram, SV will be its magnitude and direction.

294. Cor.

294. Cor. 3. If any number of forces, whose quantities and directions are given, act upon a lever, the position of a fulcrum, about which they will be in equilibrio, and the quantity and direction of pressure upon it, may be found; for, let the directions of Q and R meet in A , and taking $AB : AC :: Q : R$, and, completing the parallelogram, its diagonal AD produced, will cut the lever in a point, which being supported, Q and R will be in equilibrio (286); and combining AD with S , and the force, resulting from these, with another, &c. the diagonal will always intersect the lever in a point F , about which they will be in equilibrio.

FIG.
LXVIII.

295. LEMMA. *If right lines be drawn from any point P to the extremities of the diagonal, and sides, of the parallelogram $ABCD$, the triangle PAC , having the diagonal for its base, is equal to the difference, or sum of the triangles PAB , PAD , having the sides for their bases, according as P is situated between the lines forming the angle BAD , or those which form its supplement to two right angles.*

FIG.
LXX.

DEM. CASE I. Let P be situated between the lines forming the angle BAD ; and, drawing Pnm perpendicular to AB or CD , the triangle $APC = ADPC - ADP$; but $ADPC = ADC + DPC = AB \times \frac{mn}{2} + DC \times \frac{Pn}{2} = AB \times \frac{Pm}{2} = APB$; therefore $APC = APB - APD$.

CASE II. Let P be placed between the lines forming the supplemental angle to BAD , and the triangles $ADC \pm DPC = \frac{AB}{2} \times mn \pm Pn (Pm) = APB$; and, adding APD to both, $APC = APB + APD$. Q.E.D.

FIG.
LXXI.

PLATE

IX.

FIG.
LXXII.

296. Cor. 1. If P be in one of the sides containing the angle BAD , the triangle $APC = ADC + DPC = \frac{AB}{2} \times mn + Pn$
(Pm) = PAB .

FIG.
LXXIII.

297. Cor. 2. If P be in the diagonal, the triangle $PAB = ABC$
 $\pm BPC = BC \times \frac{mn}{2} \pm BC \times \frac{Pn}{2} = BC \times \frac{Pm}{2} = PAD$. The
upper sign is to be used when P is without the figure, the lower
when it is within.

298. Cor. 3. If perpendiculars be drawn from P to the diagonal and sides, Pd, Pm, Pq ; $AC \times Pd = AB \times Pm \pm AD \times Pq$; and when P is in the diagonal AC , $AB \times Pq = AD \times Pm$.

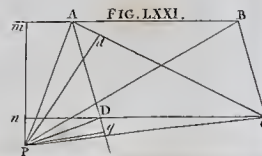
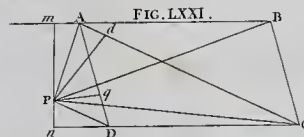
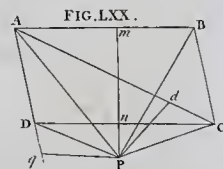
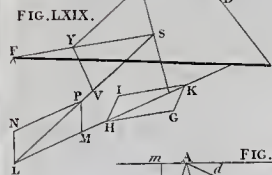
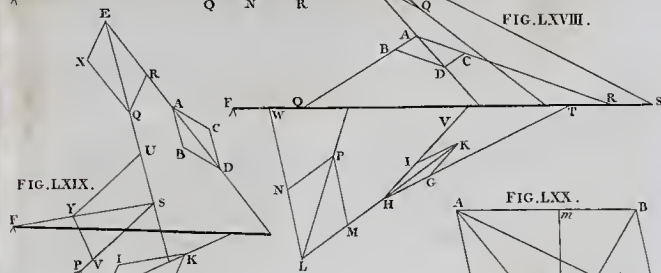
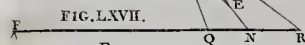
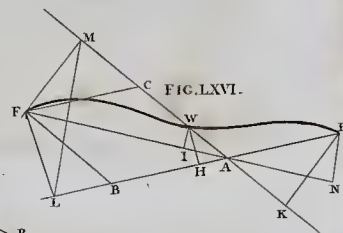
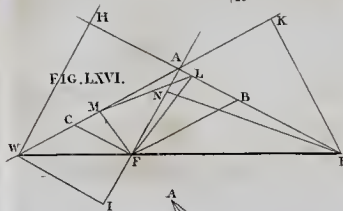
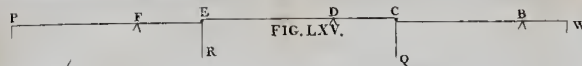
FIG.
LXXIV.

299. PROP. *If there be any number of forces Q, R, S, T , in the same plane, which are combined as before, the sums of the products, arising from multiplying each into the perpendicular distance of its direction from any point F , in the diagonal CL , are equal on each side of it.*

DEM. Let CD, CG, CH, CK be the relative magnitudes of Q, R, S, T , respectively, and compounding them, and drawing perpendiculars, from any point F in the last diagonal, upon their directions, viz. Fa, Fb, Fc, Fd, Fg, Fh ; from the preceding lemma, $CK \times Fh = CI \times Fg = CE \times Fc - CH \times Fd = CD \times Fa + CG \times Fb - CH \times Fd$, and consequently CK or $T \times Fh + CH$ or $S \times Fd = CD$ or $Q \times Fa + CG$ or $R \times Fb$. Q. E. D.

300. Cor. A lever therefore passing through any point F in the diagonal LC , produced in any direction, and in the same plane with the forces which act upon it, will be in equilibrio.

301. LEMMA.





301. LEMMA. *If a right line PW revolve round a fixed point F, and lines are drawn from its extremities to another fixed point Q, to which perpendiculars FL, FM, and Fl, Fm are drawn from F, and W is supposed to ascend, FL has to FM a greater ratio than Fl : Fm.*

DEM. Take $Pr = py = FW$, and, drawing the perpendiculars rv, yz ; $FL : rv :: \sin. \angle FWL : \sin. \angle rPv :: P\mathcal{Q} : W\mathcal{Q}$; $rv : FM :: Pr(FW) : FP$, and comp. $FL : FM :: P\mathcal{Q} \times FW : W\mathcal{Q} \times FP$. By a similar process, $Fm : Fl :: w\mathcal{Q} \times FP : p\mathcal{Q} \times FW$; and adding these analogies together, $FL \times Fm : FM \times Fl :: P\mathcal{Q} \times FW \times w\mathcal{Q} \times FP : W\mathcal{Q} \times FP \times p\mathcal{Q} \times FW :: P\mathcal{Q} \times w\mathcal{Q} : W\mathcal{Q} \times p\mathcal{Q}$. But $P\mathcal{Q}$ is greater than $p\mathcal{Q}$, and $w\mathcal{Q}$ than $W\mathcal{Q}$, and consequently $P\mathcal{Q} \times w\mathcal{Q}$ is greater than $W\mathcal{Q} \times p\mathcal{Q}$, and $FL \times Fm$ than $FM \times Fl$, and the ratio of $FL : FM$ is greater than that of $Fl : Fm$. Q. E. D.

302. PROP. *If any two powers P and W, whose directions always meet in the same point Q, be in equilibrio upon the lever PW in any one position of PW, they cannot be in equilibrio when it is in any other position pw.*

DEM. Let the lever revolve, and be in any other situation pw , and because $P : W :: FL : FM$, or in a greater ratio than $Fl : Fm$ (301) P will preponderate. Q. E. D.

303. Cor. 1. It is evident that the lever cannot rest till it pass through the points F, \mathcal{Q} .

304. Cor. 2. If the directions of P and W be parallel to each other, they will be in equilibrio in any position of the lever, because the perpendiculars drawn from F to their directions are always as the distances, and consequently in a given ratio to each other. And, for the same reason, if there be ever so many forces, acting in parallel directions upon the arms of a straight lever,
and

and in equilibrio in any one situation of the lever, they will be in equilibrio in every situation of it.

FIG.
LXXVI.

305. PROP. *If the center of motion F be placed above the straight lever PW, and P and W, acting always in parallel directions, equilibrate in any position PW, they do not equilibrate in any other, pw.*

DEM. From F draw FM and Fm perpendicular to the lever, and $P : W :: WM : PM :: wm : pm$, or in a less ratio than that of $Lw : Lp$ or $LV : LR$, and consequently much less than that of the perpendiculars from the center of motion F upon the directions, or $MV : MR$. Q. E. D.

306. Cor. 1. Because $P : W$ in a less ratio than that of $MV : MR$, $P \times MR$ is less than $W \times MV$, and consequently W will descend.

PLATE

X.
FIG.
LXXVII.

307. Cor. 2. If F be placed on the other side of the lever, the descending body will preponderate; for $P : W :: WM : PM :: mw : mp$, and consequently $P : W$ in a greater ratio than that of $Lw : Lp$ or $LV : LR$, and therefore much greater than that of the perpendiculars upon the directions or $MV : MR$; and $P \times MR$ is greater than that of $W \times MV$.

FIG.
LXXVIII.

308. PROP. *If P and W act always at the same distance from the lever, in directions parallel to FM, and in equilibrio about F in the position PFW, and the lever be moved, the descending body will continue to preponderate.*

DEM. Let FN , ps , wr , be drawn parallel to the directions in which the bodies act, pQw be the lever, and Fn be the position of FN ; and from similar Δs $Qr : Qs :: Qw : Qp$, and therefore $Qr : Qs$ in a greater ratio than $nw : np$ or $NW : NP$ or $WM : PL$ or

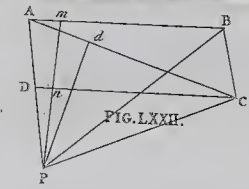


FIG. LXXII.

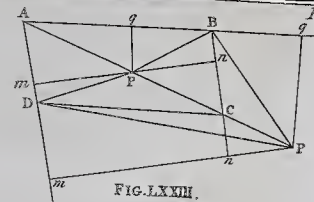


FIG. LXXIII.

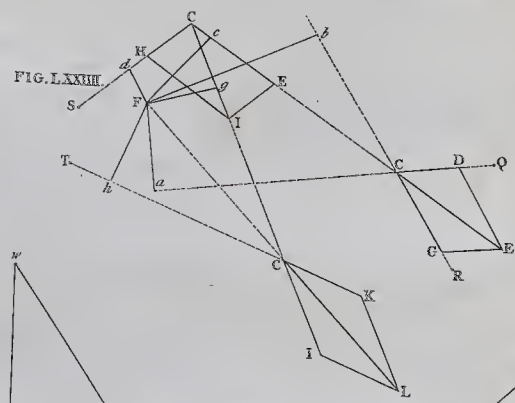


FIG. LXXIII.

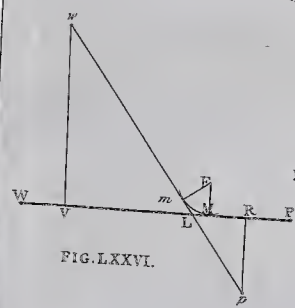


FIG. LXXVI.

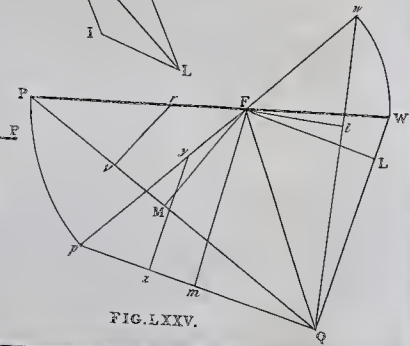


FIG. LXXV.

PL or P to W ; but $Nr : Ns$ in a greater ratio than $Qr : Qs$, and consequently $P : W$ in a less ratio than $Nr : Ns$, and the body W will continue to preponderate.

309. Cor. When the bodies are placed under the lever, the descending body will continue to descend, because $P : W$ in a less ratio than $Nr : Ns$, as easily appears by turning the figure.

S C H O L I U M.

310. In all communications of motion by impact, the quantities of motion lost and gained being equal and opposite, the quantity of motion estimated in the same direction is invariable, and the quantities of matter vary inversely as the velocities lost and gained; and if two bodies A and B act upon a lever, or any other machine, they are so connected that A cannot descend without making B ascend with the same quantity of motion, and their quantities of matter are therefore inversely as each other. These cases, having such marks of coincidence, are inferred to be similar in every respect, and the cause of an equilibrium in the mechanical powers, is often immediately assigned from this equality of momenta; but they are not exactly similar, because when A impinges upon B , some part of its motion is transferred to B , and A 's motion necessarily precedes this communication of motion; but, when they act upon any machine, the descending body A cannot be said to communicate any part of its motion to B ascending; because, from their connection, their motion must necessarily commence and be extinguished together; and besides, the power of the lever ought to be considered. The ratio of the power and weight may however be assigned, in every machine, from their incipient momenta; for P and W , acting upon the arms of a lever, exert a pressure and have a tendency to move; and if a and b be the velocities with which A and B strike the lever, estimated in such directions that their pressures are solely employed in resisting each other's efforts to produce motion, no part being lost by obliquity of direction, $A \times a \times$ into its velocity or distance from the centre of motion.

FIG.
LXXIX.

motion is equal to $B \times b \times$ into its velocity or distance, when there is an equilibrium; and by substituting pressures or P and W , for $A \times a$ and $B \times b$, the ratio of $P : W$ is found to be the same as was before collected from the resolution of motion. As the velocity of any point of a lever varies as its rectilinear distance from the center of motion, all points describing similar circular arcs round it in the same time, if the directions of P and W be perpendicular to their rectilinear distances, their velocities will be the same as those of the points where they act and wholly efficient, and when there is an equilibrium, P will be to W inversely as their rectilinear distances. But if the direction of P or W be enclined to their rectilinear distance, their efficient velocity will not be equal to that of the point where they act; if P 's direction be the line PD , it is evident that P will have two motions whilst the lever revolves, one acceding to, or receding from, F , according as the angle DPF is less or greater than a right angle, and the other producing the rotation of the lever; for, describing a circular arc with F as a center, and FD as radius, the power will act at every intermediate point in PE whilst the lever describes the angle PFD . The motion in the direction of the lever is inefficient, and if PD represent the direction and quantity of A 's velocity, and be resolved into two, PC in the direction of PF , and DC perpendicular to it, this last only is efficient; and if the angle at F be very small, $A \times DC \times$ its perpendicular velocity or into FD , or $\frac{A \times FP \times PD \times FL}{FP}$, supposing FL to be

perpendicular to the direction, is A 's efficacy to turn the lever. The effective part of B 's velocity being found in the same manner, and these values of a and b being substituted for them in the supposition of an equality of moments; whatever be their directions, the ratio of P to W is the same as that discovered by other principles. Thus in an equilibrium $A \times PD \times FL$ is given, and consequently

$A \times PD$ or the pressure of P in the line PD is as $\frac{1}{FL}$. The de-

monstration of this fundamental proposition, ascribed to Archimedes, depends upon this principle, that if a number of weights be suspended upon the arms of a lever, at points equidistant from each other, whether on the same side of the fulcrum or not, their

efficacy

efficacy to make the lever revolve is the same as if they were united in a point bisecting the distance of the points of suspension. Mr. Huygens says, that many fruitless attempts have been made to remedy the defects of this demonstration, and proposes another founded on this principle, that when two equal bodies are placed upon the arms of a lever, that which is most remote from the fulcrum will preponderate. But this principle is not more evident than that of Archimedes, and besides, his process is prolix and tiresome. Mr. Maclaurin hath given a demonstration of this proposition when the arms of the lever are commensurate, and his method might easily be extended to cases in which they are incommensurate, and be made general; but this would add to the length of a process already very long. The only principle in Sir I. Newton's demonstration that hath been controverted is this, which is taken for granted, "that the same power will have the same effect to whatever point of the direction, in which it acts, it be applied;" yet no doubts are entertained of the truth of it, though, perhaps, from its simplicity and intuitive evidence, it cannot be demonstrated by more simple principles. If the line PC and the angle PFC be invariable, the radii PF , CF being fixed to PC at the points P and C , it is evident, that two equal forces P and Q , acting upon the points P and C in the directions PC and CP , will destroy each other, and the line PC , and consequently PF and CF , will be quiescent. If therefore P make a body describe the arc Ww in any small time, Q will make it describe wW in the same time, or they would not destroy each other's effects. Equal forces therefore P and Q , acting at different points of the same direction, in the same time make the radii PF , CF describe angles at F , PFp and CFc equal to WFw , and therefore to each other. Though this demonstration of Newton be general, concise, and perfectly satisfactory, the great utility of the proposition may possibly render other demonstrations of it not undeserving of attention.

PLATE
VII.
FIG.
LIX.

DEMONSTRATION OF ARCHIMEDES.

FIG.
LXXX.

Let AB be a homogeneous cylinder, and C the bisection of its axis, and it is evident, that if a fulcrum, or power equal and opposite to the pressure upon it, be applied at C , the parts AC and CB will be in equilibrio. Let any point D be taken, and, bisecting AD in E , and DB in F , it is clear that two powers, respectively equal to the weights of AD and DB , applied at E and F , will support them, and have the same effect with the fulcrum, or power, applied at C , and be in equilibrio about C : but $CE = CA - AE = \frac{AB - AD}{2} = \frac{DB}{2}$; and $CF = CB - BF = \frac{AB - DB}{2} = \frac{AD}{2}$; and consequently $CE : CF :: DB : AD ::$ force applied at F : force applied at E . Q. E. D.

FIG.
LXXXI.

311. LEMMA. *If from any point P , in the diagonal of a parallelogram, $ABCD$, two lines Pm and Pq , be drawn perpendicular to the sides, the perpendiculars and sides are inversely as each other.*

DEM. Draw CE, CF , perpendicular to the sides, and, from similar triangles, $Pq : CF :: AP : AC :: Pm : CE$; therefore $Pq : Pm :: CF : CE :: CD (AB) : CB (AD)$ (sim. triangles). Or, this follows from (298), where it is proved, that $AB \times Pm = AD \times Pq$, and consequently, $Pm : Pq :: AD : AB$. Q. E. D.

FIG.
LXXXII.

312. PROP. *If any three forces W, P, Z , whose magnitudes and directions are AB, AC, AD , act upon the lever WP , which is at rest, W is to P inversely as the perpendiculars let fall from F upon their directions.*

DEM.

DEMONSTRATION OF ARCHIMEDES.

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DEM. The directions of these forces must be in the same plane, and meet in the same point, (216), and are respectively as the sides and diagonal of a parallelogram parallel to their directions; and $W : P :: AB : AC$ (196) $:: FL : FM$ (311). The effect will evidently be the same when a fulcrum is applied at F instead of the force Z . Q.E.D.

*Dr. Hamilton's demonstration of this proposition depends upon the same principle with the above. Let the forces W, P, Z , act upon the inflexible line WP at the points W, P, F , and let the directions of P and W meet in C , and the direction of a force Z , equivalent to them, must pass through C . Therefore P, W, Z , are to each other as the sides and diagonal of a parallelogram respectively parallel to their directions, FA, FB, FC , and $W : P :: AC : BC :: \sin. \angle AFC (FCB) :: \sin. \angle ACF :: FM : FL$.

FIG.
LXXXIII.

$W : Z :: AC : FC :: \sin. \angle FCB : \sin. \angle FBC$ or $BCA :: P\mathcal{Q} : PN$, supposing $P\mathcal{Q}$ and PN to be perpendicular to the directions of Z and W .

$P : Z :: FA : FC :: \sin. \angle FCA : \sin. \angle FAC$ or $ACB :: WG : WR$, supposing WG and WR to be perpendicular to the directions of Z and P .

The parts of W and P , which act in directions exactly opposite to those of Z , are found by resolving AC and $\mathcal{Q}A (BC)$ into two forces, one parallel to the right line joining W, P , as Am and Bn , and the other parallel to FC as $\mathcal{Q}m, \mathcal{Q}n$. The opposite parts Am, Bn , are equal and destroy each other, and the conspiring parts $\mathcal{Q}m + \mathcal{Q}n$ must be equal to Z ; and consequently when two forces W and P are in equilibrio with a third force Z , and their directions are all parallel,

FIG.
LXXXIV.

$$\begin{aligned} W : P &:: \mathcal{Q}n : Cn :: PB : BC :: P\mathcal{Q} : \mathcal{Q}W; \\ W : Z &:: Cm : \mathcal{Q}C :: CA : CW :: \mathcal{Q}P : PW; \text{ and} \\ P : Z &:: Cn : C\mathcal{Q} :: CB : CP :: W\mathcal{Q} : WP. \end{aligned}$$

† B A L A N C E.

313. DEF. The ancient balance, commonly called the *statera romana*, or steelyard, is a lever of the first kind supported at the point F , placed

PLATE
XI.
FIG.
near LXXXV.

* Essay on the Principles of Mechanics. † Helsham, Lect. VI. Muschenb. Ch. VIII. § CCCLXXXIII. Desaguliers, pag. 95.

near one extremity, about which the brachia FA , FN equiponderate. On one side of F at the extremity A , an unknown weight, W , is suspended, and on the other side a known weight, P , is moveable upon the arm FN , which is divided into parts equal to FA , each of which is also divided into 10, 100, &c. equal parts.

314. Cor. If P be at x , the fourth division from I , which is the n^{th} division from F , when an equilibrium obtains between it and W , and IK be divided into m equal parts, the weight of W will be equal to $P \times n + \frac{4}{m}$; for $W:P :: Fx:FA(274) :: FA \times n + \frac{4}{m} : FA :: n + \frac{4}{m} : 1$, and consequently $W = P \times n + \frac{4}{m}$.

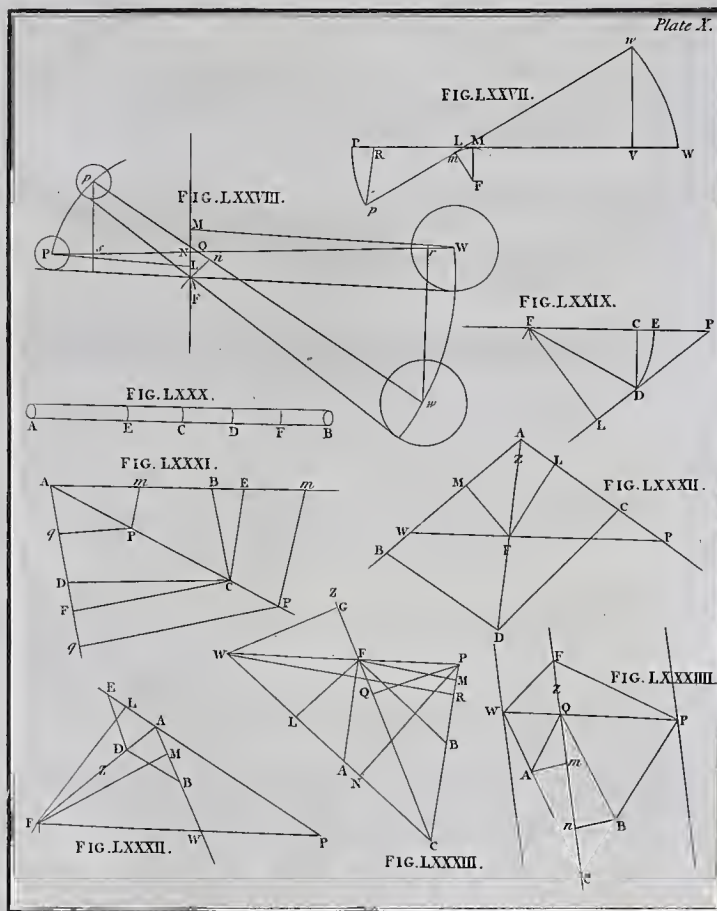
FIG.
LXXXVI.

315. DEF. The common, or modern, balance, or a pair of scales, is a lever of the first kind, as AB , supported at its bisection F : to the extremities A and B are suspended basins or scales, and the brachia and basins, on each side of F , are supposed to equiponderate.

316. Cor. 1. If any known weight, P , placed in one scale equiponderate with one unknown, W , placed in the other, the weight of W is known, being equal to that of P ; for $P:W :: AF:BF(274) :: 1:1$, and consequently $P = W$.

317. Cor. 2. A balance, whose arms are unequal in length, is fallacious; for if AF and BF be unequal, P and W , when in equilibrio, cannot be equal, for $W = \frac{P \times BF}{AF}$. But the relation of $BF:AF$ being known, W is also known.

318. Cor. 3. If a man, whose weight is equal to W , standing in one scale and in equilibrio with P placed in the other, press the



the beam upwards in D with a force equal to \mathcal{Q} , the diminution of W 's moment is equal to $\mathcal{Q} \times FD$; and because the reaction at the scale is equal to \mathcal{Q} , the encrease of W 's moment is equal to $\mathcal{Q} \times FA$, and consequently W will descend with a force equal to $\mathcal{Q} \times AD$. If the pressure be upwards at E , W will descend with a force, resulting from this pressure, equal to $\mathcal{Q} \times EF$, and, from the reaction, with a force equal to $\mathcal{Q} \times FA$; and therefore the whole force of descent is equal to $\mathcal{Q} \times EA$. When the pressure is downwards at D , the encrease of W 's moment is equal to $\mathcal{Q} \times FD$, and the diminution of its moment $= \mathcal{Q} \times FA$; and, consequently, W will ascend with a force equal to $W \times DA$. If the pressure be downwards at E , the diminution of W 's moment, or encrease of P 's moment, is equal to $\mathcal{Q} \times EF$, and a part, \mathcal{Q}_2 of W 's weight being transferred to E , the diminution of its moment, on that account, is equal to $\mathcal{Q} \times FA$; and consequently the whole diminution of W 's moment, or force of P 's ascent is equal to $\mathcal{Q} \times EA$.

319. Cor. 4. If the center of motion C be in the right line joining the centers of suspension, the equilibrium of equal weights P and W will obtain in every position; the perpendiculars let fall from C upon the directions being always equal to each other. But when C is above or below WP , an equilibrium of equal weights does not obtain, unless WP coincide with the horizontal line AB . When WP coincides with AB , the perpendiculars let fall from C upon the directions of W and P are equal to GB and GA , CG being perpendicular to AB ; but when the balance is in any other position WP , the perpendicular CI is greater than CH , because gL , which is less than CI , is equal to gM , which is greater than CH . W will descend and continue to vibrate till its motion be destroyed by friction. This corollary is also deducible from (305).

FIG.
LXXXVII.

320. Cor. 5. If P and W be unequal, and C , be in the right line WP , the heavier of them will descend till WP be perpendicular to the horizon, or, if the center of motion be not in WP , till $P \times CH = W \times CI$.

321. Cor.

FIG.
LXXXVI.

321. Cor. 6. In a balance whose arms are unequal, the weight of W may be still ascertained; for let W , suspended at A , be in equilibrio with a known weight \mathcal{Q} , suspended at B , and $W \times FA = \mathcal{Q} \times FB$, and, suspended at B , let it be in equilibrio with a known weight R suspended at A , and $W \times FB = R \times FA$; consequently, by multiplying these equations together, $W^2 \times FA \times FB = \mathcal{Q} \times R \times FA \times FB$, and $W = \sqrt{\mathcal{Q} \times R}$.

FIG.
LXXXVIII.

322. *Cor. 7. If the beam of the balance be supposed to have weight and be similar and homogeneous in every part, its center of gravity is in the bisection F ; but if it be not homogeneous, or the center of motion be not in the bisection, let G be its center of gravity, and an equilibrium will obtain when $P \times PF + B$ (weight of the balance) $\times GF = W \times FW$.

323. Cor. 8. If from F , a style FD , perpendicular to WP , be raised, the equilibrium of the balance will be affected by it, except it be in an horizontal situation, the moment of the style being measured by its weight multiplied into the distance of its center of gravity from the line FH , perpendicular to the horizon. But the equilibrium is restored by continuing the style to the other side of F , so that the moments on each side may be equal and opposite.

324. Cor. 9. If the center of gravity of the balance, scales, and weights, be in the center of motion, F , an equilibrium obtains in every position of the balance; but if this center be above or below F , the balance cannot be quiescent till the right line joining F and this center be perpendicular to the horizon. The best position therefore of the center of gravity is below F , and as little below it as possible, that the arcs described by it, during its vibrations, may be small and soon described. The points of suspension should

• These corollaries may be omitted till the chapter upon the center of gravity be read.

should be in the same right line with the center of motion, which ought accurately to bisect their distance (317).

325. Cor. 10. The arms of the balance should be as long as can be used conveniently; because the moment of a given body varies as its distance from the fulcrum, and, therefore, the greater the distance, the more distinguishable will be the moment arising from any small difference between P and W . And to distinguish very minute differences of weight, the friction upon the axis, in the motion of the beam and scales, ought to be as little as possible.

C H A P. VIII.

* W H E E L A N D A X I S.

FIG.
LXXXIX.

326. DEF. *A WHEEL and axis or axis in peritrochio, is a machine composed of a circular wheel, in whose center a cylindrical axis is inserted and fixed; and the wheel revolving by the action of a power P, the axis, whose extremities are supported, revolves with it, and the rope, to which a body W is appended, is tied to the axis and wrapped round it during its motion.*

327. Cor. It is evident that every point of the axis B, G, &c. describes a circle round its corresponding center y, z, &c. in the time of a revolution of the wheel; and that any points of the wheel and axis, A and B describe similar arcs of circles in the same time, and, consequently, their velocities are as the peripheries or radii of the circle described by them.

328. PROP. *If the directions of P and W be perpendicular to the radii of the wheel and axis respectively, they are in equilibrio when $P : W :: \text{radius of the axis} : \text{radius of the wheel}$.*

DEM. The same power is required to support W to whatever point of the axis it be applied, because the distance from the corresponding center of motion is the same, and the wheel and axis may be reduced to a bent lever, and consequently there will be an equilibrium when $P : W :: W\text{'s distance from the center of motion} : P\text{'s distance} :: \text{radius of the axis} : \text{radius of the wheel}$ (277).
Q. E. D.

This

* Keil's Physics, Lect. X. Helsham, Lect. VII. Muschenb. Ch. VIII. Sect. CCCCXLIII. Emerson, Prop. XXIV. Varignon, Tom. I. Sect. IV.

Otherwise:

This proposition is usually proved by the following process: since the directions of P and W are perpendicular to their respective distances from their centers of motion, they are wholly efficient, and P 's velocity is to W 's velocity as the periphery of the wheel to the periphery of the axis, and consequently, when there is an equilibrium, $P : W ::$ periphery of the axis : periphery of the wheel $::$ radius of the axis : radius of the wheel (272). Q. E. D.

329. Cor. 1. If the thickness of the rope, to which W is appended, be not inconsiderable, it ought not to be neglected; for, when one or more spires of the ropes are folded about the axis, the distance of W 's direction from the center of motion is increased, being equal to the semidiameters of the axis and ropes; and there is an equilibrium when $P : W ::$ the distance of W 's direction from the center of motion : semidiameter of the wheel.

330. Cor. 2. If P and W act in the same plane, and in the directions PD , and WD , meeting in D , and be in equilibrio, they are equivalent to a third force, or pressure upon the axis at A , whose direction meets PD , and WD in D (216); and, producing PD, WD , these three forces are to each other, as the sides DF , DE , and diagonal DG , of the parallelogram EF ; therefore $P : W :: DF : DE$, or drawing AN , AM , perpendicular to WD and FD respectively, $P : W :: AN : AM$ (311).

FIG.
XC.

331. Cor. 3. The pressure upon the axis at A (Pr) : $P :: DG : DF :: \sin. \angle DFG$ or $PDW : \sin. \angle FGD$ or ADW ; $Pr : W :: DG : DE :: \sin. \angle DEG$ or $PDW : \sin. \angle DGE$ or ADP ; and $P : W :: \sin. \angle ADW : \sin. \angle ADP$. When the angle PDW is infinitely small, or PD and WD are parallel, the perpendiculars AN , AM , are to each other as $AW : PA$.

FIG.
XCI.

332. PROP. *In a combination of wheels GH, FK, EL, whose axes are QM, RN, BD, an equilibrium obtains when* $P : W :: QD \times RC \times AB : DG \times CF \times AE$.

DEM. Let Q, R, W , be respectively in equilibrio with P, Q, R ; and $P : Q :: QD : DG$, (328)

$$Q : R :: RC : FC,$$

$R : W :: AB : AE$, and, componendo, $P : W :: QD \times RC \times AB : DG \times FC \times AE$. Q. E. D.

333. Cor. 1. If the ratios of $QD : DG, CR : CF, AB : AE$ be the same, and the number of wheels be n , $P : W :: QD^n : DG^n$; and if this given ratio be that of $1 : r$, $P : W :: 1 : r^n$ and $P = \frac{W}{r^n}$.

FIG.
XCII.

334. Cor. 2. If the peripheries QM, RN , were to touch the peripheries FK, EL , and operate upon them by means of teeth made in each, this analogy would still obtain, and P would be to W as the product of the semidiameters of the axes, to the product of the semidiameters of the wheels. As if XY be a combination of wheels, and W be appended upon the axis YF , and the power P act at P ; $P : W ::$ products of the semidiameters of the pinions $B, D, F : \text{product of the semidiameters of the wheels } A, C, E$.

335. Cor. 3. Because the number of teeth in the wheels and pinions are to each other as their peripheries, or radii, $P : W ::$ semidiameter of the axis to which W is appended multiplied into the number of teeth in the pinions : the length of the lever where P acts multiplied into the number of teeth in the several wheels.

336. Cor. 4. The number of revolutions of a pinion or wheel being inversely as the time of one revolution, or inversely as the periphery or number of teeth in it, the number of revolutions of
e the

the wheel where P acts, is to the number of revolutions, in the same time, of the axis to which W is appended, as the product of the number of teeth in the wheels to the product of the number of teeth in the pinions.

337. PROP. *If a wheel with teeth CDA revolve about C , and impel the wheel SBA round with it, by the action of teeth BD, bd , &c. upon B and b , whose curvature is such that B, b , describe the lines BD, bd , whilst they describe BA and bA respectively, the moments of these wheels are equal.*

FIG.
XCIII.

DEM. Because the points B and D come into contact at A , at the same time, every part of BA is applied to AD , and is therefore equal to it, and consequently B and D move towards it with equal velocities; and therefore if equal weights act at B and D perpendicular to SB, CD , they must be in equilibrio. Q. E. D.

338. Cor. 1. If $CDnB$ be a bent lever whose fulcrum is at C , equal forces, acting at D and B in directions perpendicular to CD, SB , would keep the levers $CDnB$ and SB in equilibrio.

339. Cor. 2. Because the point B moves over BD whilst the arc BA rolls over DA , the figure of the tooth BnD is an epicycloidal arc; and the effect is the same whether B describe the concave or convex side of BnD , and consequently whether the wheel CDA impel SBA by the action of the convex side of the tooth upon B ; or SBA impel CDA round with it by the action of B upon the concave side of the tooth BnD .

340. Cor. 3. If any number of epicycloidal teeth DB, db and Aa , be inserted in the periphery AD , at equal distances from each other, and teeth B, b, A , be inserted in the other wheel AB , at equal distances, and $Ab = Ad$, and $AB = AD$, the teeth

B, b, A , will all act together with equal moments; for because $Ab = Bb$, the velocities and consequently moments of B and b are equal, and they act together, because b is always found in bd and B in BnD .

341. Cor. 4. The effect is the same when the epicycloidal teeth are upon the periphery of the wheel SBA ; for, whilst B and D move to A , B describes the epicycloid BnD , and D describes another epicycloid DmB , whose base is BA .

S C H O L I U M .

342. The teeth should not act upon each other before they arrive at SC , joining their centers, and the machine is more or less complete according to the number of teeth acting together. The action of any tooth should not cease before that of the succeeding tooth begins.

P U L L E Y*.

FIG.
XCV.

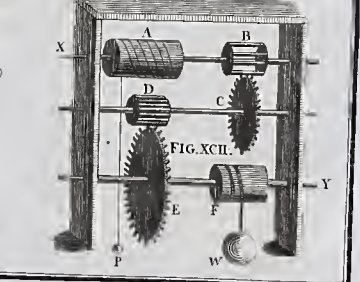
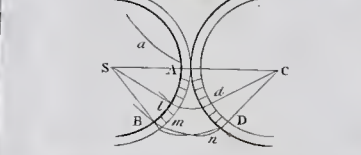
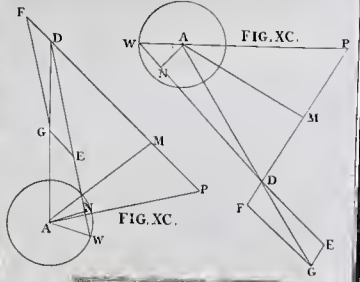
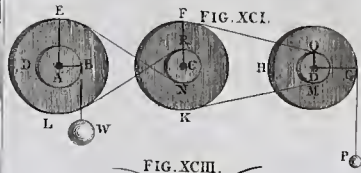
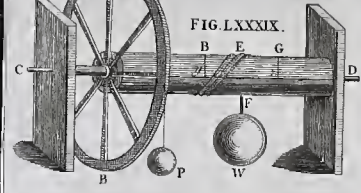
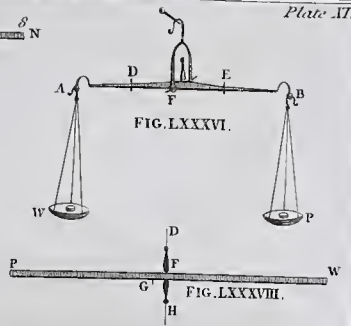
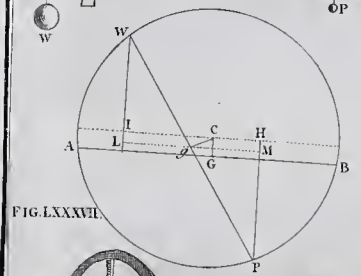
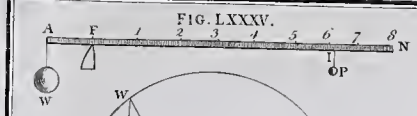
343. DEF. *A pulley is a small circular wheel, as mEn , revolving about an axis passing through its center, by the action of a power which is applied to a rope passing over the pulley.*

344. PROP. *In a single fixed pulley, an equilibrium obtains when the power P , is equal to the weight W .*

DEM. When a rope is stretched and quiescent, it is evident that the tension of every part is the same, otherwise motion would ensue; therefore the tension of Pm is equal to that of Wn and $P = W$. Q. E. D.

S C H O-

* Muschenbroek, Ch. VIII. Sect. CCCCXCIV. Varignon, Tom. I. Sect. III. Keil's Physics, Lect. X. Hellsham, Lect. VII. Hamilton's Principles of Mechanics, pag. 162. Emerson, Prop. XXVII. Desaguliers, pag. 99.



S C H O L I U M.

345. This proposition is sometimes proved by considering the pulley as a lever; for the moments of P and W being the same to whatever parts of their directions they are applied, but, if applied at m and n , mEn is a straight lever whose center of motion, E , is in its bisection, and consequently when there is an equilibrium $P = W$ (274). And the conclusion thus deduced is certainly satisfactory without any other demonstration.

346. Cor. 1. If the same rope pass over any number of fixed pulleys, P is equal to W when there is an equilibrium, because the tension of the rope is, in every part, the same.

347. PROP. If the weight, W , be sustained by a power P , applied to a rope passing over a moveable pulley E , $P : W :: 1 : 2$.

FIG.
XCVI.

DEM. Let the tension of the rope $PA = m$, and that of BC , and DF , is each equal to m , and consequently the tension of EW is equal to $2m$; therefore $P : W :: m : 2m :: 1 : 2$. Q. E. D.

348. Cor. 1. If W be supported by P , applied to a rope passing over any number of moveable pulleys (n), $P : W :: 1 : 2n$; for the number of ropes supporting the weight is equal to $2n$, each supports an equal part of it, and the tension of the rope, to which P is applied, is equal to that of one of them.

349. Cor. 2. If W be supported by P , applied to a rope passing over each pulley in two blocks, to the lower of which W is appended; $P : W :: 1 : \text{number of ropes at the lower block}$. For all the ropes A, B, C, D, E , support W , and each supports an equal part, because their tension is the same, and the tension of each is equal to that of F to which P is applied.

FIG.
XCVII.

350. Cor.

350. Cor 3. If a moveable pulley, L , be annexed to this system, P will support a weight equal to $2 \times W$; for the tensions of the ropes G and H are equal and they both sustain the weight, which is therefore equal to $2 \times W$.

351. Cor. 4. It is evident that P acquires no moment or quantity of motion by this distribution of the weight amongst a number of pulleys; for the velocity of P is to that of W as the number of ropes, supporting the weight, to unity. If W be elevated through any space equal to s , any point in each of the ropes supporting it, must move through a space equal to s , and P consequently through a space equal to s multiplied into the number of strings.

FIG.
XCVIII.

352. PROP. *If the body W be supported by the power P , in a system composed of one fixed, and any number of moveable pulleys, each having a separate string, $P : W :: \text{unity} : \text{that power of } 2 \text{ whose index is the number of moveable pulleys.}$*

DEM. Let the tension of the string, to which P is applied, be equal to m , and the tension of A and B is each equal to m ; that of C and D is each equal to $2m$; that of E and F is each equal to $4m$; and that of GH to $8m$; therefore $P : W :: m : 8m :: 1 : 8$; and, if the number of moveable pulleys be equal to n , it is evident that the tension of the string supporting W is equal to the n^{th} term of the geometric series $2, 4, 8, \&c.$ or to 2^n and $P : W :: 1 : 2^n$. Q. E. D.

353. Cor. 1. $P \times 2^n = W$ and $P = \frac{W}{2^n}$; and if the number of moveable pulleys and P or W be given, the number of moveable pulleys, or n , may be found.

354. Cor.

354. Cor. 2. The parts of W , sustained by the several moveable pulleys, &c. are to each other as 2, 4, 8, 16, &c.

355. PROP. *If the body, W , be sustained by the power, P , in a system of pulleys where the rope passing over each pulley is immediately fixed to W , $P : W ::$ unity : that power of two, diminished by unity, whose exponent is the number of ropes, or number of moveable pulleys increased by unity.*

FIG.
XCIX.

DEM. Let the tension of the rope, to which P is applied, be equal to m , and it is evident that the tensions of the ropes G, B, D, F , are equal to $m, 2m, 4m, 8m$, respectively; but these ropes entirely support W , and consequently $P : W :: m : 1 + 2 + 4 + 8 + m :: 1 : 15$. If the number of ropes fixed to W be equal to n , it is evident that $P : W :: 1 : 1 + 2 + 4 + 8, \&c.$ continued to n terms $:: 1 : 2^n - 1$. Q. E. D.

356. Cor. 1. $P \times \frac{W}{2^n - 1} = W$ and $P = \frac{W}{2^n - 1}$; and if any two of these magnitudes P, W , or n , be given, the other may be found.

357. Cor. 2. $m, 2m, 4m, 8m, \&c.$ express the ratio of the parts of W respectively, supported by $G, B, D, F, \&c.$

358. Cor. 3. If the rope, to which P is applied, instead of being fixed to W pass over a pulley to which W is appended, $P : W :: m : 4m :: 1 : 4$; for W is supported by the ropes E, G, F , whose tensions, compared with that to which P is applied, are respectively equal to $m, 2m, m$, and the sum of their tensions $= 4m$.

FIG.
C.

S C H O L I U M.

FIG.
CI.

359. When the directions of the ropes PA , QB , to which two powers Q and P are applied, and the direction in which W acts, are parallel to each other, as is supposed in the preceding propositions, it is evident that W is exactly equal to $P + Q$, because they just support it, and their force is all effective, no part being lost by obliquity of direction. This also appears by considering BCA to be a lever, for (287) $W : P + Q :: CB + CA : AB$, and consequently $P + Q = W$ and $P = Q = \frac{W}{2}$. But if the direction in which P acts be changed to ZA , touching the pulley in D , motion will ensue, parallel to the horizon, as P acts partly in that direction, and the quantity and direction of Q must be changed to restore the equilibrium. It is evident that the pulley cannot be at rest, till the horizontal parts of the forces P and Q be equal to each other, and the parts contributing to the support of W be each equal to $\frac{W}{2}$. This is also obvious from (216); for the pulley BCA is acted upon by three forces, whose directions are not parallel, and is quiescent; and these forces are therefore in the same plane, their directions meet in the same point D , and are consequently equally inclined to the direction of W , or to the horizontal line BA .

FIG.
CII.

360. PROP. If W be supported by two powers P and Q , whose directions touch the pulley in A and B , $W : P$ or $Q :: \sin. \angle$ contained between PA and $QB : \sin. \text{ of half that angle}$.

DEM. Because the pulley is acted upon by three forces P, Q, W , and kept at rest, and their directions are perpendicular to the sides of the triangle BCA , $W : P :: BA : CA$, and $W : Q :: BA : CB$; therefore $W : P$ or $Q :: \sin. \angle BCA$ or its supplement to two right angles $BDA : \sin. \angle BDC$ (equal to the $\angle CBA$ or CAB). Q. E. D.

361. Cor.

FIG. XCIV.

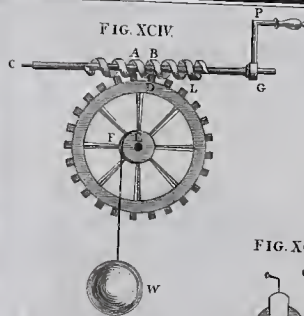


FIG. XCV.



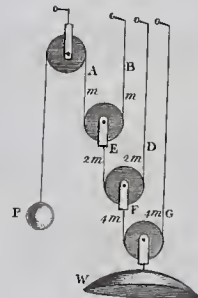
FIG. XCVII.



FIG. XCVI.



FIG. XCVIII.



361. Cor. 1. Because $P : W :: CA : BA$ and $Q : W :: CB : BA$; therefore $P + Q : W :: CA + CB : BA$.

362. Cor. 2. If AB be an arc of 60 degrees, or the angle ADB be equal to 120 degrees, AB is equal to AC or CB , and consequently $W = P$ or Q .

363. Cor. 3. Because $P : W :: AC : AB$, $P = \frac{W \times AC}{AB}$. If therefore the arc AB vanish, or the angle ADB be equal to 180° , $P = \frac{W \times AC}{0}$, or P is infinitely great compared with W . As the arc AB encreases to a semicircle, P decreases and becomes the least possible when it is a semicircle, because the chord AB is then the greatest possible. In this case PA and QB are parallel to each other and to CD , and $W = P + Q$, because $AB = AC + CB$, and consequently $P = \frac{W \times AC}{AB} = \frac{W}{2}$. As the arc BA encreases beyond the semicircle, P encreases and becomes infinite when it is equal to the periphery.

364. Cor. 4. If P be finite, W is either finite or evanescent; for $W = \frac{P \times AB}{AC} = 0$, when AB vanishes, and is finite when AB is finite.

365. PROP. If W be sustained by P in a system of moveable pulleys v, y, x , each of which has a separate rope, and the angles contained by the directions of the ropes be FSE, DTC, BRA ; $W : P :: EF \times DC \times BA : vE \times yC \times xA$.

FIG.
CIII.

DEM. Let W, T, R, P represent the tensions of the strings vW, yT, xR, PG respectively; and (360) $W : T :: EF : vE$

$$T : R :: CD : yC$$

$R : P :: AB : xA$; and consequently

$$W : P :: EF \times CD \times AB : vE \times yC \times xA.$$

Q. E. D.

366. Cor. 1. If the pulleys and the angles FSE, DTC, BRA be equal, $W : P :: EF^3 : vE^3$; and, if the number of moveable pulleys be equal to n , $W : P :: EF^n : vE^n$.

367. Cor. 2. If the directions of the ropes become parallel, $W : P :: 2Ev \times 2yC \times 2xA : Ev \times Cy \times Ax$, and, when the pulleys are equal, $W : P :: 2 \times 2 \times 2 : 1$; or as that power of two, whose exponent is the number of moveable pulleys, to unity, which coincides with (352).

368. Cor. 3. If the tension of the rope PG , to which P is applied, be equal to m , $R = \frac{m \times AB}{Ax}$, $T = \frac{m \times AB \times CD}{Ax \times Cy}$, $W = \frac{m \times AB \times CD \times EF}{Ax \times Cy \times Ev}$.

S C H O L I U M.

369. The conclusions, derived from considering the tensions of the several ropes in any system of pulleys, may also be investigated by supposing the moments of P and W to be equal, the velocities involved in these moments being reduced to opposite directions.

1. In fig. 99. the ropes being all parallel to the direction in which W acts, the velocities of P and W are in opposite directions and entirely efficient; and, when an equilibrium obtains, the product of P and its velocity is equal to the product of W and its velocity: for let $Ll = Mm = Nn = Oo = v$, and if the point L be

be elevated to l , through a space equal to v , the pulley EF will descend through a space equal to v , and any point of the rope E through a space equal to $2v$, and, when the point M is elevated to m , through $3v$; therefore the pulley BA will descend through a space equal to $3v$, and when the ropes A and B are stretched, A must descend through a space equal to $6v$, and when the point N is elevated to n , through $7v$, &c.; consequently the velocities of W and the several strings E, C, A, P , are respectively as $v, 3v, 7v, 15v$, and when an equilibrium obtains $P : W :: 1 : 15$. 2. In fig. 98. let W ascend through a space equal to v , and the ropes F, E, A will evidently be elevated through spaces equal to $2v, 4v, 8v$, respectively, and if the number of moveable pulleys be equal to n , it is clear that the velocity of $P : \text{velocity of } W :: 2^n \times v : v :: 2^n : 1$, and when there is an equilibrium P and W are inversely as their velocities, or as $1 : 2^n$. 3. If the ropes, sustaining the pulleys be not parallel to each other; let EFV be a pulley sustained by two powers E and F , whose directions meet in S ; and if W be raised through a very small space equal to SV , E and F , acting parallel to ES and FS , will describe the spaces SE and SF respectively, and the velocities of W, E, F , are to each other as SV, SE, SF . Resolve SF and SE , each into two, Sm, mF , and Sm, mE , of which mF and mE being equal and opposite destroy each other, and the remainders are wholly efficient; therefore, in an equilibrium, $\overline{E + F} \times Sm = W \times SV$; but $\overline{E + F} \times SF : \overline{E + F} \times Sm$ (or $W \times SV$) :: $SF : Sm : FV : Fm$, and $\frac{\overline{E + F} \times SF}{2} (E \times SE) : W \times SV :: \frac{FV}{2} : Fm :: FV : FE$, or the power or tension at E is to the moment of W as $EV : FE$, which coincides with what is proved in (360).

FIG.
CIV.

INCLINED PLANE.

PLATE
XIII.
FIG.
CV.

370. DEF. *A plane surface inclined, in any angle, to an horizontal plane, is called an inclined plane. If BA be drawn upon the inclined plane, and BC upon the horizontal plane, from the same point B of their common intersection, perpendicular to it, the angle ABC is the angle of elevation or inclination; and if PW be the direction of the power, the angle PWA is called the angle of traction.*

371. PROP. *If a body, whose weight is W, be just supported upon the inclined plane AB by a power P, acting in the direction PW; $P : W :: PW : PC$ (supposing PC to be perpendicular to the horizontal line BC) $:: \sin. \angle \text{inclination } ABC : \cos. \angle AWP$.*

DEM. The weight W , acting in the direction WL or PC perpendicular to the horizon, is supported by the power P , and reaction of the plane, acting respectively in the directions PW , and WC perpendicular to the plane; and consequently P , W , and pressure upon the plane, are to each other respectively as the sides PW , PC and WC of the triangles PWC (196), and $P : W :: PW : PC :: \sin. \angle PCW$ or $ABC : \sin. \angle PWC$ or $\cos. \angle PWA$. Q. E. D.

Another demonstration :

FIG.
XVI.

Let WZ , representing the weight of W , or its tendency to descend in a direction perpendicular to the horizon, be resolved into two, one SZ perpendicular to the plane, and the other WS parallel to it. WS represents the tendency of W to descend upon the plane, and, taking $Wx = WS$, and drawing $x E$ perpendicular to the plane, any force WD , WE , &c. drawn from W and terminated by $x E$, will be in equilibrio with W , because the only efficient

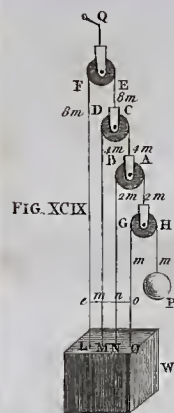


FIG. XCIX

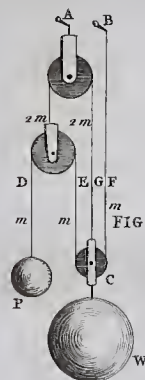


FIG. C

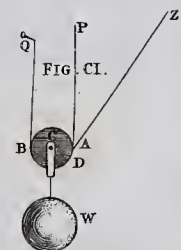


FIG. CI.

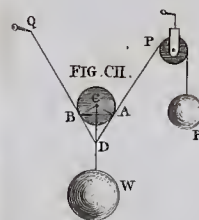


FIG. CII.

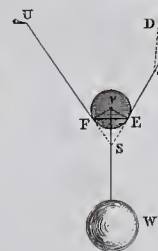


FIG. CIII.

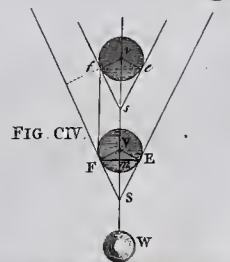


FIG. CIV.

cient part of WE , resulting from resolution, is equal and opposite to WS ; but $P : W :: WE : WZ :: \frac{WX \times \text{rad.}}{\text{cof. } \angle EWX} : \frac{WS \times \text{rad.}}{\text{fin. } \angle SZW} :: \text{fin. } \angle SZW \text{ or } ACB : \text{cof. } \angle EWA$.

372. Cor. 1. The same force, acting in parallel directions, is required to support the same weight upon every part of the plane; because the angle of traction is the same, and $P = \frac{W \times \text{fin. } \angle \text{inclin.}}{\text{cof. } \angle \text{of traction}}$, which is a given quantity.

373. Cor. 2. Because $P = \frac{W \times \text{fin. of } \angle \text{inclin.}}{\text{cof. } \angle \text{traction}}$, if W and the plane be given, P is the least possible when the cosine of the angle AWP is the greatest possible; or when PW is parallel to the plane: as the angle PWA encreases from hence, its cosine decreases, and consequently P encreases, and becomes equal to W , when PW is perpendicular to the horizon, and infinitely greater than W when PW coincides with WG or WC .

FIG.
CV.

374. Cor. 3. When PW is parallel to the plane, $P : W :: \text{fin. } \angle \text{of inclin.} : \text{radius} :: AC : AB$; when it is parallel to the base, $P : W :: \text{fin. } \angle \text{of inclin.} : \text{fin. } \angle BAC$ (this angle being the complement of the angle of traction) $:: AC : BC$.

375. Cor. 4. If AB be perpendicular to the horizon, and the angle $PWA = 0$, $P = W$; for the cof. of $\angle PWA = \text{fin. of } 90^\circ = \text{fin. of } \angle \text{of inclination}$. If, in this supposition, PW be parallel to the horizon, P is infinite; for cof. of $\angle \text{of traction} = 0$.

376. Cor. 5. Let Pr represent the pressure upon the plane, and $Pr : W :: \text{fin. } \angle WPC : \text{fin. } \angle PWC \text{ or cof. } \angle \text{of traction}$, and $Pr = W \times$

$\frac{W \times \sin. \angle WPC}{\cos. \angle PWA} = 0$, when PW is perpendicular to the horizon, and infinite when the $\angle PWA$ is a right angle.

377. Cor. 6. Let CI be perpendicular to P 's direction, and the sides of the triangle CIB , CI , CB , and IB , are respectively perpendicular to the directions of P , W , and pressure upon the plane, and consequently P , W , and $P r$, are respectively as CI , CB and IB .

378. Cor. 7. If WCP were an angular balance revolving upon the fulcrum C , and acted upon by two forces in the directions PW and WL , an equilibrium would obtain between them when $P : W$ inversely as the perpendiculars let fall from C upon their directions, or $P : W :: CL : CK$ (272); but $CL : CK :: \sin. \angle CWL$ or $WCP : \sin. \angle CWK$ or $CWP :: PW : PC$.

FIG.
CVII.

379. Cor. 8. If P and W be in equilibrio, their perpendicular velocities are to each other inversely as their magnitudes; for let W descend through an infinitely small space WA , and its perpendicular descent is WY , and P 's perpendicular ascent is equal to the difference between PA and PW , or Am , if Em be perpendicular to PA . Draw FDn perpendicular to PA , and DZ to FA , and $Am : WY :: P$'s velocity : W 's velocity :: $An : DZ :: AF : DF$ (sim. triangles) :: $W : P$ (377).

FIG.
CVIII.

380. PROP. If two bodies W and V , placed upon the inclined planes AB , AD , support each other by means of a rope passing over the pulley P , $V : W :: \cos. \angle AVP \times \sin. \angle ABC : \cos. \angle PWA \times \sin. \angle ADC$.

DEM. The tension of the rope VPW is every where the same, or the same power, P , is required to support V and W ; but
 $P :$

$P : W :: \sin. \angle ABC : \cos. \angle PWA$; and
 $V : P :: \cos. \angle AVP : \sin. \angle ADC$; therefore
 $V : W :: \sin. \angle ABC \times \cos. \angle AVP : \cos. \angle PWA \times \sin. \angle ADC$.
 Q. E. D.

381. Cor. 1. If PV and PW be parallel to AD and AB respectively, $P : W :: AC : AB$,
 $V : P :: AD : AC$; therefore
 $V : W :: AD : AB$.

This conclusion is also derived from the general expression, the cosines of the angles PVA and PWA becoming AD and AB .

382. Cor. 2. If PV and PW be parallel to the bases of the planes, $P : W :: AC : BC$,
 $V : P :: DC : AC$, and
 $V : W :: DC : BC$.

This also follows from the general expression, the complements of the \angle s AVP and AWP , being, in this supposition, respectively equal to DAC and BAC .

WEDGE.

383. DEF. A wedge is a hard body, generally of a triangular prismatic figure as AF , which is generated by the motion of the triangle AED upon the right line EF always perpendicular to its plane.

 FIG.
CIX.

If AED be an isosceles triangle, it is called an isosceles wedge, if scalene, a scalene wedge. $ABCD$ is the back of the wedge, upon which a force is impressed usually by percussion; and $ABFE$, $DEFC$, are the sides of the wedge, upon which the resistances of wood, &c. act and counterpoise the force of percussion. The angle AED is the vertical angle of the wedge.

384. PROP. If two equal resistances, acting in equal angles upon the sides of an isosceles wedge ABV , be in equilibrio with a power, P , acting

FIG.
CX.

ing in a direction perpendicular to the back; P is to the resistances, as the rectangle of the sines of half the vertical angle of the wedge and the inclination of the resistance to the sides, to the square of the radius.

DEM. Let the directions and quantities of the resistances be CD and Cd ; resolve each into two, DE and de , in the directions of the sides, and CE , Ce , perpendicular to them, and DE , de , are lost by their obliquity. Resolve EC , eC into two forces, EF , eF , parallel to the base, and FC , Fc perpendicular to it; EF , eF , being equal and opposite, destroy each other, and the forces $2FC$, being exactly opposite to, and in equilibrio with P , are equal to it. But $FC:EC::\sin.\angle FEC$ or $\angle BVC:\text{rad.}$

$EC:DC::\sin.\angle EDC:\text{rad.}$; and

ex æquo $FC:DC::\sin.\angle BVC \times \sin.\angle EDC:\text{rad.}^2::2FC$ or $P:2DC$ or resistances. Q. E. D.

385. Cor. 1. If the resistances be perpendicular to the sides, the sine of the angle ADC becomes the radius; and consequently $P:\text{resistances}::\sin.\angle CVA:\text{radius}::AC:AV::AB:AV+BV$.

386. Cor. 2. If the resistances be perpendicular to the axis, $P:\text{resist.}::BC \times CV:BV^2$; and, when the angle AVB is equal to two right angles, $P:\text{resist.}::BC \times 0:BC^2::0:BC$, or P is infinitely less than the resistances.

387. Cor. 3. When the resistances are perpendicular to the back AB , the angle $EDC=BVC$, and $P:\text{resist.}::\sin.^2\angle BVC:\text{rad.}^2::BC^2:BV^2$. If the angle BVA be equal to two right angles $BC=BV$, and $P=\text{resistances}$; and if the angle BVA be diminished without limit, the resistances are encreased without limit.

388. Cor.

388. Cor. 4. If P be the same in the two suppositions of cor. 1. and 2; resist. in cor. 1 : resist. in cor. 2 :: $CV : BV$.

389. Cor. 5. If the resistances be given in cor. 2, P is the greatest possible when the angle BVA is a right angle.

390. Cor. 6. If the quantities of the resistances be given, and their directions be variable; P is the greatest possible in the common supposition of cor. 1.; for it varies as the sine of the angle EDC , which is the greatest when that angle is right.

S C H O L I U M.

391. The relation, subsisting between the power and resistances in the wedge, has been traced through various processes and modes of reasoning, by different philosophical writers*; and different, and frequently contradictory, conclusions have been deduced from different demonstrations, some of which must consequently be erroneous. As in all equilibria of forces, they must either obliquely, or directly, oppose each other, and their intensities, estimated in opposite directions, must be equal; one source of error hath resulted from a misapplication of this equality of moments, which is universally true and will never lead to false conclusions, if the velocities be reduced to opposite directions, and the efficient parts only of the moments be supposed equal. 1. Let two equal resistances R and r act perpendicularly upon the sides of an isosceles wedge ABD , and be in equilibrio with a power P acting perpendicularly upon the base AB ; and taking Dd , the continuation of PD , very small, and drawing dE , de parallel, and DE , De perpendicular, to the sides, the velocities of P , R , r , are to each other as Dd , DE , De respectively; these being described by them in the same time. Resolve DE , De into two, EL and eL parallel

FIG.
CXI.

* Rohault Not. ad a, 9. Maclaurin's Newt. Chap. III. Art. 21. Graves, Lect. I. C. X. Desaguliers, p. 107. Emerson, Prop. 30. Muschenbroek, Sect. CCCCLXIII. and CCCCXCI.

rallel to the base, which being equal and opposite destroy each other, and DL , DL perpendicular to the base, which being opposite to P 's direction, are wholly efficient; and because there is an equilibrium $\overline{R+r} \times DL = P \times Dd$, but $\overline{R+r} \times DL$ or $P \times Dd : \overline{R+r} \times DE :: DL : DE :: AP : AD :: AB : AB + AD$, the same as (385).

When R and r act in any other direction, the relation of their moments to that of P is discoverable by a similar process. If they be parallel to the base, draw Mdm parallel to their direction, and the velocities of P , R and r , are respectively as Dd , DM and dm ; and resolving DM , dm into two DE , de perpendicular, and ME , me parallel, to the sides, the last not being opposed to the sides are quite inefficient; and repeating the process above $\overline{R+r} \times DL = P \times Dd$; and $P \times Dd : \overline{R+r} \times DE :: AP : AD$

$$\overline{R+r} \times DE : \overline{R+r} \times DM :: DE : DM :: PD : AD$$

and $P \times Dd : \overline{R+r} \times DM :: AP \times PD : AD^2$; which is the same as (386).

A different analogy is investigated by Rowning, from the same principle, viz. the power : the resistances :: $AP : PD$.

2. Another source of error is an improper use of this principle, "that if three forces act upon a body, which remains at rest, they are to each other as the three sides of a triangle, respectively parallel or perpendicular to their directions;" which certainly is untrue, if any parts of them be inefficient and lost by obliquity of direction; for the force FG will have exactly the same effect upon the sides of the wedge AD with FH , FI , or the perpendicular FL , supposing GL to be parallel to AD , and of the parts, resulting from a resolution of each into two forces, one perpendicular to, and the other coincident with, the sides, these last, Fx , &c. are lost. When the wedge does not fill the cleft, and the resistances act perpendicularly upon its sides, Maclaurin and Varignon have applied the above principle and deduced different conclusions, both of which cannot be true. Dr. Hamilton, Gravesands, Desaguliers, Emerson, Muschenbroek assert, that the power is to the resistances as half the back is to the height, and Keil, Whiston and

and Nicholson, assign this ratio to be that of the whole back to the height. The first analogy coincides with this general theorem, "that the power is to the resistances, when in equilibrio, as a line drawn from the bisection of the base to one side, parallel to the resistance upon that side, to the height of the wedge; and this theorem is presumed to be erroneous for the following reasons.

1. If the directions of the resistances be parallel to the axis, their sum is equal to the power, according to the theorem, which cannot be true; because the power acting perpendicularly is wholly efficient, but part of the resistance is lost by obliquity of direction.

2. If the vertical angle of the wedge be diminished without limit, a line drawn, from the bisection of the base to a side, parallel to the resistance, accedes to equality with the height when the direction of the resistance is parallel to the height, and the power and resistance become equal, according to the theorem; but, the power continuing the same, the directions of the resistances accede to a parallelism with the sides, and they become wholly inefficient.

3. When the resistances act perpendicularly to the axis, and the vertical angle of the wedge accedes to two right angles; according to the theorem, the ratio of the power to the resistances encreases without limit, which cannot be true, because the inefficiency of the resistances evidently encreases without limit, and consequently the ratio of the resistances to the power encreases without limit, which is directly the reverse of the theorem.

4. If the power be the same in the two cases where the directions of the resistances are perpendicular to the sides and axis, of the wedge, the resistances are, according to the theorem, as the side and height of the wedge respectively; or greater resistances are required to sustain a given power, when their directions are perpendicular to the sides, and therefore entirely efficient, than when they are oblique and consequently not wholly efficient, which certainly is untrue.

It is presumed, for these reasons, that the analogy, derived from the theorem, and consequently the theorem itself, cannot be generally true; and were these reasons less decisive, it would not be difficult to point out the several oversights in the demonstration which induced a false conclusion*.

FIG.
CXIII.

392. PROP. *If the power and resistances P , R , and r , act perpendicularly upon the back and sides of a scalene wedge ABD , and be in equilibrio, $P : R + r :: \text{back of the wedge} : \text{sum of its sides}$.*

DEM.

FIG.
CXII.

*THEOREM. If a power P , acting perpendicularly upon the back of an isosceles wedge ABC , be in equilibrio with the equal resistances E and F , acting perpendicularly upon its sides, $P : E + F$ as a line PE drawn from the bisection of the base P to one side, parallel to the resistance upon it, to the height of the wedge PB .

DEM. Let the directions of E and F meet in the bisection of the base P , and $PE = PF$, and, completing the parallelogram, $P : E + F :: PN : PE + PF$ (196) :: $PH : PE :: PE : PB$.

Cor. If PN be the force of P , PE and PF will represent the force with which P protrudes the resistances in directions perpendicular to the sides.

THEOREM II. If the directions of the resistances be any other lines PD , PO , equally inclined to the sides, the power is to the resistances ($E + F$) as PD to PB .

DEM. Resolve PE and PF into two, PG and PK in the directions of PD and PO , and GE , FK perpendicular to them; and these last forces, acting perpendicularly to the resistances, are lost. If PN be the force of P , PE and PF are its force perpendicular to the sides, and $PG + PK$ the force in equilibrio with the resistances; therefore $P : E + F :: PN : PG + PK :: PH : PG :: PD : PB$, because $PE^2 = PG \times PD = PH \times PB$, and consequently $PH : PG :: PD : PB$. Q. E. D.

In this demonstration, EG and FK , being inclined to the sides of the wedge, are not inefficient, and, not being opposite, they do not destroy each other, and consequently ought not to be neglected. And, besides, if the power estimated in the directions of PD and PO be equal to $PG + PK$, these will be equal to the power estimated in the direction PN , or equal to PN . Resolve PG and PK into two, GL , OL , perpendicular to the axis, which being equal and opposite, destroy each other, and PL and PL are the only remaining parts of the force, and these are never equal to PN , unless PD and PE coincide. *Hamilton on the Principles of Mechanics.*

DEM. Because the power and resistances act perpendicularly upon the back and sides, they are wholly efficient; their directions are in the same plane, meet in the same point (216); and their magnitudes are to each other as the three sides of a triangle parallel, or perpendicular, to their directions; therefore $P : R + r :: AB : AD + DB$. Q. E. D.

393. Cor. If PM, RE, re , be the relative magnitudes of P, R, r , in equilibrio, and lines be drawn through M, E , and e , parallel to the sides of the wedge respectively, any forces, whose magnitudes are PL, RF, rf , drawn from P, R and r , terminated by the lines drawn parallel to the sides, will be in equilibrio; for their perpendicular and only efficient parts PM, RE and re are in equilibrio.

S C R E W.

*394. DEF. If a right line, divided into equal parts AC, CE, EG , &c. representing equal inclined planes, be so wrapped round the convex and concave surfaces of two cylinders with equal bases, that AB, CD, EF , &c. the horizontal bases of the planes, may be bent into the peripheries of circles parallel and equal to the bases of the cylinders; equal spirals will be formed upon their surfaces, whose lengths are AC, CE, EG , &c. and distances CB, ED, FG , &c. the perpendicular heights of the planes. The convex cylinder is inserted in the concave, and so adapted to it that the spirals protuberant upon one cylinder may exactly fill the excavated spiral or groove upon the other: the first is called the external, and the second the internal, screw.

PLATE
XIV.
FIG.
CXIV.

Another definition:

If a cylinder move uniformly about its axis, whilst a point moves uniformly upon its surface in a right line parallel to the axis, the line, described by this compound motion, is a spiral, which, being raised upon the external surface of the cylinder, forms the external screw; and a similar spiral groove being cut upon the internal surface of a hollow cylinder;

* Maclaurin, Chap. III. XXII. Keil's Physics, Lect. X. Rohault, Not. ad 9. Emerson's Mechanics, Prop. 29.

linder, of the same diameter, to receive the protuberant spiral, forms the internal screw. Either of these screws is fixed, and the other is moveable by a lever passing through the center, and in the plane, of its base.

FIG.
CXV.

395. PROP. *There is an equilibrium upon the screw, when the power, P, is to the weight, W, or resistances acting parallel to the axis, as the distance between two contiguous spirals, to the periphery of the circle described by the power.*

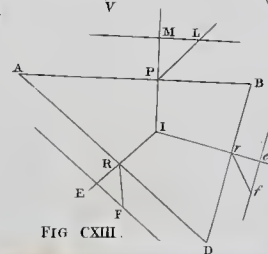
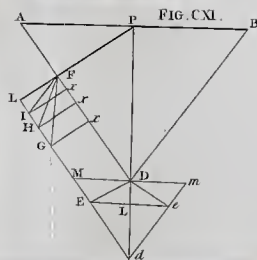
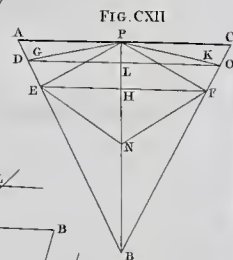
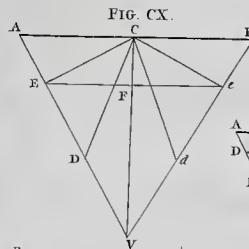
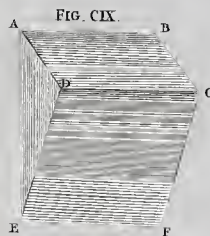
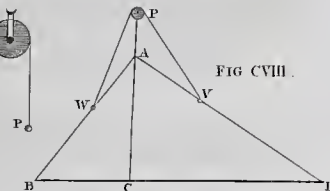
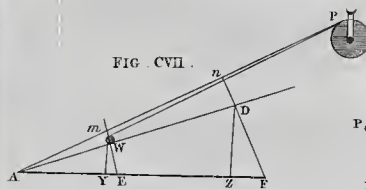
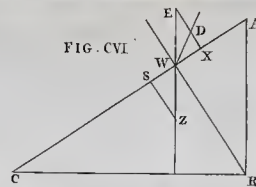
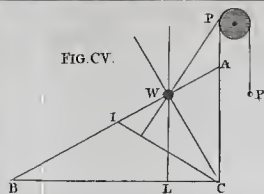
DEM. It is evident that W , acting upon either of the screws, in a direction parallel to its axis, will equally press every point of the spirals of the other in contact with it, in directions perpendicular to their bases. If ABD be the base of one spiral, p , the magnitude of a power acting at B perpendicular to BC , and in equilibrio with the pressure pr , upon one point of the spiral, P' a power acting at P perpendicularly to PC , and in equilibrio with p or pr ; (374) $p : pr :: \text{distance between two contiguous spirals } (d) : ABD$; $P' : p :: BC : PC :: ABD : \text{periphery of the circle described by the power, therefore } P' : pr :: d : \text{periphery of this circle, and the sum of all the } P'\text{'s, or } P \text{ is to the sum of all the } pr\text{'s, or } W \text{ in the same ratio, that is, } P : W :: d : \text{periphery of the circle described by } P.$ Q. E. D.

Another demonstration from def. 2.

Whilst B or P makes one revolution, W is elevated through a space equal to d , therefore the velocity of P is to the velocity of W , as the periphery of the circle described by P to the distance between two contiguous spirals, and consequently when there is an equilibrium, $P : W :: \text{dist. between two spirals} : \text{the periphery described by } P.$ Q. E. D.

396. Cor. 1. If the direction of W be inclined to the axis, and P do not act in the plane of the base, but in any other direction, the ratio of P to W may be found by art. 371.

397. Cor.



397. Cor. 2. The pressure sustained by any point of the spirals is to the part of W incumbent upon it, as the length of a spiral to its base ABD .

398. Cor. 3. In an endless or perpetual screw, acting upon a wheel, the distance of whose teeth is equal to AB , the distance of two contiguous spirals, and in equilibrio with a body W , suspended from the axis EF ; $P : W :: AB \times EF : \text{diameter of the wheel} \times \text{periphery of a circle whose radius is } PG$; for if R be the resistance of a tooth to a spiral of the screw, $P : R :: AB : \text{per. of the circle described by } P$, $R : W :: EF : \text{diameter of the wheel}$; therefore $P : W :: AB \times EF : \text{per. of the circle described by } P \times \text{diameter of the wheel}$.

PLATE
XII.
FIG.
XCIV.

S C H O L I U M.

399. The second demonstration of this proposition is generally deemed unsatisfactory, because the moments of P and W are not reduced to opposite directions; but the ratio of P to W may be investigated by finding the opposite and efficient parts of a power sustaining a body upon an inclined plane, and combining it with the power of the lever. Let WC be an inclined plane, or spiral of the screw, W a weight acting perpendicularly to the horizon WE , and supported by a power P , acting parallel to WB ; and if W be elevated through a very small space Ww , or perpendicular height Hw , it is evident that P , acting parallel to WE , will descend through a space equal to $WE - we$ or WH ; and consequently the velocities of W and P , estimated in opposite directions, are as Hw to HW , or as the height of the plane to its base, and when an equilibrium obtains, $P : W :: \text{height of the plane} : \text{its base or } ABD$. But, from the nature of the lever, a power, acting at the distance PC , must be diminished in the ratio of $BC : PC$, or of the circumference $ABD : \text{circumference described by the power acting at } P$.

FIG.
CXVI.

FIG.
CXV.

CHAP. IX.

CENTER OF GRAVITY.

400. DEF. **T**HE center of gravity of a body or system of bodies is a point about which the parts of the body or system are in equilibrio.

401. PROP. To find the center of gravity of a body.

PLATE
XIV.
FIG.
CXVII.

Let $A, B, C, D, \&c.$ be particles of the body, and finding the centers of equilibrium p and q , of A and B, C and D respectively (274); let $A + B$ be placed in p , and $C + D$ in q , and their center of equilibrium, G , will be the center of gravity of the particles $A, B, C, D, \&c.$ Because the force of gravity acts upon the particles in parallel directions, the efficacy of A to communicate motion to G is $A \times AG$, and that of B is $B \times BG$ (279) or $A \times \overline{Ap + pG}$ and $B \times \overline{Bp + pG}$, which are equivalent to them (181), or $\overline{A + B} \times pG$, since $A \times \overline{Ap}$ and $B \times \overline{Bp}$ are equal and opposite, and consequently destroy each other. The sum of the moments of C and D is found, by a similar process, to be the same as if they were placed in q ; and consequently G , which is the center of gravity of $A + B$ and $C + D$, placed in p and q respectively, is the center of gravity of A, B, C, D , placed at the points $A, B, C, D, \&c.$ Q. E. D.

402. Cor. 1. The particles of the body cannot be in equilibrio about any other point except G ; for, if possible, let X be such a point, and it is proved, as before, that the efforts of A and B to move $X = \overline{A + B} \times pX$, and of C and $D = \overline{C + D} \times qX$; therefore the point X is kept in equilibrio by two forces, $\overline{A + B} \times pX$ and

and $\overline{C + D} \times qX$, not acting in opposite directions, which is impossible (187).

403. Cor. 2. In every situation of the body composed of the particles A, B, C, D , &c. if the point G be supported, the body will be at rest; for the force of gravity acting always in parallel directions upon the particles, their moments, or efforts to move G , will always be as $A \times AG, B \times BG$, &c. which by the process used in this proposition, will always be reduced to two forces that are equal and opposite.

404. Cor. 3. If $A + B + C + D$, &c. be equal to \mathcal{Q} , and the pressure of each in parallel directions be equal to q , a force as $\mathcal{Q} \times q$, acting at the point G , in a direction opposite to that in which the particles press, will remove their pressure. Or if A, B, C , &c. be destitute of gravity, and only resist the action of a force by their inertia, a force P acting at G will communicate equal velocities to every particle; because their resistances, being exerted in directions opposite to that of P (3d law of motion) and therefore parallel to each other, vary as their distance from G , and consequently the sums of the resistances on each side of G are equal. And $v.v$, if \mathcal{Q} be moving and without gravity, a force applied, at G (the center of inertia) equal to the moment of \mathcal{Q} , will destroy all motion.

S C H O L I U M.

405. The particles which compose a body, being connected together by the force of cohesion, every line of particles may be considered as a lever, impressed in different points by the action of gravity, or any force which acts in parallel directions. The particle A therefore is connected with G by the cohesion of the intermediate particles, and, an infinite number of levers terminating in G will therefore be formed, upon which the particles, in any one line, act with intensities varying as their magnitudes multiplied into their distances.

FIG.
CXVIII.

406. PROP. *The sum, or difference, of the products, which results from multiplying each particle A, B, C, D into its perpendicular distance from any plane LN, according as they are on the same or different sides of the plane, is equal to the product of all the particles multiplied into the distance of their center of gravity, G, from that plane.*

DEM. Let P and Q be the centers of gravity of A and B , C and D , and drawing right lines through P , Q , G , parallel to the plane, which intersect the perpendiculars drawn from those points respectively; and $A : B :: BP : AB :: Bn$ or $Bb - Pp : Am$ or $Pp - Aa$,

and $A \times Pp - Aa = B \times Bb - Pp$, or

$$A \times Aa + B \times Bb = A + B \times Pp.$$

By a similar process it appears, that $C \times Cc + D \times Dd = C + D \times Qq$. But $A + B : C + D :: QG : PG :: Gv$ or $Gg \mp Qq : Pp$ or $Pp - Gg$, and $A + B \times Pp - Gg = C + D \times Gg \pm Qq$, or, by transposition and a substitution of equals, $A \times Aa + B \times Bb \pm C \times Cc \pm D \times Dd = A + B + C + D \times Gg$, where the higher or lower signs are to be used, according as the bodies are on the same, or a different, side of the plane. Q. E. D.

FIG.
CXIX.

407. Cor. 1. If the particles be placed upon the same right line, or Aa, Bb, Cc, Dd, Gg , become Ag, Bg, Cg, Dg, Gg , respectively, it is evident, that $A \times Ag + B \times Bg \pm C \times Cg \pm D \times Dg = A + B + C + D \times Gg$; or the sum, or difference, of the products resulting from the multiplication of each particle into its distance from any point g , according as they are on the same, or a different, side of that point, is equal to the product of their sum multiplied into the distance of their center of gravity from that point.

408. Cor. 2. The whole moment of a body, acting upon a lever, being equal to that of every particle, or to the sum of the products which results from the multiplication of each particle into its distance from the center of motion (279), is equal therefore to the

the product of the whole body into the distance of the center of gravity from the center of motion, and is consequently the same as if it were collected in the center of gravity. The demonstration of this proposition obtains therefore when A, B, C, D , are collections of particles or bodies, whose centers of gravity are the points A, B, C, D . And to find the center of gravity of a system of bodies, it is evident, that, in the proposition (401), bodies, whose centers of gravity are A, B, C, D , &c. may be substituted for particles.

409. Cor. 3. If A, B, C, D , be bodies acting upon any plane LN , in parallel directions, the sum of their efforts to move it is the same as if they were collected in their center of gravity; for, if A, B, C, D , be the respective centers of gravity of each body, this sum is equal to $A \times Aa \times Bb \pm C \times Cc \pm D \times Dd = \overline{A + B + C + D} \times Gg$; or, if they be placed upon a lever, the sum of their efforts to make it revolve is the same, as if they were placed at G . When the center of gravity therefore is in the plane, or at the fulcrum of a lever, the plane and lever are quiescent. And if any point Z be taken in NL , $A \times aZ + B \times bZ + C \times cZ + D \times dZ = \overline{A + B + C + D} \times gZ$; for if a plane pass through Z , the proof is the same as that of this proposition.

 FIG.
CXVIII.

410. Cor. 4. The distance of any plane from the common center of gravity of A, B, C, D , &c. or Gg is equal to $\frac{A \times Aa + B \times Bb \pm C \times Cc \pm D \times Dd}{A + B + C + D, \&c.}$; and its distance from a plane passing through any point Z is equal to $\frac{A \times Za + B \times Zb \pm C \times Zc \pm D \times Zd}{A + B + C + D}$, where the lower signs are to be used for those bodies not on the same side of Z with A and B .

FIG.
CXX.

411. Cor. 5. A right line drawn from A through the center of gravity G of any number of bodies $A, B, C, D, \&c.$ will pass through the center of gravity of the remainder; for $B \times Bb + D \times Dd = C \times Cc$, and consequently the center of gravity of B, C, D is in the plane passing through AG , and if this plane revolve, their center of gravity is always in the plane passing through AG , and consequently it must be in the line AG produced, which is the common intersection of the planes. If r be this center, $\overline{B + C + D} \times Gr = A \times AG$, and if the bodies be equal and n their number, $AG = n - 1 \times Gr$.

FIG.
CXXI.

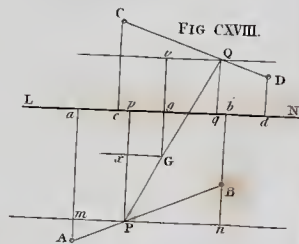
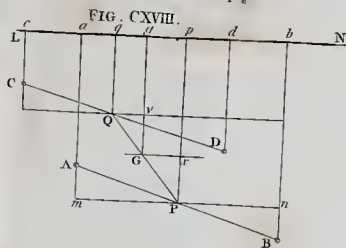
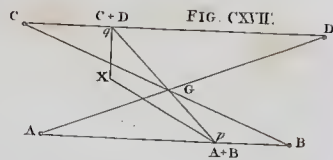
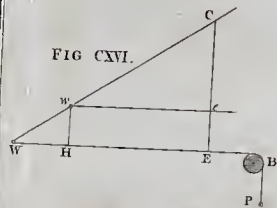
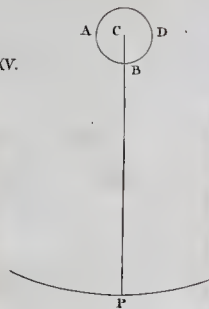
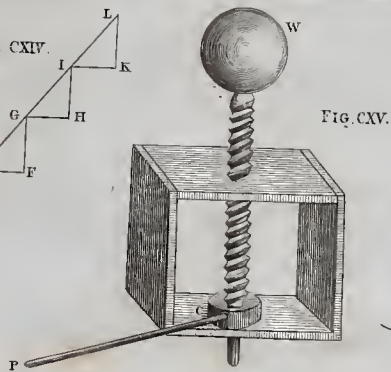
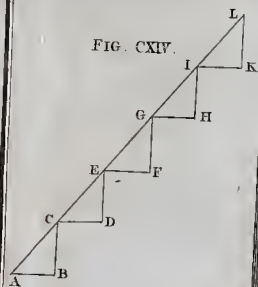
412. Cor. 6. If a circle or sphere be described about the center of gravity G , of any number of bodies $A, B, C, \&c.$ and any point P be taken in the periphery of the circle, or surface of the sphere, $PA^2 \times A + PB^2 \times B + PC^2 \times C, \&c.$ is a given quantity; for, drawing GP and the perpendiculars to it Aa, Bb, Cc , $A \times Ga = B \times Gb + C \times Gc$ (409), or, by substitution of equals, $A \times \frac{GA^2 - PA^2 + GP^2}{2GP} = B \times \frac{PB^2 - BG^2 - GP^2}{2GP} + C \times \frac{PC^2 - GP^2 - GC^2}{2GP}$, or, $A \times PA^2 + B \times PB^2 + C \times PC^2 = A \times GA^2 + GP^2 + B \times GB^2 + GP^2 + C \times GC^2 + GP^2$, and this side of the equation is invariable in whatever point of the periphery or surface P be placed.

FIG.
CXXII.

413. PROP. If $A, B, C, D, \&c.$ be particles of a body urged by forces in parallel directions, whose magnitudes are $Aa, Bb, Cc, \&c.$ the sum of their weights is equal to the weight of $A + B + C, \&c.$ acted upon by a force whose magnitude is Gg .

DEM. The weights of $A, B, C, \&c.$ are $A \times Aa, B \times Bb, C \times Cc, \&c.$ (235), and consequently the sum of their weights is equal to $A + B + C, \&c. \times Gg$; but this product is the weight of $A + B + C, \&c.$ acted upon by the force Gg . Q. E. D.

414. Cor.





414. Cor. If the forces $Aa, Bb, Cc, \&c.$ be equal to each other, Gg is equal to one of them, or if the particles $A, B, C, \&c.$ be acted upon by the same force, their weight is the same as if they were collected in their center of gravity and acted upon by that force. The tendency therefore of a body to descend is the same as if it were collected in its center of gravity, and, consequently, if a line drawn from that center perpendicular to the horizon, fall within the base of the body, it cannot fall, and, if without the base, it cannot stand.

415. PROP. *If any number of bodies $A, B, C, \&c.$ move in parallel directions, with any velocities, the center of gravity will describe a right line parallel to them.* FIG. CXXIII.

DEM. Let A and B, a and $b,$ be cotemporary positions of the bodies A and $B,$ and $G, g,$ their centers of gravity, and through g draw a line xy parallel, and consequently equal, to $AB.$ From the nature of the center of gravity, $A : B :: BG : AG :: bg : ag :: yg : xg$ (sim. triangles); and the point g divides the parallel and equal lines $AB, xy,$ in the same ratio, and Gg is a right line parallel to Aa or $Bb.$ If H be the center of gravity of $A, B, C,$ it is proved in the same manner, that it cuts the parallel and equal lines $GC, vz,$ in the same ratio, and Hb is consequently a right line parallel to $Gg.$ Q. E. D.

416. Cor. 1. If any number of bodies $A, B, C, \&c.$ ascend, or descend in parallel right lines, the sum of the products resulting from the multiplication of each body into the space described by it, is equal to the product of their sum and the space described by their center of gravity $G;$ for, let $a, b, c, g,$ be cotemporary positions of $A, B, C, G,$ and, drawing any plane $NL, A \times Am + B \times Bn + C \times Cq = \overline{A + B + C} \times Gb$ (406), and $A \times am + B \times bv + C \times cq = \overline{A + B + C} \times gb;$ and consequently by addition $A \times Aa + B \times Bb + C \times Cc = \overline{A + B + C} \times Gg.$ FIG. CXXII.

417. Cor.

FIG.
CXXIII.
CXXIV.

417. Cor. 2. If any number of bodies therefore move in parallel directions with any unequal velocities, or they be placed upon the lever XY , and receive unequal impulses from any force at the same time in parallel directions, the center of gravity will, in the beginning of its motion, move uniformly in a right line parallel to them, and its velocity is equal to the products of each body into its velocity, divided by the sum of the bodies; for the spaces Aa , Bb , Cc , Gg are described in the same time, and vary as the velocities, and Gg (or velocity of G) = $\frac{A \times Aa + B \times Bb + C \times Cc}{A + B + C}$.

FIG.
CXXV.

418. PROP. *The surface, or solid, described by the line or surface AB , moving round an axis passing through C , is equal to the surface or solid formed by the multiplication of AB into the line described by the center of gravity, G , of AB .*

DEM. For if AB be covered with physical points of the same thickness, AD , DE , EF , &c. $AD \times CA + DE \times CD + EF \times CE$, &c. = $\overline{AD + DE + EF}$, &c. or $AB \times CG$ (407), or because the arcs Aa , Dd , Ee , &c. are similar, $AD \times Aa + DE \times Dd + EF \times Ee$, &c. (= the sum of all the superficial, or solid, annuli Ad , De , &c.) = $AB \times Gg$.

FIG.
CXXVI.

419. PROP. *The content of the solid generated by the revolution of a plane figure ABC round an axis AC , is equal to that of the solid whose base is ABC , and perpendicular height the periphery of the circle described by the center of gravity, G .*

DEM. For the moment of ABC or $ABC \times GD$, is equal to the sum of the moments of the elementary parts of ABC , or to the sum of all the small rectangles En multiplied into the distance of their center of gravity from AC , or to the sum of all the $y \times mn \times \frac{y}{2}$ or $\frac{y^2}{2} \times mn$; therefore $ABC \times GD$ is equal to the sum of
all

all the $\frac{y^2}{2} \times mn$, and $p \times ABC \times GD$ to the sum of all the $\frac{p y^2}{2} \times mn$; but if p be the periphery of a circle whose radius is unity, $p \times GD$ is the periphery of a circle whose radius is GD ; $\frac{p \times y^2}{2}$ is the area of a circle whose radius is y , and the sum of all the $\frac{p y^2}{2} \times mn$, is therefore equal to the solid described by ABC .

420. DEF. *The center of gravity of a line or surface, is the center of gravity of equal physical points, which are supposed to compose that line or surface.*

421. Cor. 1. The center of gravity of a right line is evidently in its bisection, and the center of gravity of a parallelogram, prism, cylinder, are in the bisection of the diameter or axis, because the number of physical lines in a parallelogram, or of physical laminæ in the prism and cylinder, on each side of this point, are equal.

422. Cor. 2. The center of gravity of any number of lines AB , CD , EF , &c. is found by joining the centers of gravity m , n , of any two AB , CD , and taking a point p , so that $pm : pn :: CD : AB$, and joining p , and q the center of gravity of EF , and taking a point r so that $pr : rq :: EF : AB + CD$, &c. (401). The center of gravity of the sides of a triangle ABC may be found as above, or it may be found by bisecting the sides in m , n , p , joining the points of bisection, and bisecting the angles mpn , and nmp , by the lines pr and mq , whose intersection, G , is the center of gravity; for $pq : qn :: pm : mn :: AC : BC$, or the center of gravity is in qm ; and $mr : rn :: mp : pn :: AC : AB$, or this center is in pr , and therefore must be in their intersection G .

FIG. CXXVII.

FIG. CXXVIII.

FIG.
CXXIX.

423. PROP. *If a regular polygon ABCDE be inscribed in the segment of a circle whose center is F, half the sum of its sides is to the sine of half the arc :: perpendicular let fall from F upon a side : the distance of G from F.*

DEM. For let m, n, p, q , be respectively the centers of gravity of the sides, and joining vs , the centers of AB, BC , and CD, DE , it is clear that the center of gravity of all the sides will be in vs , and it will be also in FC , and consequently at G the intersection of FC and vs . But from the similar triangles $sFq, EID, ED : EI (:: ED + DC : EC) :: Fq : Fs$, and from the similar triangles $GsF, CEH, CE : HE :: Fs : FG$, and, by adding these analogies together, $ED + DC : EH :: Fq : FG$. Q. E. D.

424. Cor. 1. If a regular polygon be inscribed in a circle, its center of gravity is the center of the circle; for $FG = \frac{\text{fine of half the periphery} \times Fq}{\frac{1}{2} \text{ the sum of the sides}} = \frac{0 \times Fq}{\frac{1}{2} \text{ the sum of the sides}}.$

425. Cor. 2. As the arc $EC : \text{fine of that arc} :: FE : \text{distance of the center of gravity of twice } EC \text{ from the center } F$; for, let the sides of the polygon inscribed, be diminished without limit, so as to coincide with the circular arc and be equal to it, and Fq is then equal to the radius. If the arc be a semicircle, the quadrant $CK : FC :: FC : FG$; and the center of gravity of the whole periphery is in F , because $FG = \frac{0 \times FC}{KCL} = 0.$

426. PROP. *The distance of the center of gravity of a triangle from any of its angles, is equal to two third parts of the line drawn from that angle to the bisection of the side opposite to it.*

DEM.

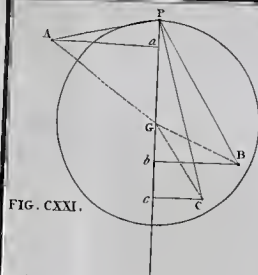


FIG. CXXI.

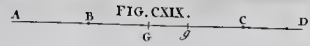


FIG. CXXII.

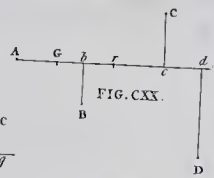


FIG. CXXIII.

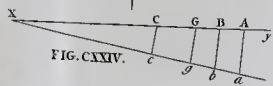


FIG. CXXIV.

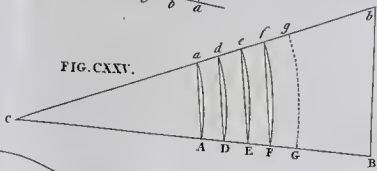


FIG. CXXV.

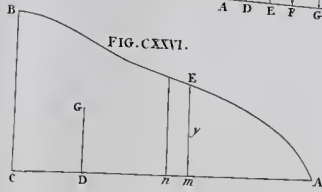


FIG. CXXVI.

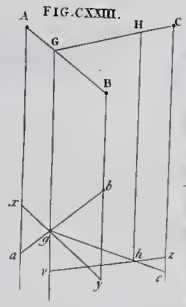


FIG. CXXVII.

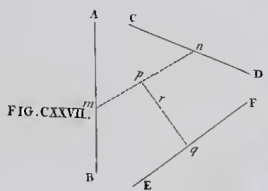


FIG. CXXVIII.

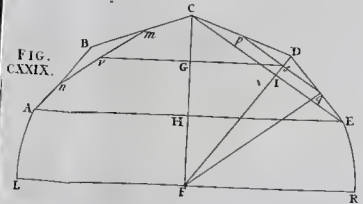


FIG. CXXIX.

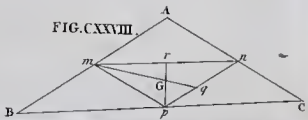


FIG. CXXX.

DEM. Let Ap bisect the opposite side BC , and drawing $ab, cd, ef, \&c.$ infinitely near and parallel to BC , each will be bisected, and the center of gravity of each, and consequently of all, which make up the triangular surface, will be in Ap ; and, for the same reason, it is in Bq bisecting AC , and therefore must be in their intersection G ; but, from similar triangles, $AG : Gp :: AB : pq :: BC : pC :: 2 : 1$. Q. E. D.

FIG.
CXXX.

427. Cor. 1. The center of gravity of any right lined figure is found by dividing it into triangles, and considering them as bodies: let m be the center of gravity of ABC , and n of BDC and joining mn , and taking a point G , so that $mG : nG :: DBC : BAC$, it will be the center of gravity of the trapezium $ABDC$ (401). The process is similar when the figure has more sides.

FIG.
CXXXI.

428. Cor. 2. The center of gravity of a regular polygon, inscribed in a circle, is the center of that circle; for drawing a right line BH , from any of the angles through the center, the figures on each side are similar and equal, and the center of gravity is in BH ; and, for the same reason, it is in DI , and therefore must be at C , and consequently the center of gravity of a circle is its center.

FIG.
CXXXII.

429. PROP. The distance of the center of gravity of a regular polygon, inscribed in a sector of a circle $FEAC$, from its center, is equal to two thirds of distance of the center of gravity of the periphery $EDABC$.

FIG.
CXXXIII.

DEM. Find the center of gravity, H , of the periphery (423), and, taking $IE = \frac{1}{3}FE$, describe a polygon similar to $EDABC$. The centers of gravity of the triangles FED, FDA, FAB, FBC , are at L, S, T, V , the bisections of IK, KN, NP, PQ ; and the centers of gravity of $FEDA$ and $FABC$, are at R , and X , the bisections of SL, VT , and the center of the whole figure is at G , the bisection of RX ; but $FE : EI :: FM : LM$, therefore $LM = \frac{1}{3}FM$, and, for the same reason, $RO = \frac{1}{3}FO$ and $GH = \frac{1}{3}FH$. Q. E. D.

Y

430. Cor.

430. Cor. 1. The distance of the center of gravity of a sector of a circle from the center, is therefore equal to two thirds of the distance of the center of gravity of the circular arc upon which it stands.

431. Cor. 2. As the arc AE : sine of AE :: radius : distance of the center of gravity of twice AE from F , or FH (425); but $FG = \frac{2}{3} \times FH$; therefore AE : sine of AE ($:: EABC$: its subtense EC) $:: \frac{2}{3}$ radius : FG .

S C H O L I U M.

432. The distance, from the vertex, of the center of gravity of a parabolic space is equal to three fifths, and of a paraboloid to two thirds, of the axis; but these, and the most difficult propositions upon the center of gravity, are best investigated by a fluxional process.

FIG.
CXXXIV.

433. PROP. *If one or more of the bodies $A, B, C, \&c.$ move uniformly in the same right line, with velocities equal to $a, b, c, \&c.$ their common center of gravity will move uniformly.*

DEM. Let A and B move uniformly in the same, or an opposite, direction, P be their center of gravity, and D their distance: and because the motions of A and B are uniform, D either continues the same, or encreases and decreases uniformly; but $AP = \frac{D \times B}{A+B}$ and consequently varies as D , and P moves uniformly. If another body C move uniformly in the same right line, and R be the center of gravity of A, B, C ; the distance CP either continues the same, or encreases and decreases uniformly, because C and P move uniformly; but $PR = \frac{CP \times C}{A+B+C}$, and consequently varies as CP , or R moves uniformly. Q. E. D.

434. Cor.

434. Cor. 1. The velocity of the center of gravity is equal to $\frac{Aa \pm Bb \pm Cc}{A + B + C}$; for, let p, r, a, b, c , be cotemporary positions of P, R , and the bodies, and (407) $A \times Ap$ or $\times \overline{Aa + ap} - B \times Bp$ or $\times \overline{bp \pm Bb} = \overline{A + B} \times Pp$, and $Pp = \frac{A \times Aa \pm B \times Bb}{A + B}$, because $A \times ap - B \times bp = 0$. And, placing $A + B$ in P and repeating the above process, it appears that $Rr =$ the velocity of R (106) $= \frac{A \times Aa \pm B \times Bb \pm C \times Cc}{A + B + C}$. It is from hence again collected, that the velocity of R is uniform; because Aa, Bb, Cc are constant, and consequently their sum, or difference, multiplied into the same given quantities, or the velocity of R , is always the same.

435. Cor. 2. Because $\overline{A + B + C} \times Rr = A \times Aa \pm B \times Bb \pm C \times Cc$, the velocity of the center of gravity is such as would be communicated to the sum of the bodies acted upon by a force equal to $A \times Aa \pm B \times Bb \pm G \times Cc$.

436. *PROP. If one or more bodies $A, B, C, \&c.$ move uniformly in right lines, either in the same or different planes, their common center of gravity S will move uniformly in a right line.

FIG.
CXXXV.

DEM.

* This proposition is proved very neatly, though, perhaps, less intelligibly by the following process.

CASE I. Let two bodies move, in the same plane, in the directions DE, AB ; and let D and A, E and B be cotemporary positions, and H, K , the centers of gravity in those positions, respectively; and taking $BP = AD$, joining EP and drawing DL parallel to HK , $DE : AB$ in the given ratio of the motions of the bodies; and, because the angle EDP is given, all the angles of the triangle EDP are given, and DP is to PE in a given ratio: but $PE : PL :: BE : BK$, which is a given ratio, therefore $DP : PL$ in a given ratio; and, because all the angles of the triangle DPL are given, the angle PDL is given, and L is always in DL . By the nature of the center of gravity, $DA : DH :: EB : EK :: PB$ or $DA : LK$; therefore $DH = LK$, and $DHKL$ is a parallelogram, HK is parallel to DL , and the angle BHK is given, and the center of gravity K is always in the right line HK given in position. And, because all the angles of the triangles DPL and DLE are given, the lines $DP, DE,$
 $DL,$

FIG.
CXXXVI.

C E N T E R O F G R A V I T Y.

DEM. I. Let B describe Bb uniformly in the time T , and P, Q be the centers of gravity of A and B ; and $A + B : B :: AB : AP :: Ab : AQ :: Bb : PQ$ (EUC. B. 6. P. 5.), and PQ is parallel to Bb , and equal to $\frac{Bb \times B}{A + B}$, and varies therefore as Bb or uniformly.

Let A describe Aa uniformly in the time T , either in the same plane with Bb , or not, and R be the center of gravity of A , and B placed at b ; and QR , the path of the center of gravity, will appear, by the same process with the above, to be parallel to Aa , and equal to $\frac{Aa \times A}{A + B}$, and consequently it varies as Aa , or encreases uniformly.

When both bodies move at the same time, the point P will have two motions PQ and QR , and will consequently describe the diagonal PR uniformly in the time T (185).

2. Let a third body C be added, and the common center of gravity be S , and CS produced will pass through the center of gravity of A and B (411). From the nature of the center of gravity, $A + B + C : A + B :: CP : CS :: CQ : CT :: QP : ST$, and $ST =$
 $A + B$

DL , that is, AB, DE, HK , are in a given ratio, and consequently the point K moves uniformly in HK . The demonstration is the same if one of the bodies moves from B towards A .

CASE II. Let the paths of the bodies AB and DE be in different planes; and through AB draw a plane Bde parallel to DE , and through DE draw the plane $DdeE$ perpendicular to Bde ; produce BA to d , and let Dd, Ee , be perpendicular to de , and the planes DdA, EeB , will be perpendicular to the plane edB . Let A and D, B and E be cotemporary positions of the bodies. If the body at D were to move in de , the center of gravity would move uniformly in some line HK (case 1.); through HK erect the plane $HKkb$ perpendicular to HBK . From similar triangles, and the nature of the center of gravity, $Ab : bD :: AH : Hd :: BK : Ke :: Bk : kE$; therefore bk is the path of the center of gravity of the bodies moving in AB, DE . And, because $Dd : Hb :: Ad : AH :: Be : BK :: Ee$ or $Dd : Kk, Hb = Kk$ and kb is equal and parallel to HK ; therefore the center of gravity of the bodies, moving uniformly in AB, DE , moves uniformly in bk .

CASE III. The common center of gravity of two bodies and a third body, is either at rest, or moves uniformly in a right line; for two may be placed in their common center of gravity, which was proved to move uniformly, and the center of gravity of the three or more bodies is proved, by the same process as before, to move uniformly.

$\frac{A+B \times QP}{A+B+C}$, and varies as QP or uniformly; and for the same reason TV , the motion of T arising from A 's motion, is equal to $\frac{QR \times A+B}{A+B+C}$, and therefore varies as QR or uniformly (case 1.).

When A and B move together, the motions ST , TV , will be combined into one, SV ; and if C describe Cc uniformly in the time T , the common center of gravity will describe VX , and this new motion, combined with SV , will make it describe SX uniformly in the time T . Q. E. D.

437. Cor. 1. It is evident, that the paths of the center of gravity, arising from the motion of any one body, is always parallel to that of the moving body: PQ and ST are parallel to Bb ; QR and TV are parallel to Aa and VX to Cc .

438. Cor. 2. The centers of gravity of two, three, &c. bodies will describe polygons or curves similar to that of the moving body to which their motion is owing; and if the velocity of the body be variable, the velocity of each center will be variable according to the same law.

439. Cor. 3. The velocity of the center of gravity of two, three, &c. bodies is the same as if they were placed in it, and acted upon by forces, equal to the moments of the moving bodies, in their respective planes and directions; for $B \times Bb = A+B \times PQ$ and $A \times Aa = A+B \times QR$; and if $A+B$ were placed at P , and acted upon by forces equal to $B \times Bb$ and $A \times Aa$ in the planes and directions of Bb and Aa , they would describe the diagonal PR (185).

440. PROP. The common center of gravity of two or more bodies is not affected by any action of the bodies upon each other.

DEM.

FIG.
CXXXVIII.

DEM. Let A and B be two bodies in a system, acting upon each other, G their common center of gravity, and Aa , Bb , the velocities lost by A and gained by B respectively in opposite directions, and $A \times Aa = B \times Bb$ (3d law of motion), or $A : B :: Bb : Aa :: BG : AG :: bg : aG$, or $A : B :: B$'s distance from the center of gravity : A 's distance from it; and consequently the same point G is still the center of gravity of A and B , or it has been immovable. What is proved of these two, is true of every two bodies, and therefore of all. Q. E. D.

441. Cor. If two parts of a system A and B , attract or repel each other, or moving with unequal rectilinear motions, disturb each others motion by the force of their inertia, the center of gravity will not be affected by their mutual action.

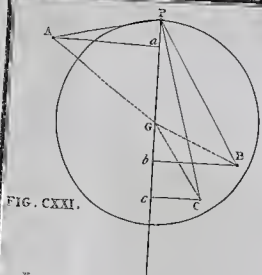


FIG. CXXI.

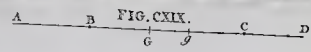


FIG. CXX.

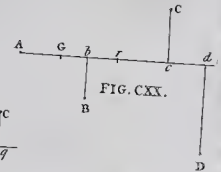


FIG. CXX.

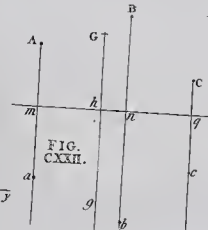


FIG. CXXII.

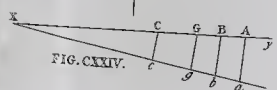


FIG. CXXIV.

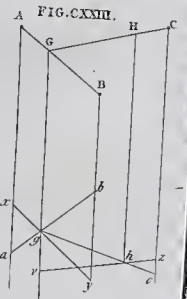


FIG. CXXIII.

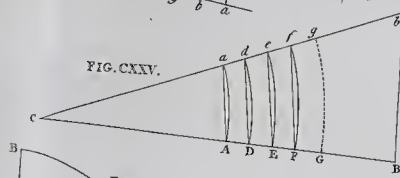


FIG. CXXV.

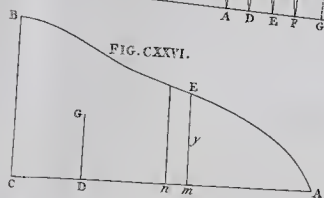


FIG. CXXVI.

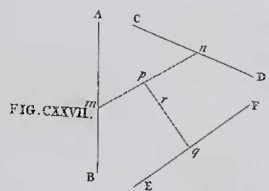


FIG. CXXVII.

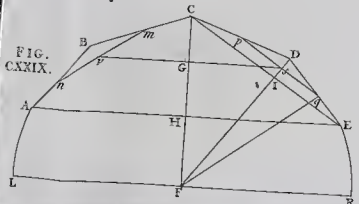


FIG. CXXIX.

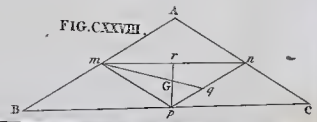


FIG. CXXVIII.

C H A P. X.

COMMUNICATION OF MOTION BY DIRECT IMPACT.

442. DEF. *T*WO bodies are said to impinge directly, when the right line, in which their centers of gravity move, passes through the point of contact.

443. DEF. The center of gravity of a body, or system of bodies, supposed to be without gravity, is the same point with that center when its influence obtains.

LEMMA. If two bodies X and Y , consisting of any number of particles without gravity A, B, C , and D, E , &c. be immoveably connected to the right line SR , passing through their centers of gravity G and H , which is perpendicularly impressed at F by a force F , the moments communicated to these bodies are inversely as the distances of their centers of gravity from the point where F acts.

FIG.
CXXXIX

DEM. The sum of the resistances of D and E to the action of F is $D \times Fd + E \times Fe$ (279) $= \overline{D + E} \times FH$ (407), and consequently if d, e, b , be the respective velocities of D, E , and H , and $D + E = Y$ be placed in H , $D \times d + E \times e = Y \times b$; and for the same reason $A \times a + B \times b + C \times c$, &c. $= X \times g$, supposing a, b, c, g , to be the respective velocities of A, B, C , and of X collected at G . Let a fulcrum be placed at G , and because F is in equilibrio with

with $D \times d + E \times e$ or $Y \times b$, $F : Y \times b :: GH : FG$, and if a fulcrum be placed at H , $X \times g : F :: FH : GH$, and ex æquo $X \times g : Y \times b :: FH : FG$. Q. E. D.

444. Cor. Because $g : b :: FH \times Y : FG \times X :: D \times Fd + E \times Fe : A \times Fa + B \times Fb + C \times Fc$, the velocities of G and H , and of the points A, B, C , &c. are the same, whether the particles be placed in those points, or collected in their centers of gravity G and H ; and the effect is evidently the same, whether G and H be in the right line SR , or in PQ , making an invariable angle with SR , and at the same perpendicular distances from F 's direction.

445. PROP. *If a material surface SUR, composed of the particles A, B, C, D, &c. be impressed by a force equal to F, acting in the plane of the surface at the center of gravity F, and perpendicularly to a right line SFR passing through it, the particles will move with equal velocities in directions parallel to that of F.*

DEM. Let the particles, on one side of F 's direction, $A + B + C$, &c. $= X$ be placed at their center of gravity G , and $D + E$, &c. on the other side at their center of gravity H , and $G FH$ is a right line (411). But $g : b :: Y \times FH : X \times FG :: 1 : 1$ and $g = b$; and the incipient motions of X and Y are parallel to F 's direction (2d law of motion). Since X and Y begin to move with equal velocities, in parallel directions, they are relatively at rest, and consequently will not disturb each others motions; and because the relative situations of the points A, B, C , &c. are always the same, they will move with velocities equal to those of G and H , and in parallel directions. The effects are the same whether A, B, C , &c. be in their real places, or collected in their centers of gravity (444); and the initial motions of G and H are not disturbed by the mutual actions of the particles (441). Q. E. D.

446. Cor. 1. If any particles $B, C, \&c.$ be not in the superficies SR , but on different sides of it, their center of gravity will be in q (411), and their efforts, resulting from their inertia, to disturb the incipient motion of the particles, are measured by the products of each particle into its perpendicular distance from the plane SR (279), and these being equal and opposite, and consequently destroying each other, the particles will continue to move with their initial velocities. If a body therefore be impelled by any power F , in a direction passing through its center of gravity, its constituent particles will move with equal velocities, in directions parallel to that of F ; and if the particles be equal, and their number n , the velocity of each, and of the center of gravity, is the same whether that center be impelled by the force F , or each particle by a force equal to $\frac{F}{n}$, in the direction of F .

FIG.
CXXXIX.

447. Cor. 2. In the impact of one body A upon another body B , either moving or quiescent, when the right line joining their centers of gravity, and the lines in which they move, are in the same right line, the forces of impact and resistance produce a change of velocity, in each body, only in this line, or their impact is direct.

448. PROP.* *Let a body A , moving with a velocity equal to a , impinge directly upon B , moving in the same, or an opposite, direction, with a velocity equal to b , to find their common velocity.*

If A and B be inelastic, (3d law of motion) $Aa \pm Bb$ is the same before and after impact (where the higher sign is to be used when the bodies move in the same direction, and the lower when they move in opposite directions,) and, because they move together, if v be their common velocity, $Aa \pm Bb = \overline{A+B} \times v$ and $v = \frac{Aa \pm Bb}{A+B}$. Q. E. I.

449. Cor.

* Newt. Cor. 3. ad leges motus. Maclaurin's Newt. Chap. IV. Rohault. Not. Part I. Ch. II. Art. 6. Helsham. Lect. V. and Appendix. Nov. Com. Petropol. Tom. XVII. pag. 315.

449. Cor. 1. The velocity, which A loses by impact, is equal to the difference between its velocities before and after impact, or to $a - v$; and the velocity which B gains by impact, in the direction of A 's motion, is equal to the difference between its velocities after and before impact, or to $v \mp b$. Substituting the value of v ,

$$a - v = a - \frac{Aa \mp Bb}{A + B} = \frac{Aa + Ba - Aa \mp Bb}{A + B} = \frac{B \times \overline{a \mp b}}{A + B}$$

$$= l$$
, and, resolving this equation into a proportion, $A + B : B :: a \mp b : l$. The velocity gained by $B = g = v \mp b = \frac{Aa \pm Bb}{A + B}$

$$\mp b = \frac{Aa \pm Bb \mp Ab \mp Bb}{A + B} = \frac{A \times \overline{a \mp b}}{A + B}$$
, and $A + B : A :: a \mp b : g$. From hence we have the following rules.,

1.* The sum of the bodies is to the struck body as the difference or sum (according as they move in the same, or an opposite, direction) of the velocities, before impact, to the velocity lost by the striking body.

2. The sum of the bodies is to the striking body as the difference or sum of the velocities, before impact, to the velocity gained by the struck body.

450. Cor. 2. If A and B move in opposite directions with velocities that are always inversely as the bodies, their common center of gravity will not be affected by their motion; but, l and g , representing the velocities lost by A and gained by B , are made in opposite

* These rules are investigated differently by the following process:

1. $A \times l = B \times g$ (3d law of motion) and $g = \frac{A \times l}{B}$; but $a - l$ (A 's velocity after impact) $= g \pm b$ (B 's velocity after impact) $= \frac{A \times l}{B} \pm b$ and $B \times a - B \times l = A \times l \pm B \times b$ and $l \times \overline{A + B} = B \times \overline{a \mp b}$, and $\overline{A + B} : B :: \overline{a \mp b} : l$.

2. $l = \frac{B \times g}{A}$ and $a - l = a - \frac{B \times g}{A} = g \pm b$, (because the bodies move on together) and $Aa - B \times g = Ag \pm A \times b$ and $\overline{A + B} \times g = A \times \overline{a \mp b}$, and $A + B : A :: a \mp b : g$.

opposite directions, and $A \times l = B \times g$ or $A : B :: g : l$, or g and l are inversely as the bodies, and consequently the velocity of the center of gravity of A and B is not affected by impact, and is, therefore, both before and after impact, equal to their common velocity, v .

451. PROP. *If A, moving with a velocity equal to a, impinge upon B, moving in the same, or an opposite, direction with a velocity equal to b, and A and B be perfectly elastic, or one of them be perfectly hard, and the other perfectly elastic, to find the velocity lost by A (L), and that gained by B (G).*

When A and B are inelastic, $A + B : B :: a \mp b : l$, and
 $A + B : A :: a \mp b : g$; but $L = 2l$,
 and $G = 2g$ (257); therefore $A + B : 2B :: a \mp b : L$, and $A + B : 2A :: a \mp b : G$; and we have these rules.

1. The sum of the bodies is to twice the struck body as the difference or sum of the velocities, before impact, to the velocity lost by the striking body.

2. The sum of the bodies is to twice the striking body as the difference or sum of the velocities, before impact, to the velocity gained by the struck body. Q. E. I.

452. Cor. 1. $L : G :: 2l : 2g :: B : A$ and $A \times L = B \times G$; therefore in all perfectly elastic bodies $L = \frac{B \times G}{A}$ and $G = \frac{A \times L}{B}$.

453. Cor. 2. A 's velocity after impact is equal to the difference between its velocity before impact and the velocity lost, or to $a - \frac{2B \times a \mp b}{A + B} = \frac{Aa + Ba - 2Ba \pm 2Bb}{A + B} = \frac{A - B \times a \pm 2Bb}{A + B}$.

B 's velocity after impact is equal to the sum of its velocity before, and

and the velocity gained, or to $\pm b + \frac{2A \times \overline{a \mp b}}{A+B} = \frac{2Aa \pm Bb \mp Ab}{A+B}$.

If $A = B$, $\frac{\overline{A-B} \times a \pm 2Bb}{A+B} = \pm b$, and $\frac{2Aa \pm Bb \mp Ab}{A+B} = a$,
or the bodies move with interchanged velocities.

454. Cor. 3. If A be greater or less than B , L is less or greater than $a \mp b$, and G greater or less than $a \mp b$. If A be to B , as m to n , $L = \overline{a \mp b} \times \frac{2n}{m+n}$, and $G = \overline{a \mp b} \times \frac{2m}{m+n}$.

455. Cor. 4. If there be any number of bodies A, B, C, D , &c. in geometric progression, encreasing or decreasing, and A , moving with a velocity equal to a , communicate motion to B at rest, and B , moving with the velocity gained, communicate motion to C at rest, &c. the velocities communicated to B, C, D , &c. will be in geometric progression, decreasing or encreasing. For (451)

$$A+B : 2A :: a : \text{velocity of } B$$

$$B+C : 2B :: \text{vel. of } B : \text{vel. of } C$$

$$C+D : 2C :: \text{vel. of } C : \text{vel. of } D$$

$$D+E : 2D :: \text{vel. of } D : \text{vel. of } E. \text{ And, because } A, B, C,$$

&c. are in geometric progression, $A+B : 2A :: B+C : 2B :: C+D : 2C :: D+E : 2D$; and consequently $a : \text{vel. of } B :: \text{vel. of } B : \text{vel. of } C :: \text{vel. of } C : \text{vel. of } D$, &c. And if the bodies encrease or decrease in magnitude, a is greater or less than the velocity of B , or the velocities communicated decrease or encrease.

456. Cor. 5. If the number of bodies in geometric progression be equal to n , $a : \text{velocity of the last body} :: \overline{A+B}^{n-1} : 2A^{n-1}$; for $a : \text{vel. of the last body}$ in a ratio compounded of the ratios of $a : \text{vel. of } B$, $\text{vel. of } B : \text{vel. of } C$, $\text{vel. of } C : \text{vel. of } D$, &c. each of which is equal to the ratio of $A+B : 2A$, and the number of ratios is equal to $n-1$. The velocity of A , or a , is also to the velocity of the last body $:: \overline{A+B} \times \overline{B+C} \times \overline{C+D}$, &c. $: A \times B \times C$, &c. (the product of all except the last) $\times 2^{n-1}$.

457. Cor.

457. Cor. 6. The velocities lost by $A, B, C, \&c.$ are in geometric progression encreasing or decreasing, according as the bodies decrease or encrease; for,

$$A + B : 2B :: a : \text{vel. lost by } A$$

$$B + C : 2C :: \text{vel. of } B : \text{vel. lost by } B$$

$$C + D : 2D :: \text{vel. of } C : \text{vel. lost by } C, \&c.$$

But $A, B, C, \&c.$ being in geometric progression, $A + B : 2B :: B + C : 2C :: C + D : 2D, \&c.$ and consequently $a : \text{vel. lost by } A :: \text{vel. of } B : \text{vel. lost by } B :: \text{vel. of } C : \text{lost by } C, \&c.$; therefore the ratios of the velocities lost and gained are the same, and the last being in geometric progression (456), the first are so too.

458. PROP. *If there be three perfectly elastic bodies, $A, X, Q,$ and $A,$ moving with a velocity equal to a , impinge upon X at rest, and X impinge upon Q at rest, Q 's velocity is the greatest possible when X is a mean proportional between A and Q .*

DEM. Let q be the velocity of Q , and $a : q :: \overline{A + X} \times \overline{X + Q} : 4AX$ (456) $:: A + X + Y + Q : 4A$ (supposing A to be to X as Y to Q , and consequently $A + X : X :: Y + Q : Q$, and for $\overline{A + X} \times Q$, substituting its equal $\overline{Y + Q} \times X$); and $q = \frac{4Aa}{A + X + Y + Q}$, and, $4Aa$

being given, q varies $\frac{1}{A + X + Y + Q}$, and is greatest when their sum is least, or when $X = Y$; for $A \times Q$ being given, $X \times Y$ is given, and the sum of two factors containing a given product is least* when they are equal. Q. E. D.

459. Cor. 1. The velocity, communicated to Q , will be encreased, by encreasing the number of bodies in geometric progression, between it and A , and arrive at its limit when that number

* Let the given product be represented by the given rectangle AB , contained by two unequal lines AL, LB . Make the square LE equal to the rectangle AB and $AL + LB = AD = 2CM$ is always greater than $2LM$ or $LM + LF$.

number is infinitely great, and then $a : q :: \sqrt{Q} : \sqrt{A}$. Let the successive differences of the bodies be $x, y, z, \&c.$ or $B = A + x$, $C = B + y$, $D = C + z$, $\&c.$ and let their respective velocities be $a, b, c, d, \&c.$; because $A + B : 2A :: a : b$ and $B = A + x$, by substitution, $A + A + x : 2A$ or $A + \frac{x}{2} : A :: a : b :: \sqrt{B} : \sqrt{A}$ (31); $B + C : 2B$ or $2B + y : 2B$ or $B + \frac{y}{2} : B :: b : c :: \sqrt{C} : \sqrt{B}$; and, for the same reason, $C + \frac{z}{2} : C :: c : d :: \sqrt{D} : \sqrt{C}$, $\&c.$ but $a : d$ in a ratio compounded of the ratios of $a : b, b : c, c : d, \&c.$ or their equals $\sqrt{B} : \sqrt{A}, \sqrt{C} : \sqrt{B}, \sqrt{D} : \sqrt{C}, \&c.$ or of the square root of the last body to the square root of the first (41); and consequently $a : q :: \sqrt{Q} : \sqrt{A}$.

460. Cor.2. Because $A \times L = B \times G$, $A : B :: G : L$; but G and L are made in opposite directions, and are to each other as the spaces described by the bodies in the same time, and being inversely as the distances from the center of gravity, this center will evidently not be affected by impact, but will continue to move with a velocity equal to v , their common velocity after impact when inelastic. This also appears from article 440.

461. PROP. *If two perfectly elastic bodies A and B impinge directly, the sums of the products, resulting from multiplying each body into the square of its velocity, is the same before and after impact.*

DEM. Retaining the same notation, $a - v = l$, and $2a - 2v = L$, and $2v - 2b = G$; and consequently A 's velocity after impact is equal to $a - 2a + 2v = 2v - a$, and B 's velocity $= \pm b + G = 2v - b$; but $A \times (2v - a)^2 + B \times (2v - b)^2 = A \times 4v \times v - a + a^2 + B \times 4v \times v - b + b^2 = A \times a^2 + B \times b^2$, because $A \times a + B \times \pm b = A + B \times v$, and consequently $A \times v - a +$
 $B \times$

$B \times \overline{v \mp b} = 0$, or, multiplying this equation by $4v$, $A \times 4v \times \overline{v - a} + B 4v \times \overline{v \mp b} = 0^*$. Q. E. D.

462. Cor. 1. Whatever therefore be the number of elastic bodies A, B, C , &c. which impinge successively, beginning with A , the sums of the products, resulting from the multiplication of each body into the square of its velocity, are the same before and after impact.

463. Cor. 2. The relative velocities of A and B are the same before and after impact; for, before impact, the relative velocity is equal to the velocity of A diminished or encreased by that of B , according as they move in the same, or an opposite, direction, or equal to $a \mp b$; and, after impact, it is equal to B 's velocity diminished or encreased by that of A , or equal to $2v \mp b - 2v \mp a = a \mp b$, the same as before impact.

464. Cor. 3. In this proposition the force of elasticity is equal to that of compression, and this obtains very nearly in some kinds of bodies, as glass, ivory, tempered steel, &c. and, in these, the conservatio vis vivæ is preserved, or the products of the bodies into the

* This proposition may be proved a little differently by the following process: let p and q be equal to the velocities of A and B respectively, after impact, and $A \times a \pm B \times b = B \times q \pm A \times p$ (3d law of motion) and $Aa \mp Ap = Bq \mp Bb$. But the relative velocities of A and B are the same before and after impact, or $a \mp b = q \mp p$, and $a \pm p = q \pm b$, and, multiplying $Aa \mp Ap$ into $a \pm p$, and $Bq \mp Bb$ into $q \pm b$, the results are

$$\begin{aligned} A \times a^2 \mp A p a \\ &= A p a - A p^2 = B \times q^2 \mp B q b \\ &= B q b - B b^2, \text{ or} \\ A \times a^2 + B \times b^2 &= B \times q^2 + A \times p^2. \end{aligned}$$

EXAMP. Let the ratios of $A : B$, and $a : b$, be respectively as $1 : 9$, and $6 : 1$, and, before impact, $A \times a^2 + B \times b^2 = 1 \times 36 + 9 \times 1 = 45$, and, after impact, $A \times p^2 = A \times \left[a - \frac{2B \times a - b}{A + B} \right]^2$, and $B \times q^2 = B \times b + \frac{2A \times a - b}{A + B} \right]^2$; the first $= 6 - \frac{18 \times 5}{10} = \frac{60 - 90}{10} = -3$ or 9 , and the second $= 9 \times 1 + \frac{2 \times 5}{10} = 9 \times 4 = 36$, and $A \times p^2 + B \times q^2 = 45$.

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the square of their velocities, are the same before and after impact. But in all bodies, endued with imperfect degrees of elasticity, the equality of these products is not preserved.

465. PROP. *When either A or B is perfectly hard, and the other imperfectly elastic, or when they both are imperfectly elastic, and the whole force of restitution is to that of compression, as any number $n:1$, to find the velocities of A and B after impact.*

DEM. When A and B are inelastic, let the velocities lost by A and gained by B, respectively, by the force of impact, be l and g ; and $1:n :: l:nl =$ the velocity which A loses by the force of restitution, and $nl + l = A$'s whole loss of velocity. Let $m = n + 1$ and $a - ml = A$'s velocity after impact, $=$ (because $l = \frac{B \times g}{A}$) $\frac{Aa - mgB}{A}$. B's velocity after impact $= \pm b + g + ng = \pm b + mg$. Q. E. I.

466. Cor. 1. Because $l = \frac{B \times \overline{a \mp b}}{A + B}$, $\overline{n + 1} \times l$, or A's whole loss of velocity, $= \overline{n + 1} \times \frac{B \times \overline{a \mp b}}{A + B}$, and consequently $A + B : \overline{n + 1} \times B :: a \mp b : \text{velocity lost by A}$. And because $\overline{n + 1} \times g$, or B's whole gain of velocity, $= \overline{n + 1} \times \frac{A \times \overline{a \mp b}}{A + B}$, $A + B : \overline{n + 1} \times A :: a \mp b : \text{velocity gained by B}$.

467. Cor. 2. It is proved in the same manner, as when the bodies were perfectly elastic, that the velocities lost and gained, by a series of imperfectly elastic bodies in geometric progression, which are all at rest, except the first, are in geometric progression; for let
the

the velocities of $A, B, C, D, \&c.$ be $a, b, c, \&c.$ and

$$A + B : n + 1 \times A :: a : b$$

$$B + C : n + 1 \times B :: b : c$$

$$C + D : n + 1 \times C :: c : d,$$

and because $A : B :: B : C :: C : D, \&c.$

$$A + B : n + 1 \times A :: B + C : n + 1 \times B :: C + D : n + 1 \times C, \&c.,$$

or $a : b :: b : c :: c : d, \&c.$

468. Cor. 3. If elasticity be equal to the resistance made to compression, $n = 1$, and, substituting p and q for the velocities of A and B after impact, $p = a - \frac{mBg}{A}$, and $q = mg \pm b$. Let elasticity be $\frac{1}{2}, \frac{1}{3}, \frac{1}{4}, \&c.$ of the resistance, or $n = \frac{1}{2}, \frac{1}{3}, \frac{1}{4}, \&c.$; and $p = a - \frac{\frac{1}{2}Bg}{A}, a - \frac{\frac{1}{3}Bg}{A}, a - \frac{\frac{1}{4}Bg}{A}, \&c.$ and $q = \frac{1}{2}g \pm b, \frac{1}{3}g \pm b, \frac{1}{4}g \pm b, \&c.$: if therefore A, B, g and n be given, the ratio of the relative velocities, or of $a \mp b : q \mp p$ may be found; and $v : v$, if the ratio of the relative velocities, and A, B, a, g be known, the degree of elasticity, or n , may be found: for $p = a - \frac{mBg}{A}$ and $q = \pm b + mg$, and the ratio of $q - p (\pm b - a + g + \frac{Bg}{A} \times m) : a - b$ is a given ratio, and consequently $m(n + 1)$ is known.

469. Cor. 4. The converse of prop. (461) is true, and if the products of A and B into the squares of their velocities, be the same before and after impact, the force of restitution is equal to that of compression; and generally if the products of A and B into that power of the velocity, whose exponent is any number r , be supposed equal before and after impact, the relation between the forces of elasticity and resistance may be found: for $a - \frac{mBg}{A} = p$ and $\pm b + mg = q$, and by solving the equation $A \times a^r + B \times$

$B \times b' = A \times a - \frac{mBg}{A} \Big| + B \times \pm b + mg \Big|, m$, and consequently n may be found.

FIG.
CXLI.

470. LEMMA. *If two sperical bodies A and B move from A and B at the same time, with velocities which are to each other as $a : b$, or $AC : BD$, to find a plane that shall touch both bodies at the point of impact.*

Join AB , and describe the parallogram AH , whose sides are AC and AB ; join DH , and from C , as a center, describe a circular arc, with a radius equal to the semidiaters of the bodies, cutting DH in L ; draw LE parallel to AC , and EF parallel to the line joining C and L , and the bodies will at the same time be in E and F , and a plane perpendicular to EF will touch both bodies at the point of impact; for $DE : EL$ or $FC :: DB : BH$ or AC , and

$$DB : DB \pm DE \text{ or } BE :: AC : AC \pm FC \text{ or } AF,$$

and $DB : AC :: BE : AF :: b : a$, and BE, AF are therefore described in the same time by B and A , and because $EF = CL =$ the sum of the semidiameters of the bodies, they will be in contact at P , with a plane perpendicular to EF . Q. E. I.

FIG.
CXLI.

471. LEMMA. *If two sperical bodies A and B inpinge obliquely, to determine the velocities lost by A and gained by B, and the velocity of each after impact.*

Case 1. Suppose A and B to be inelastic, and to meet in D , and Ll to be a plane in contact with both the bodies, and the velocities to be FD and GD . Resolve each velocity into two, FL , and GM , perpendicular to the plane, and DL, DM , parallel to it: DL and DM are not affected by impact, and FL, GM , are directly opposite to each other. Taking therefore $DH = \frac{A \times FL - B \times GM}{A + B}$

(448),

(448), and Dm , Dl respectively equal to DM , DL , A and B will describe the lines DQ , DR respectively.

Case 2. If A and B be either perfectly or imperfectly elastic, they will be reflected, and the velocities DE , DH be equal respectively

 FIG.
CXLIII.

to $FL - \frac{2B \times FL + GM}{A+B}$ and $-GM + \frac{2A \times FL + GM}{A+B}$ (453):

taking therefore $Dm = DM$ and $Dl = DL$, A will be reflected along DK , and B along DP ; for the velocities DM , DL remain after impact, and these, compounded with DH and DE respectively, will make A and B describe the diagonals of parallelograms, whose sides are DE , Dl , and DH , Dm . Q. E. I.

472. *LEMMA. To find the spontaneous center of conversion of the particles P and Q , without gravity, or a point about which the right line SR , adhering to those particles, begins to revolve, when impressed perpendicularly by a force F , acting at a point F which is not the center of gravity G .

 FIG.
CXLIV.

Let pqs be a new position of SR , infinitely near to the former, and the point S is evidently stationary whilst P and Q describe the small spaces Pp and Qq : but $Q \times FQ : P \times FP :: Pp : Qq$ (443)

$:: SP$ or $SF + FP : SQ$ or $SF - FQ$, and $Q \times FQ \times SF - Q \times$

$FQ^2 = P \times FP^2 + P \times FP \times SF$, and $SF = \frac{P \times FP^2 + Q \times FQ^2}{Q \times FQ - P \times FP}$

$= \frac{P \times FP^2 + Q \times FQ^2}{P + Q \times FQ}$. Q. E. I.

473. Cor. I. Whilst F acts at the same point, the point S is always the same whatever be the magnitude of F , because $\frac{P \times FP^2 + Q \times FQ^2}{P + Q \times F}$ is a given quantity.

474. Cor.

* Philos. Transf. for 1780, pag. 550, where this subject is treated with great perspicuity by Mr. Vince.

474. Cor. 2. To whatever point of $\mathcal{Q}P$ the same force F be applied, the incipient motions of \mathcal{Q} and P or $\mathcal{Q} \times \mathcal{Q}q + P \times Pp$, are invariably the same (2d law of motion); and consequently the velocity of the center of gravity, or Gg , is always the same as if both particles were placed at G and acted upon by the same force F ; for

$$Gg = \frac{P \times Pp + \mathcal{Q} \times \mathcal{Q}q}{P + \mathcal{Q}} \quad (406).$$

FIG.
CXLV.

475. Cor. 3. If SUR be a material surface composed of the particles A, B, C , &c. and a force F act in the line PF placed in the surface, and perpendicular to a line SG drawn through the center of gravity G , each particle will receive an impulse in a direction parallel to PF (2d law of motion); and because G is not affected by the action of the particles upon each other, it will move in a right line (440), and the center of spontaneous conversion will be in some point of the line SG perpendicular to PF . Let S be that center, and the whole plane will begin to revolve about S ; but if a force \mathcal{Q} equal to F were to act in an opposite direction, all motion would evidently be destroyed, and consequently the combined force of every particle in the surface (revolving round S , which results from the action of F), to make it revolve about F and destroy the action of \mathcal{Q} would be equal to nothing. But the force of any particle $A = A \times SA$ multiplied into the perpendicular distance of its direction from $F = A \times SA \times Fm$, or, because (letting fall Sa perpendicular to SF) $Fm : Fn :: Sa : SA$, and $Fm = \frac{Fn \times Sa}{SA}$ and $A \times SA \times Fm = A \times Fn \times Sa = A \times Sa \times \overline{SF - Sn} = A \times Sa \times SF - A \times SA^2$: and by a similar process $B \times Sb \times SF - B \times SB^2 = B$'s force: consequently $SF = \frac{A \times SA^2 + B \times SB^2 \&c.}{A \times Sa + B \times Sb \&c.} = \frac{A \times SA^2 + B \times SB^2 \&c.}{A + B \&c. \times SG}$. If S be the center of suspension, F will therefore be the center of oscillation (art. 512).

476. Cor.

FIG. CXXXVII.

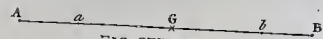
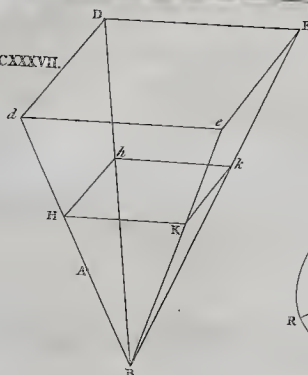


FIG. CXXXVIII.

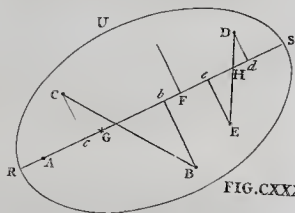


FIG. CXXXIX.

FIG. CXXXIX.

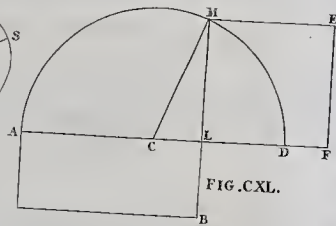
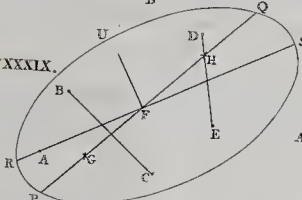


FIG. CXL.

FIG. CXLI

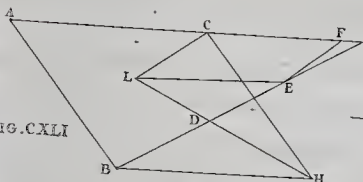
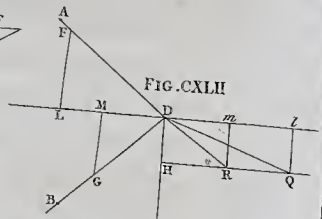


FIG. CXLI



476. Cor. 4. If $p q S$ be the next position of SR , and $x x$ be drawn parallel to it, P has been progressive through $x p$, and Q regressive through $x q$, subtending an angle at g equal to that at S . The initial motion of the center of gravity is not affected by the unequal motions of P and Q , because their actions upon it, or $P \times p g$ and $Q \times q g$ are equal and opposite, and the motion of G is consequently rectilineal and uniform (1st law of motion); and because the incipient angular motions of P and Q about g continue to be uniform, it is evident that when the angle $G S g$ would become equal to four right angles, or P and Q have made one revolution round g , $G g$, the line described by the center of gravity, would be equal to the periphery of a circle whose radius is $S G$.

FIG.
CXLIV.

477. PROP. *Let a body B, be impelled by a force F, acting in the direction FD not passing through its center of gravity G, to define its motion.*

FIG.
CXLVI.

Through G draw a right line SR perpendicular to FD produced, and, if S be the spontaneous center of conversion of the plane surface SUR , every particle will receive an impulse, and begin to move in a direction parallel to this surface (2d law of motion), and consequently to revolve round an axis passing through S perpendicular to it. And if the parts of every section of the body, perpendicular to DF , on each side of the plane SUR , be similar and similarly situated, the plane of conversion SUR will not be affected by either the progressive or rotatory motion of the particles: for their effects upon SUR , arising from their progressive motion, are as the sums of the products resulting from the multiplication of each particle into its distance from this plane, which are equal and opposite; and if p and q be two equal particles, similarly situated in any plane perpendicular to DF , their effects, arising from their rotation, being measured by the product of each particle into the distance from the axis, will be in that plane, and equal and opposite to each other. Consequently the incipient plane of rotation remains unaltered, and if a right line be drawn through G , perpendicular to the plane SUR , every particle

ticle will move round this line as an axis, whilst G moves uniformly in a right line.

478. Cor.1. In a body, B , revolving by the action of a force impelling it in a direction not passing through the center of gravity, the distance of the spontaneous center of conversion from the intersection of F 's direction with a right line passing through the center of gravity, that is, SD is equal to the sum of the particles multiplied into the square of the distance of each, divided by B multiplied into SG .

FIG.
CXLV.

$$479. \text{ Cor.2. } \begin{aligned} A \times SA^2 &= A \times \overline{SG^2} + GA^2 + 2SG \times Ga \text{ (Euc.B.2.P.12)} \\ B \times SB^2 &= B \times \overline{SG^2} + GB^2 + 2SG \times Gb \\ C \times SC^2 &= C \times \overline{SG^2} + GC^2 + 2SG \times Gc \\ D \times SD^2 &= D \times \overline{SG^2} + GD^2 - 2SG \times Gd, \text{ \&c.} \end{aligned}$$

and, because $A \times 2SG \times Ga + B \times 2SG \times Gb + C \times 2SG \times Gc - D \times 2SG \times GD, \text{ \&c.} = 0$, $A \times SA^2 + B \times SB^2 + C \times SC^2 + D \times SD^2 = A \times \overline{SG^2} + GA^2 + B \times \overline{SG^2} + GB^2 + C \times \overline{SG^2} + GC^2, \text{ \&c.}$;

therefore $SF = SG + GF = \frac{A \times \overline{SG^2} + GA^2 + B \times \overline{SG^2} + GB^2}{A + B + C, \text{ \&c.} \times SG}, \text{ \&c.}$

 \&c. ; and $FG = \frac{A \times GA^2 + B \times GB^2 + C \times GC^2}{A + B + C, \text{ \&c.} \times SG}, \text{ \&c.}$, $SG \times GF$

is equal to a given quantity, or SG and GF vary inverfely as each other.

FIG.
CXLVI.

480. Cor. 3. Because the velocity of G is the same as if the particles were concentrated in G , and acted upon by the force F , the time of describing the angle GSg or qgz is the same to whatever point F be applied: but, Gg being given, the angle GSg varies inverfely as SG , or directly as GF (last cor.). If different forces act at the same point, Gg , or the angle of rotation, will evidently be as the magnitude of the force; and, consequently, the angular velocity will vary generally as $F \times GF$, or as the magnitude of the force multiplied into the perpendicular distance of the center of gravity from its direction.

481. PROP.

481. PROP. Let a body, whose quantity of matter is Q , moving with a velocity equal to V , impinge upon the body B , in the direction FC passing through the center of gravity of Q , to find the velocity of the center of gravity, G , of the body B .

FIG.
CXLVI

Let g be the velocity of G , and S the spontaneous center of conversion of B ; and $SG : SD :: g : \text{velocity of } D$, which is therefore equal to $\frac{SD \times g}{SG}$. $V - \frac{SD \times g}{SG}$ is therefore equal to the velocity lost by Q in the direction FC , and (3d law of motiou) $Q \times V - \frac{SD \times g}{SG} = B \times g$, and $\frac{Q \times V \times SG - SD \times g}{SG} = B \times g$, and $g = \frac{Q \times V \times SG}{B \times SG + Q \times SD}$. Q. E. I.

482. Cor. 1. If the bodies be perfectly elastic, the velocity of G , after impact, is equal to $\frac{2Q \times V \times SG}{B \times SG + Q \times SD}$.

483. Cor. 2. Because the velocity of Q , after impact, is equal to that of D , or to $\frac{SD \times g}{SG}$, or, substituting the velocity of g , equal to $\frac{Q \times V \times SD}{B \times SG + Q \times SD}$; $V - \frac{Q \times V \times SD}{B \times SG + Q \times SD}$ is Q 's loss of velocity, when the bodies are inelastic, which is equal to $\frac{B \times V \times SG}{B \times SG + Q \times SD} = Q$'s loss of velocity, and consequently Q 's velocity after impact, when the bodies are perfectly elastic, is equal to $V - \frac{2B \times V \times SG}{B \times SG + Q \times SD} = \frac{V \times Q \times SD - B \times SG}{B \times SG + Q \times SD}$.

484. Cor.

484. Cor. 3. If FC pass through the center of gravity G , of the body B , or the impact become direct, $SG = SD$, and G 's velocity $= \frac{2Q \times V \times SG}{B \times SG + Q \times SG} = \frac{2Q \times V}{B + Q}$. Q 's velocity $= \frac{V \times Q \times SD - B \times SD}{B \times SG + Q \times SD} = V \times \frac{Q - B}{B + Q}$; and rules are deduced from these expressions, which are the same as in direct impact before investigated.

485. Cor. 4. Because $Q \times V = \overline{Q + B} \times g$ (3d law of motion) $g = \frac{Q \times V}{Q + B}$, or the progressive motion of B is the same upon whatever point of B the body Q impinges.

486. Cor. 5. If \mathcal{V} be the magnitude of a body placed at D , which receives the same velocity from the impact of Q , that was communicated to the point D , and, consequently, the velocity of \mathcal{V} , after impact, will be equal to $\frac{Q \times V}{Q + \mathcal{V}} = \frac{Q \times V \times SD}{B \times SG + Q \times SD}$, and $\mathcal{V} = \frac{B \times SG}{SD}$. Hence if the direction does not pass through the center of gravity of Q , find the distance of the spontaneous center of conversion of Q , s , and its center of gravity I , and a body whose magnitude is equal to $Q \times \frac{sI}{sD}$ impinging directly, would have the same effect upon B , whose velocity consequently may be estimated as above.

C H A P. XI.

*CENTERS OF PERCUSSION, GYRATION
AND OSCILLATION.

487. DEF. *THE center of percussion of a body, or system of bodies, is a point, which being stopped by an immoveable obstacle, the body or system is perfectly quiescent.*

If $A, B, C, \&c.$ be particles of a body, or bodies, whose centers of gravity are the points $A, B, C, \&c.$ connected by inflexible lines, and moving with equal velocities, in directions parallel to any right line DN , an immoveable obstacle opposed to the center of gravity G in the direction ND will destroy all motion. For, when G is stopped, each particle will endeavour, by its inertia, to proceed in the line in which it was moving, with a moment equal to the product of its quantity of matter and velocity; and, if any plane be drawn through DGN , the effort of each particle to communicate motion to it varies as its moment multiplied into its perpendicular distance (279). But, the velocities of $A, B, C, \&c.$ being equal, the sums of these efforts to move the plane, on each side of it, are equal (409), and being opposite they consequently destroy each other, and $A, B, C, \&c.$ are perfectly quiescent. If different bodies, or particles of the same body $A, B, C, \&c.$, whose relative situation is unchangeable, revolve about an axis passing through any point S , and a plane be drawn through their center of gravity G perpendicular to the axis, each particle will describe a plane parallel to this plane, their angular velocities will be equal, and their lineal velocities will be as their perpendicular distances from the axis of suspension.

FIG.
CXLVII.FIG.
CXLVIII.

* Simpson's Fluxions, pag. 210. Lyons's Fluxions, pag. 244. Emerson's Mechanics, Sect. VI.

fion. And if, as before, a plane be drawn through G , parallel to the axis, the sums of the moments of A, B, C , &c. multiplied into their perpendicular distances, on each side of this plane, are not equal, because the velocities are not equal; and consequently if G be stopped, A, B, C , &c. will not be quiescent. But the lateral efforts to move the plane passing through SG perpendicularly to the axis, being supposed equal, it is evident that the center of percussion will be in SG .

FIG.
CXLVIII.

488. PROP. *Let A, B, C , &c. be particles of a body, revolving about an axis passing through the point S , it is required to find the center of percussion.*

Let A, B, C , &c. be the places of the particles reduced to the plane passing through SG , perpendicular to the axis, by letting fall perpendiculars from them upon it; and let A 's moment or $A \times SA$, because its velocity is as SA , be represented in quantity and direction by Ap perpendicular to SA , and resolved into two forces Aq perpendicular, and pq parallel to SG . The whole moment of A : its force perpendicular to SG :: $Ap (A \times SA) : Aq$:: $SA : Sa$ (sim. triangles), and $Aq = \frac{A \times SA \times Sa}{SA} = A \times Sa$; and, supposing P to be the center of percussion, the efficacy of A to communicate motion to P is equal to its perpendicular moment multiplied into the distance at which it acts =

$$A \times Sa \times Pm = A \times Sa \times \overline{SP - Sm} = A \times Sa \times SP - \frac{SA^2}{Sa} =$$

$A \times Sa \times SP - A \times SA^2$. By a similar process it appears, that the efficacy of B, C , &c. to communicate motion to P , is $B \times Sb \times SP - B \times SB^2$, $C \times Sc \times SP - C \times SC^2$, &c.; and, because P is quiescent, these forces are equal on each side of it, or $A \times Sa \times SP - A \times SA^2 + B \times Sb \times SP - B \times SB^2 = C \times Sc \times SP - C \times SC^2$, &c. and consequently $SP = \frac{A \times SA^2 + B \times SB^2 + C \times SC^2}{A \times Sa + B \times Sb + C \times Sc}$, &c.

Q. E. I.

489. Cor.

489. Cor. 1. Because $A \times Sa + B \times Sb + C \times Sc, \&c. = \overline{A+B+C},$
 $\&c. \times SG$ (410), $SP = \frac{A \times SA^2 + B \times SB^2 + C \times SC^2}{A + B + C, \&c. \times SG}, \&c.$

490. Cor. 2. The distance of the center of percussion from the center of gravity, or GP , is equal to $\frac{A \times GA^2 + B \times GB^2 + C \times GC^2}{A + B + C, \&c. \times SG},$
 $\&c.$ (479).

491. Cor. 3. In the same body, or system of bodies, $A, B, C,$
 $\&c. \overline{A + B + C}, \&c. \times SG \times GP = A \times GA^2 + B \times GB^2 + C \times GC^2, \&c.$ $SG \times GP$ is also a given quantity, wherever the point of suspension be placed, because $\frac{A \times GA^2 + B \times GB^2}{A + B, \&c.}$ is given;
 and if SG be infinitely great, or the velocities of $A, B, C, \&c.$ be equal to each other, GP is evanescent.

492. Cor. 4. If a circle be described from G as a center, with a radius equal to SG , and, the plane of motion remaining as before, the points of suspension be in different points of the periphery of this circle, the distance between the centers of gravity and percussion will be invariable; for SG being given, GP is given (last cor.)

493. Cor. 5. If particles, whose magnitudes are $\frac{A \times SA^2}{SP^2}, \frac{B \times SB^2}{SP^2},$
 $\frac{C \times SC^2}{SP^2}, \&c.$ be concentrated in P , the same angular velocity will be generated, in these particles, by a force F acting for a given time, and in $A, B, C, \&c.$ at their respective distances $SA, SB, SC, \&c.$ For, let $X, Y, Z, \&c.$ be respectively in equilibrio with $A, B, C, \&c.$ and their moments are inversely as their perpendicular distances (272), or $A \times SA : X \times SP :: SP : SA$, and $A \times SA^2 = X$
 B b 2 $\times SP^2,$

$\times SP^2$, and $X = \frac{A \times SA^2}{SP^2}$. By a similar process it is proved that $Y = \frac{B \times SB^2}{SP^2}$, and $Z = \frac{C \times SC^2}{SP^2}$.

494. Cor. 6. A pendulous body moving with a given angular velocity, will have the greatest effect upon another body against which it impinges, when the point of impact is the center of percussion; for in this case all its motion is communicated. But when the direction of impact does not pass through this center, it will have a lateral motion, or endeavour to continue its rotation.

495. Cor. 7. If the points A, B, C , &c. and G be reduced to any plane perpendicular to the axis, and the center of percussion be supposed to be in this plane, its distance from the axis of suspension is found, by substituting the distances of their places in this plane from the axis, instead of SA, SB , &c. in the expression $\frac{A \times SA^2 + B \times SB^2}{A + B, \text{ \&c.} \times GS}$, &c.

496. LEMMA. *The velocity V , generated in a body whose quantity of matter is Q , by the action of a constant force F , for a given time, will vary as the force directly and quantity of matter inversely.*

DEM. Let V, Q, F , be variable, and it is evident that V will be encreased and diminished in the same ratio with F directly, and in the same ratio with which the inertia or (pag. 63.) Q is diminished and encreased; or V varies as $\frac{F}{Q}$. Q. E. D.

497. Cor. Because the velocity is equal to the moment of a body divided by its quantity of matter, if m represent the quantity of motion generated in a given time by F , or the numeral product of
of

of the velocity and quantity of matter in a given body, V will always be equal to $\frac{m}{Q}$.

498. PROP. If any constant force F act, for a given time, perpendicularly upon the line SV , at the point V , and SV and the particles $A, B, C, \&c.$ whose center of gravity is G , are connected to an axis passing through a fixed point S , the velocity communicated to the point V will be as $\frac{F \times SV^2}{A \times SA^2 + B \times SB^2 + C \times SC^2, \&c.}$

FIG.
CXLVIII.

DEM. Suppose particles, whose magnitudes are $\frac{A \times SA^2}{SV^2}, \frac{B \times SB^2}{SV^2}, \frac{C \times SC^2}{SV^2}, \&c.$ to be concentrated in V ; and they would be in equilibrio with $A, B, C, \&c.$ respectively (493), and consequently equally resist the action of F . But the velocity generated in these particles, collected in V , is as $\frac{F}{Q}$ (496), or as $\frac{F}{A \times SA^2 + B \times SB^2 + C \times SC^2, \&c.}$ or as

$$\frac{F \times SV^2}{A \times SA^2 + B \times SB^2 + C \times SC^2, \&c.} \quad Q. E. D.$$

499. Cor. 1. If m , as before (497), represent the absolute quantity of motion generated by F in a given time, then the velocity of

$$V = \frac{m \times SV^2}{A \times SA^2 + B \times SB^2 + C \times SC^2, \&c.}$$

500. Cor. 2. Because all points describe equal angles at the axis, in equal times, the velocity of V is to the velocity of $A :: SV : SA$, and the velocity of $A =$ velocity of $V \times SA =$

$$\frac{m \times SV \times SA}{A \times SA^2 + B \times SB^2 + C \times SC^2, \&c.}; \text{ and by a similar process, the velocity}$$

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velocity of $B = \frac{m \times SV \times SB}{A \times SA^2 + B \times SB^2 + C \times SC^2}$, &c.; and the

velocity of $C = \frac{m \times SV}{A \times SA^2 + B \times SB^2 + C \times SC^2}$, &c.

501. Cor. 3. The angular velocity of a body, being as its lineal velocity directly and distance inversely, the angular velocity of A , or

B , &c. $= \frac{m \times SV}{A \times SA^2 + B \times SB^2 + C \times SC^2}$, &c. And the moment of a body being measured by the quantity of matter and velocity, the sum of the moments communicated to A , B , C , &c. will be equal to $\frac{A \times SA + B \times SB + C \times SC, \&c. \times m \times SV}{A \times SA^2 + B \times SB^2 + C \times SC^2}$ (last cor.), &c.

FIG.
CXLIX.

502. Cor. 4. If A , B , C , &c. be collected in any point L , the quantity of motion generated in them, in a given time, by the action of the constant force F at V is equal to $\frac{A+B+C, \&c. \times SL \times m \times SV}{A+B+C, \&c. \times SL^2}$

(500) $= \frac{A \times SA + B \times SB + C \times SC, \&c. \times m \times SV}{A \times SA^2 + B \times SB^2 + C \times SC^2}$, &c. and con-

sequently $SL = \frac{A \times SA^2 + B \times SB^2 + C \times SC^2, \&c.}{A \times SA + B \times SB + C \times SC, \&c.}$: but $A \times SA + B \times SB + C \times SC = A + B + C, \&c. \times SG$, and $SL = \frac{A \times SA^2 + B \times SB^2 + C \times SC^2}{A + B + C, \&c. \times SG}$, &c., and L is the center of per-

cussion. If therefore any system of particles be concentrated in the center of percussion, the quantities of motion generated in them, in a given time, placed in that point and at their respective distances, SA , SB , &c., by the action of F at any point V , are the same.

S C H O L I U M.

503. In the preceding proposition and corollaries A , B , C , &c. are supposed to be unaffected by gravity, and motion to be communicated

municated by some constant external force F , which is only resisted by their inertia; but the rules investigated are applicable when F is the force of gravity, or that force acts in conjunction with F : for during an infinitely small time the force of gravity may be considered as constant. If a system of bodies A, B, C , &c. whose center of gravity is G , be suspended from an horizontal axis passing through S , the sum of their efforts to descend is the same as if they were collected in G , and they will therefore descend till SG be perpendicular to the horizon. In any other position of SG , let GL represent the force of descent in a direction perpendicular to the horizon, and be resolved into two forces, GE touching the arc, and EL in the direction of SG . EL is the tendency from S , and GE is that part of the whole force of gravity which produces a rotation about S , and may be deemed constant whilst G describes an infinitely small arc, and is to be added to, or subtracted from, F , according as they conspire with, or oppose, each other.

 FIG.
CL.

CENTER OF GYRATION.

504. DEF. *The center of gyration of a body, or system of bodies, A, B, C , &c. is a point Y , in which, if $A + B + C$, &c. be collected, the same angular velocity will be generated, in a given time, by any constant force F , acting at any point V in a given direction, whether $A + B + C$, &c. be collected in Y , or be at their original distances SA, SB, SC , &c.*

 FIG.
CLI.

505. PROP. *To find the center of gyration of the particles A, B, C , &c. which are connected to an axis passing through the point S , and preserve their relative situation.*

Let SA, SB, SC , &c. be the perpendicular distances of A, B, C , &c. from the axis; and, supposing the quantity of motion generated by F , in a given time, to be m , their angular velocity is equal

to.

CENTERS OF PERCUSSION,

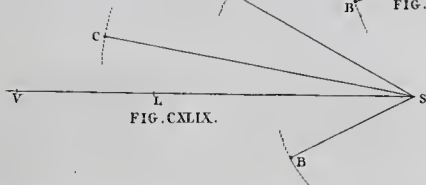
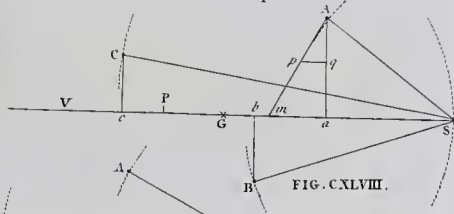
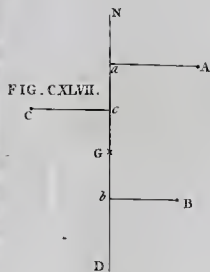
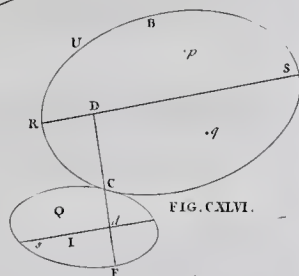
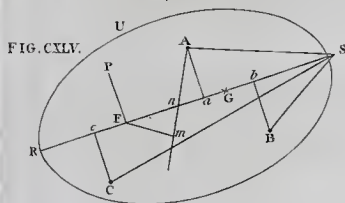
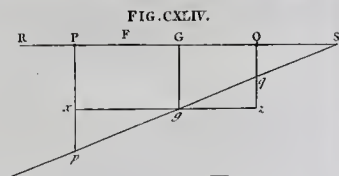
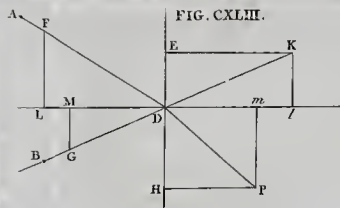
to $\frac{m \times SV}{A \times SA^2 + B \times SB^2 + C \times SC^2, \&c.}$ (501); and, when $A + B + C, \&c.$ are concentrated in Y , their angular velocity = $\frac{m \times SV}{A + B + C, \&c. \times SY^2}$. But, from the definition of the center of gyration, $\frac{m \times SV}{A \times SA^2 + B \times SB^2, \&c.} = \frac{m \times SV}{A + B + C \times SY^2}$, and $SY^2 = \frac{A \times SA^2 + B \times SB^2 + C \times SC^2, \&c.}{A + B + C, \&c.}$, and SY = the distance of the center of gyration from $S = \sqrt{\frac{A \times SA^2 + B \times SB^2 + C \times SC^2}{A + B + C, \&c.}}$, &c.

506. Cor. 1. If therefore p = a particle of a material surface or body, B , d = its perpendicular distance from the axis of suspension passing through S , the distance of the center of gyration from the axis is equal to the sum of all the $\frac{d^2 \times p}{B}$.

507. Cor. 2. If any part of a system of particles $A, B, \&c.$ be collected in their center of gyration, the center of gyration of the whole system will continue the same; for the same force will communicate an equal angular velocity to A and B and $A + B$ placed in their center of gyration. $A, B, C, \&c.$ may consequently be considered as bodies of any magnitude, whose centers of gyration are the points $A, B, C, \&c.$ respectively.

508. Cor. 3. Because $\overline{A + B + C, \&c.} \times SY^2 = A \times SA^2 + B \times SB^2 + C \times SC^2, \&c. = A + B + C \times SG \times SP$ (489), therefore $SY^2 = SG \times SP$, or the distance of the center of gyration from the axis of suspension is a mean proportional between the distances of the center of gravity and the center of percussion from that axis.

509. Cor.



509. Cor. 4. If a system of bodies $A, B, C, \&c.$ whose sum $= Q$, connected as in this proposition, be struck by a given moment at the point V , in a direction perpendicular to a line SV , drawn from V through the center of gravity and perpendicular to an axis passing through S , their angular velocity may be found: for, let R be the center of gyration, G the center of gravity, and the angular velocity of $A, B, C, \&c.$ is the same, whether Q be collected at R , or a body equal to $\frac{Q \times SR^2}{SV^2}$ (489), or $\frac{Q \times SG \times SP}{SV^2}$ be directly opposed to the impinging body at the point V . But $A + B + C, \&c.$ or Q being given, and also the quantity of matter and velocity of the striking body, the velocity of Q may be found by the common rules in the direct impact of bodies (449). And the arc described in a given time by Q and the distance SV , being known, the angle which it subtends may be found. The converse of this, or the velocity of the impinging body, may be found, if its quantity of matter, and Q and its velocity, be given.

510. Cor. 5. In a given system of bodies $A, B, C, \&c.$ whose center of gravity is G , if a circle be described from G as a center with any radius, and the center of suspension S be in its periphery, the distance of the center of gyration from S is always the same; for $SR^2 = \frac{A \times SA^2 + B \times SB^2 + C \times SC^2, \&c.}{A + B + C, \&c.}$, which is a given quantity (412).

511. Cor. 6. If the periphery of a circle, composed of the particles $A, B, C, \&c.$ were to revolve about an axis perpendicular to its plane and passing through its center, the center of gyration will be in the periphery; for $SR = \sqrt{\frac{A + B + C, \&c. \times \text{rad.}^2}{A + B + C}} =$ radius. The angular velocity is therefore the same as if all the matter were collected in any one point of the periphery; and if all the matter in a gyrating body is to be placed in the center of gyration R , it may be placed in any point of the periphery of a
C c circle

CENTERS OF PERCUSSION,

circle whose radius is SY , or collected into two equal portions, and placed in two points diametrically opposite to each other, and whose distances from $S = SY$; for then the center of gravity will be in S , and there will be no lateral motion.

CENTER OF OSCILLATION.

FIG.
CLII.

512. DEF. *The center of oscillation of a body, or system of bodies, A, B, C, &c. connected to an axis passing through S, and moving by the action of gravity, is a point, in which, if the body or system be collected, it, and A, B, C, &c. placed at their original distances SA, SB, SC, &c. will describe equal angles at the axis in the same times.*

513. PROP. *Suppose A, B, C, &c. to be particles of a body connected to an horizontal axis passing through S, and acted upon by the force of gravity in a direction parallel to the vertical line SV, to find their center of oscillation.*

Let SGO be drawn in the plane described by the center of gravity G , and O be the center of oscillation, and the lines SO , SA , SB , &c. will describe equal angles at the axis in equal times. If the weight of A , which is as A (234), acted perpendicularly to SA , its moment would be $A \times SA$: let this be represented by Ap , perpendicular to the horizon, and resolved into two, Aq perpendicular to SA , and qp parallel to it, and this last being lost, Aq is the only efficient part of A 's moment. But $Ap (A \times SA) : Aq :: SA : Aa$ (sim. triangles) and $Aq = \frac{A \times SA \times Aa}{SA} = A \times Aa^*$; and the angular velocity of A generated in a given time is (501) $\frac{A \times Aa}{A \times SA^2 + B \times SB^2 + C \times SC^2}$, &c. By a similar process the angular velocities of B , C , &c. generated in the same time, are — $B \times b$

* This follows immediately from (279), for Aa is equal to the perpendicular let fall from the center of motion upon the direction of A 's weight.

$$\frac{-B \times Bb}{A \times SA^2 + B \times SB^2 + C \times SC^2}, \&c., \frac{-C \times Cc}{A \times SA^2 + B \times SB^2 + C \times SC^2}, \&c.$$
 When $A+B+C, \&c.$ are concentrated in O , the angular velocity of O generated in the same time, would be as $\frac{A+B+C, \&c. \times Oo}{A+B+C, \&c. \times SO^2}$

$$= \frac{Oo}{SO^2} = \frac{Gg}{SG \times SO} \text{ (sim.trian.) : therefore } \frac{A \times Aa - B \times Bb - C \times Cc}{A \times SA^2 + B \times SB^2 + C \times SC^2}$$

$$= \frac{Gg}{SG \times SO}; \text{ but } A \times Aa - B \times Bb - C \times Cc = \overline{A+B+C, \&c.} \times Gg,$$

 and consequently by substitution, $SO = \frac{A \times SA^2 + B \times SB^2 + C \times SC^2}{A+B+C, \&c. \times SG}$.

514. Cor. 1. The distances of the centers of oscillation and percussion from the axis are equal, each being equal to $\frac{A \times SA^2 + B \times SB^2 + C \times SC^2}{A+B+C}$
 $\times SG$, &c., and consequently $SG \times GO$ equal is to $\frac{A \times GA^2 + B \times GB^2}{A+B+C, \&c.}$, &c. (490), and SG and GO are inversely as each other. If SG be infinitely great, or $A, B, C, \&c.$ move with equal velocities in parallel lines, GO is evanescent, or the centers of gravity and oscillation coincide.

515. Cor. 2. If any number of particles be connected to an axis and revolve round it, retaining their relative situation, by the action of a force whose direction is in the plane of motion and perpendicular to the axis, the time of describing any given angle is the same as if they were concentrated in O and urged in the same direction by the same force; and the points where the forces act are transferred from $A, B, C, \&c.$ to O . But in the center of gyration \mathcal{Y} , the same force is supposed to act at the same point, when $A, B, C, \&c.$ are placed at their respective distances $SA, SB, SC, \&c.$ and concentrated in \mathcal{Y} .

516. Cor. 3. If $A, B, C, \&c.$ be bodies composed of any number of particles, let l, m, n be the respective distances of their centers of gravity, and p, q, r , the distances of their centers of oscillation from the axis, and SO will be equal to $\frac{l \times p \times A + m \times q \times B}{A + B + C, \&c.}$

$\frac{+ n \times r \times C}{\times SG}$: for let a, b, c , be particles of A, B, C , and x, y, z , their respective distances from the axis, and, by this proposition, $SO = \frac{\text{sum of all the } a \times x^2 + \text{sum of } b \times y^2 + \text{sum of } c \times z^2}{SG \times A + B + C}$; but the sums

of all the $\frac{a \times x^2}{l \times A}, \frac{b \times y^2}{m \times B}, \frac{c \times z^2}{n \times C}$ are respectively equal to p, q, r , and consequently, the sums of $a \times x^2 + b \times y^2 + c \times z^2 = p \times l \times A + q \times m \times B + r \times n \times C$, and $SO = \frac{l \times p \times A + q \times m \times B + r \times n \times C}{A + B + C, \&c.}$

FIG.
CXLVIII.

517. Cor. 4. If A, B, C , consisting of a number of particles, were urged by a constant force F at the point V , as in (498): let $q, r, s, \&c.$ be particles in A , $t, u, v, \&c.$ particles in B , and $x, y, z, \&c.$ particles in C , and $A = q + r + s, \&c.$, $B = t + u + v, \&c.$, $C = x + y + z, \&c.$; and the velocity communicated to V in a given time is as

$$F \times SV^2$$

$\frac{q \times Sq^2 + r \times Sr^2 + s \times Ss^2, \&c. + t \times St^2 + u \times Su^2, \&c. + x \times Sx^2 + y \times Sy^2, \&c.}{}$; but if P, p, π be the centers of percussion, and G, g, γ the centers of gravity of $A, B, C, \&c.$ respectively, then (489)

$$q \times Sq^2 + r \times Sr^2 + s \times Ss^2, \&c. = SP \times SG \times A$$

$$t \times St^2 + u \times Su^2 + v \times Sv^2, \&c. = Sp \times Sg \times B$$

$$x \times Sx^2 + y \times Sy^2 + z \times Sz^2, \&c. = S\pi \times S\gamma \times C; \text{ and the}$$

$$F \times SV^2$$

velocity of V is as $\frac{SP \times SG \times A + Sp \times Sg \times B + S\pi \times S\gamma \times C}{}$.

FIG.
CLIII.

518. PROP. If O be the center of oscillation when the axis of suspension passes through S , the point S will be the center of oscillation when the axis of suspension passes through O , the plane of motion being supposed to be unaltered.

DEM.

DEM. When S is the point of suspension, the distance of the center of oscillation from the center of gravity G , or $GO = \frac{A \times AG^2 + B \times BG^2 + C \times CG^2}{A + B + C, \&c. \times SG}$, &c. (514); and when the axis passes through O , the distance of the center of oscillation from $G = D = \frac{A \times AG^2 + B \times BG^2}{A + B + C \times OG}$, &c.; and consequently $SG \times GO = D \times GO$ and $SG = D$, and the distance of the center of oscillation from $O = SO$. Q. E. D.

519. Cor. 1. If two circles be described from G as a center with the radii GS and GO , the distances of the centers of oscillation from the points of suspension would be the same in whatever parts of the peripheries they were placed; for SG , or GO being given, the other is given (491), and consequently the times of describing equal angles at the axis will be equal, if the points of suspension be any where in these peripheries.

520. Cor. 2. If the axis of rotation passed through the center of gravity G , and \mathcal{Y} were the center of gyration, $G\mathcal{Y}^2$ would be equal to $\frac{A \times GA^2 + B \times GB^2 + C \times GC^2}{A + B + C}$ (505) $= SG \times GO$, and consequently GO is a third proportional to SG and $G\mathcal{Y}$.

521. Cor. 3. Because $SG \times GO$ is a given quantity, $SG + GO$ is the least possible when the peripheries of the circle, whose radii are SG and GO , meet in \mathcal{Y} , because in that supposition $SG = GO$ (note to 458). And because in this case SO is the least possible, the time of describing a given angle at the axis is the least possible. But because GO encreases as GS decreases, the time of describing a given

a given angle at the axis decreases and arrives at its limit when V and S coincide in the center of gyration \mathcal{X} , and no point of suspension can be assumed, where the time would not be greater than when the axis passes through \mathcal{X} . As GS encreases from $G\mathcal{X}$ to infinity, GO is diminished without limit, and the time of describing a given angle, whether the axis passes through S or O , is perpetually encreased.

C H A P. XII.

RECTILINEAL MOTION OF BODIES.

522. LEMMA. *IF a side AB of a right-angled triangle ABC be divided into very small equal parts, Ab, bc, cd, de, &c. and rectangular parallelograms be described upon them about the triangle, the sum of the rectangles is equal to the area of the triangle, when Ab, bc, &c. are diminished without limit.*

FIG.
CLIV.

For the difference between the triangle and the sum of the rectangles, is the sum of the triangles $Alm + mnp + pqr, \&c.$ which is equal to $\frac{Ab \times BC}{2} = 0$ when Ab vanishes. Q. E. D.

523. PROP.* *If a quiescent body be acted upon by a constant force in the same right line, and any right line AB represent the time, and LM, perpendicular to it, the velocity communicated in the time AL, the space described in the time AL will vary as the triangle ALM.*

DEM. Let AB be divided into very small equal parts $Ab, bc, cd, \&c.$ and let the forces act only at the beginning of each part, and the velocity consequently during each will be uniform. If the velocity communicated at A be represented by bm , the increments communicated at $b, c, d, \&c.$ will be each equal to bm , and consequently LM encreases in the same ratio with AL , and AM is a right line. But the spaces described in the particles of time $Ab,$

* Keil's Physics, Lect. XI. Graves, Lect. I. Ch. XIV: Muschenbroek, Ch. VI. CLVIII. Cotes de Descensu Gravium. Morgan's Notes to Rohault, Maclaurin's Phil. Disc. B. II. Ch. V. Emerson, p. 5, &c.

RECTILINEAL MOTION OF BODIES.

$Ab, bc, cd, ef, \&c.$ are as the rectangles $Am, bp, cr, ds, \&c.$ (107), and the whole space described in the time AL , is as the sum of these rectangles, which, when $Ab, bc, cd, \&c.$ vanish and the force acts incessantly, are equal to the triangular area (ALM). Q. E. D.

524. Cor. 1. If the velocity and time be represented by V and T , and be divided into very small increments V' and T' , whose number is n , $V = n \times V'$ and $T = n \times T'$.

525. Cor. 2. If V, T and v, t represent corresponding velocities and times, $V : v :: T : t$ and $\frac{V}{v} = \frac{T}{t}$; and, if for v a number of feet uniformly described in t seconds be substituted, $V = \frac{v \times T''}{t''}$ will be the number of feet described in t'' by the velocity V .

526. Cor. 3. Whatever be the magnitude of the constant force, the space described, S , will be represented by a triangular area, and consequently S , or the spaces described by the action of different constant forces, are generally as $V \times T$, and, when numbers are substituted for V and T , S is equal to the product of $\frac{V \times T^*}{2}$. In the same constant force, the triangular areas ALM and ABC are always similar, and V and T are directly as each other, and consequently S is as V^2 or T^2 . If therefore the time be divided into equal parts, the spaces described in these times are as the odd numbers 1, 3, 5, 7, &c.; for the spaces described in 1, 2, 3, 4, &c. parts of time, are as 1, 4, 9, 16, &c. and consequently the spaces described in the 1st, 2d, 3d, &c. alone are as 1, 3, 5, 7, &c.

527. Cor.

* This equation may be deduced from finding the sum of an arithmetic progression: let V' and T' be increments of V and T , and $S = 1 + 2 + 3 + 4 \dots n \times V' T' = \frac{n^2}{2} \times V' T'$ (because n is infinitely great, and consequently vanishes compared with n^2) $= \frac{n V \times n T}{2} = \frac{V \times T}{2}$.

527. Cor. 4. The space which is described in any time T from a state of rest, is half of that described, in the same time, with the last velocity, V , continued uniform: for, in the first case, the space $= \frac{V \times T}{2}$ (526), and, in the second, the space $= V \times T$ (107).

528. Cor. 5. If a body, projected with any velocity, be acted upon by a constant force, in a direction opposite to its motion, it will be uniformly retarded: for the force will evidently destroy equal parts of velocity in equal times. A body therefore projected with the last acquired velocity, will ascend to the point from whence it fell, and the velocities in the ascent and descent are the same at the same point. And, if bodies be projected with different velocities and resisted by the same constant force, the time that elapses, T , till the velocity be destroyed, will vary as the velocity of projection V ; the space will vary as the square of the time T^2 , or square of velocity V^2 ; the spaces described in equal portions of time, will be as the odd numbers in a retrograde order 7, 5, 3, 1; and, whether these retarding forces be the same or different, the space will vary as $V \times T$, and be equal to $\frac{V \times T}{2}$.

S C H O L I U M.

529. At the surface of the earth the force of gravity is constant, for a body descends through $16\frac{1}{2}$ feet nearly in the first second, and consequently acquires a velocity that would make it describe $32\frac{2}{12}$ feet uniformly in 1"; and this is found to be the velocity communicated in every succeeding second. Any of these quantities, space described S , velocity acquired, or time T , being known, the others may be discovered.

1. The proper measure of velocity is the space uniformly described by it in any given time: if therefore a body fall by the force of gravity for t'' , and acquire a velocity which would make

D d

it

RECTILINEAL MOTION OF BODIES.

it describe V feet uniformly in $1''$, V and 32 are the proper measures of the velocities communicated in t'' and $1''$ (106), and $V : 32 :: t'' : 1''$ (525), and $V = 32 \times t''$. If $t = 2, 3, 4, \&c.$ seconds, $V = 2 \times 32, 3 \times 32, 4 \times 32, \&c.$ feet respectively.

2. If V , or the number of feet uniformly described in $1''$, by the velocity acquired in falling for an unknown time, be given, this time may be found; for it is equal $\frac{V}{32}$ seconds. If $V = 192$ feet, the time of falling $= \frac{192''}{32} = 6''$.

3. If \mathcal{V} be the space described in falling t'' , $16 : \mathcal{V} :: 1^2 : t''^2$ (526), and $\mathcal{V} = 16 \times t''^2$ feet. If $t = 2, 3, 4, \&c.$ seconds, $\mathcal{V} = 16 \times 2^2, 16 \times 3^2, 16 \times 4^2, \&c.$ feet respectively. Or \mathcal{V} may be found by knowing the velocity (V) at the end of the descent; for $\mathcal{V} : 16 :: V^2 : 32^2$, and $\mathcal{V} = \frac{16 \times V^2}{32^2} = \frac{16 \times V^2}{4 \times 16^2} = \frac{V^2}{64}$ feet. If $V = 192$ feet, $\mathcal{V} = \frac{192^2}{64} = 576$ feet.

4. If \mathcal{V} , or the number of feet described in falling, be known, the time of descent, and velocity acquired, may be found; for $\mathcal{V} = 16 \times t''^2$, and consequently $t'' = \frac{\sqrt{\mathcal{V}}}{4}$. If $\mathcal{V} = 576$ feet, $t = \frac{\sqrt{576}}{4} = \frac{24}{4} = 6$ seconds. And because $\mathcal{V} = \frac{V^2}{64}$, $V = \sqrt{\mathcal{V}} \times 8$. If $\mathcal{V} = 576$ feet, $V = \sqrt{576} \times 8 = 24 \times 8 = 192$ feet.

530. PROP. *If bodies be acted upon by different constant forces, the velocities communicated will vary in a ratio compounded of the forces and times.*

Let

Let S , V , T , represent force, velocity and time, and be supposed variable, and it is evident that the velocity will be encreased and diminished in the same ratio with both the force and time, and, these being independent of each other, V will be as $F \times T$. Q.E.D.

531. Cor. 1. V therefore is as $F \times T$, and if F be compared with the force of gravity f , or any other known force, capable of generating a velocity equal to v in the time t , then $V : v :: F \times T : f \times t$, and $\frac{V}{v} = \frac{F \times T}{f \times t}$.

532. Cor. 2. Because, in all constant forces, the space varies generally in a ratio compounded of the velocity and time, and the velocity varies generally as the force and time; the following analogies are general :

S is as $V \times T$, or as $F \times T^2$, or as $\frac{V^2}{F}$;

T varies as $\frac{S}{V}$, or as $\sqrt{\frac{S}{F}}$;

V varies as $\frac{S}{T}$, or as $\sqrt{S \times F}$; and

F varies as $\frac{V}{T}$, or as $\frac{S}{T^2}$, or as $\frac{V^2}{S}$.

If numbers be substituted for V , F , and T ; or the number of feet described in a given time 1, by the action of any constant force F , be substituted for the force; the number of feet which would be uniformly described in this given time 1, by the velocity acquired in the time T , $= V$; and $T =$ the number of parts of time each equal to 1, contained in the whole time in which the velocity is communicated: then (106) V and $2 F$ are proper measures of the velocity acquired in the times T and 1, and $V : 2 F :: T : 1$ (525) and $V = 2 F \times T$; and substituting this value of V in the general equation ($S = \frac{V \times T}{2}$), the following general equations are deduced :

D d 2

$S =$

$$S = \frac{V \times T}{2} = F \times T^2 = \frac{V^2}{4F} \text{ feet;}$$

$$T = \frac{2S}{V} = \sqrt{\frac{S}{F}}, \text{ equal portions of time, each of which is equal to } 1; \text{ and}$$

$$V = \frac{2S}{T} = \sqrt{S \times 4F} \text{ feet.}$$

533. Cor. 3. These equations and analogies are applicable to all accelerating and retarding forces, such as gravity, and to all resistances that are constant and produce equal decrements of velocity in equal times. If the motion of a ball shot into a bank of earth or piece of wood, be uniformly resisted, and the magnitude of this resistance compared with gravity, and the velocity of impact, be known; the depth, or space through which it moves before all motion is destroyed, may be discovered; and vice versâ*.

FIG. 534. PROP. *If a body descend down an inclined plane AB by the*
CLV. *action of a constant force F, whose direction is perpendicular to any right*
line BC; F will be to that part of F which makes the body descend, as
the radius to the sine of the plane's inclination to BC.

DEM. Let any given line LM , perpendicular to BC , represent the constant force F , in quantity and direction, and be resolved into two forces, LN parallel, and MN perpendicular, to the plane; MN is destroyed by the resistance of the plane, and LN is the only part of F that communicates motion to the body; therefore

$F:$

* EXAMPLE. Suppose a resistance were equal to the resistance of gravity F , when a body is projected perpendicularly upwards, multiplied into any number n , and all motion were destroyed in t'' , then S (being equal to the force multiplied into the square of the time) $= n \times 16 \times t^2$. If the velocity of projection were such as would make a body describe f feet in a second uniformly, then S (being equal to $\frac{\text{vel.}^2}{4 \times \text{force}}$) $= \frac{f^2}{4 \times n \times 16}$
 $= \frac{f^2}{n \times 64}$. If the space described and force be known, then the time is equal to $\sqrt{\frac{S}{n \times 16}}$
 and the velocity $= \sqrt{S \times n \times 64}$ feet.

F : that part of F which makes the body move (A) :: $LM : LN$:: $AB : AC$ (sim. triangles) :: rad. : fin. of $\angle ABC$. Q. E. D.

535. Cor. 1. F is to that part of F which is destroyed by the reaction of the plane as $LM : MN$:: $AB : BC$ (sim. triangles) :: radius : cof. $\angle ABC$; and consequently the part of F , destroyed by the resistance of the plane, $= \frac{F \times \text{cof. } \angle ABC}{\text{radius}}$, which is constant upon the same, or parallel, planes, and upon different planes, not parallel, varies as the $\frac{\text{base}}{\text{length}}$, if F be given.

536. Cor. 2. Because the accelerating force $LN = A = \frac{F \times AC}{AB}$, and F is given, A is constant upon the same, or parallel planes, and, upon different planes, inclined in unequal angles to BC , it varies as the height divided by the length, or as $\frac{H}{L}$, calling H the perpendicular height and L the length.

537. Cor. 3. Because the force upon the same plane is constant, the analogies and equations in (526) are applicable to the motion of a body upon the same plane; and the analogies and equations in (532) obtain when bodies move upon different planes. If BC be horizontal, and F be the force of gravity, then the space described upon the plane in $T'' = S = A \times T''^2 = F \times T^2 \times \frac{H}{L}$

$$(536) = \frac{16 T^2 \times H}{L}; \quad V = \frac{2F \times T'' \times H}{L} = \frac{32 T'' \times H}{L} = 8 \times \sqrt{H}$$

feet; and $T'' = \frac{V \times L}{32 \times H} = \frac{L}{4 \times \sqrt{H}}$.

538. Because the space described is generally as the force multiplied into the square of the time (532), the spaces that are to each:

each other as the forces must be described in the same time, or when S is as F , T is given. If therefore two bodies descend at the same instant, one down the plane AB , and the other in AC perpendicular to BC , and CP be drawn perpendicular to AB , P and C are cotemporary positions of the bodies; for $AP : AC :: LN : LM :: A : F$, that is, the space is as the force, and T^2 and T are given. If the diameter AD of a circle be perpendicular to the horizon, a body will descend through it, and any chord AP drawn from its extremity A , in the same time, because the angle APD is a right one; and the times of descent through all the chords drawn from A are consequently equal to each other.

FIG.
CLVI.

539. Cor. 5. The velocities acquired in falling down different inclined planes, are as the square roots of their perpendicular heights; for V^2 is always as the space multiplied into the force, or

FIG.
CLV.

as $L \times A$, or as $\frac{L \times H}{L}$ (536), or as H , and V is as \sqrt{H} ; this also follows from art. 537. The velocities therefore acquired in falling down the perpendicular AC , and any planes AB , AE , AG , &c. drawn from A and terminated by the base CB , are equal to each other. And because V is always equal to $2A \times T$, and in this cor. V is given, the times of falling are as the forces inversely, or as the lengths of the planes. Or because S is as $V \times T$, and V is given, T is as S , that is, as the length.

540. Cor. 6. The times of descending through different planes are as the lengths directly and square roots of the heights inversely; for the time is generally as the space described directly and last velocity inversely (532), or T is as $\frac{L}{V}$, or as $\frac{L}{\sqrt{H}}$ (last cor.); this also follows from art. 537. The times therefore of describing different planes, equally inclined to the direction of F , are as the square roots of their lengths; for H is as L (sim. triangles), and consequently $\frac{L}{\sqrt{H}}$ is as $\frac{L}{\sqrt{L}}$, or as \sqrt{L} or T .

541. PROP.

541. PROP. *The velocity acquired in falling down any number of contiguous planes, supposing none to be lost in passing from one plane to another, is equal to the velocity acquired in falling down the same perpendicular altitude.*

FIG.
CLVII.

DEM. The velocities acquired in falling down AB and EB , $AB + BC$ and EC or PC , $AB + BC + CD$ and PD , or down the perpendicular PR , are equal (539). Q. E. D.

542. Cor. The velocities, acquired in falling down any systems of planes, are therefore as the square roots of their perpendicular heights; and if the body be projected from D with the velocity acquired, it will ascend through this or any other system of planes to the same perpendicular altitude.

543. PROP. *The velocity lost in passing from any plane to the next, is to the whole velocity as the sine of the angle of their inclination to the radius.*

FIG.
CLVIII.

DEM. Let CB represent the velocity acquired in falling down AC , and be resolved into two, CP coincident with the plane CD , and BP perpendicular to it; and this last is evidently destroyed by the resistance of the plane: but $BP : CB :: \sin. \text{ of } \angle PCB \text{ or } \angle PCA : \text{radius}$. Q. E. D.

544. Cor. 1. $CB : CP$ (velocity upon the plane CD) $:: \text{radius} :: \cosine \text{ of } \angle PCB \text{ or } PCA$. If the angle ACD become equal to two right angles, or ACD be any circular arc, CP and CB coincide, or the velocity lost is equal to nothing; and consequently a body moving in any circular arc sustains no loss of velocity by changing the direction of its motion; and a body, projected with the velocity acquired, up any curvilinear arc, will rise to the same height from whence it fell.

545. Cor.

545. Cor. 2. If ACD be any circular arc, the velocity acquired in falling from A to D is consequently equal to that acquired in falling through the same perpendicular height; and the velocities acquired or destroyed in descending or ascending through two curve lines are as the square roots of their perpendicular altitudes: for every curve may be considered as consisting of an indefinite number of inclined planes.

FIG.
CLVII.

546. PROP. *The times of describing two systems of inclined planes $ABCD$ and $abcd$, whose number, inclinations, and ratios of their lengths are the same, are to each other as the square roots of the lengths of the planes.*

DEM. Because the planes are equally inclined to the direction of the force, the time of falling down AB is as \sqrt{AB} (540), or as \sqrt{EB} , or \sqrt{EC} (hypoth.), or as the time down EB , or EC ; and, dividendo, the time of falling down BC , after having fallen down AB or EB , is as \sqrt{AB} ; consequently the time of falling down $AB + BC$ is as \sqrt{AB} , or as $\sqrt{AB + BC}$, &c.: for from the supposition $\sqrt{AB} : \sqrt{ab} :: \sqrt{AB + BC} : \sqrt{ab + bc}$, &c. Q. E. D.

547. Cor. 1. The times of describing similar parts of similar curves, equally inclined, in similar parts, to the direction of the force, are as the square roots of their lengths.

548. Cor. 2. If two bodies vibrate in similar circular arcs, the times of performing these vibrations are as the square roots of the lengths of these arcs, or as the square roots of their radii.

SCHOLIUM.

549. In the preceding propositions and corollaries, inclined planes and bodies, descending or ascending upon them, are supposed to be without any asperities upon their surfaces, and the bodies are supposed to move by the action of a force in a given direction,

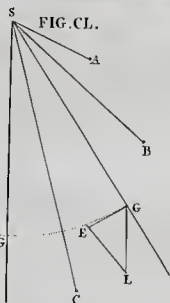


FIG. CL.

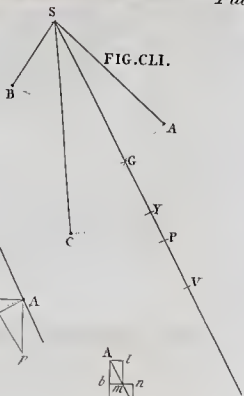


FIG. CLI.

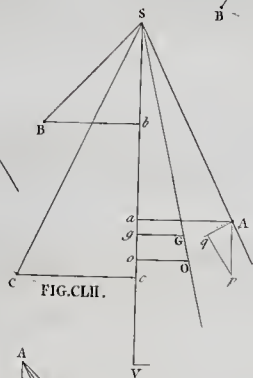


FIG. CLII.

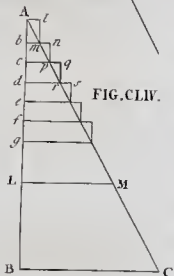


FIG. CLIV.

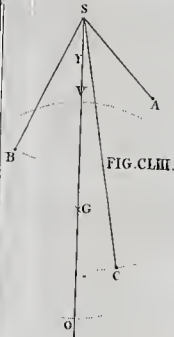


FIG. CLIII.

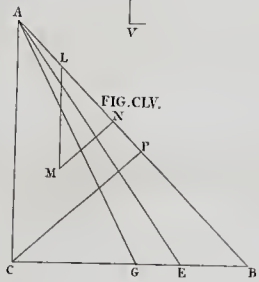


FIG. CLV.

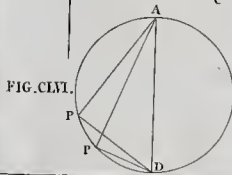


FIG. CLVI.

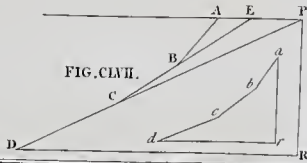


FIG. CLVII.



direction, without any retardation from friction, or rotation about an axis. In this supposition, the tendency of each particle of a body to descend, when impelled by the force of gravity, is equal to the particle multiplied into $\frac{H}{L}$, and the tendency of the body W is equal to $W \times \frac{H}{L}$, and is the same as if it were collected in its center of gravity. But if the parts of W and the inclined plane adhere together, the force of this adhesion must be estimated and allowed for in practice. If W be a spherical body, and the parts of its surface in contact with the plane adhere by the pressure of W upon it, their tendency to descend will be diminished by this adhesion, and consequently they will be regressive with regard to the center of gravity G , and revolve round it. The absolute velocities of different particles in W , and consequently their accelerating forces, are different; and the tendency to descend of the center of gravity is diminished by this motion of rotation round it. The angular velocity of every particle round G is the same, and the same as if all the particles were concentrated in their center of gyration Y , or a body whose magnitude is $\frac{W \times GY^2}{GA^2}$ were collected in A (498).

550. PROP. *Suppose the same spherical body to slide and roll down the same inclined plane, to find the ratio of the forces acting upon it.*

FIG.
CLIX.

From the nature of the circle, the initial regressive velocity of the point A is equal to the progressive velocity of the center of the sphere, or center of gravity G ; but the regressive velocity of a body whose magnitude is $\frac{W \times GY^2}{GA^2}$ placed at A (last scholium), would be destroy-

ed by the action of a force in a contrary direction equal to $\frac{H}{L}$; and in this case G would not be impeded by the regressive motion of A , but G and A would slide down with equal velocities; consequently

E e

the

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the force acting upon W when it slides : the force when it rolls down the plane $:: W + \frac{W \times GY^2}{GA^2} \times \frac{H}{L} : W \times \frac{H}{L}$. Q. E. I.

551. Cor. If the center of gyration be found, and GY , and GA be expressed in numbers, the ratio of these forces will be expressed in terms of the weight.

S C H O L I U M.

552. In the communication of motion by the external application of forces such as impact, protrusion, &c. their magnitude, or capacity to communicate motion, is not always measurable by their cotemporary effects; but in this chapter every particle of matter is supposed to be impelled by a force, similar to that of gravity, with the same intensity in parallel directions, which will consequently communicate the same velocity, or change of velocity, to bodies of different magnitudes, whether quiescent or moving. The magnitude of these forces at any instant, or the magnitude of their accumulated action, is measurable by their cotemporary effects, or by the velocity generated or destroyed in equal times; for the increments or decrements of velocity, produced in equal times, by the same constant force, being equal, the whole velocity is increased or diminished uniformly; and when the forces are unequal, the changes of velocity being as the forces, it is evident that the velocities are proper measures of the intensity of forces, and will vary as the forces multiplied into the times in which they are generated or destroyed. And the magnitudes of variable forces, whose intensities and cotemporary effects are perpetually increased or diminished, are measured at any instant of time by their constant action at that instant, or by the velocities which they would produce in the same time were their intensity constant during that time; and the magnitude of their accumulated actions for any finite time is measured by the addition of these cotemporary effects.

553. PROP.

453. *PROP. If F, V, T , represent any finite variable force, velocity and time, respectively, and \dot{V}, \dot{T} , be any small changes of V and T , \dot{V} will vary as $F \times \dot{T}$.

DEM. The force F being always finite, and being supposed to encrease or decrease according to the same law, an encrease or diminution of it in any finite time must be finite, and consequently, in an infinitely small time, evanescent compared with the whole force F , which therefore may be considered as constant during the time \dot{T} in which \dot{V} is generated or destroyed; therefore (530) \dot{V} varies as $F \times \dot{T}$. Q. E. D.

554. Cor. 1. If \dot{V} be the fluxion of the velocity, or the change of velocity generated by F at any instant, supposed to act constantly for any time \dot{T} , \dot{V} will therefore vary as $F \times \dot{T}$.

555. Cor. 2. \dot{T} is therefore as $\frac{\dot{V}}{F}$ and F as $\frac{\dot{V}}{\dot{T}}$, and the relation of any two being known, the other may be found; and because \dot{T} is as $\frac{\dot{S}}{\dot{V}}$, $V\dot{V}$ is as $F \times \dot{S}$.

556. PROP. If any right line AS represent the time, and OR be as the force at any point O , the change of velocity at O will be as the area AR .

FIG.
CLX.

DEM. The increment of velocity \dot{V} is as $F \times \dot{T}$, or as $AB \times Ap$ or as Aq ; and the next increment of velocity is as ps , and the sum of the increments, or the whole change of velocity, is as the sum of these areas; consequently the velocity communicated in the time AO is as AR . Q. E. D.

557. Cor. 1. If the time in acquiring any given velocity be known, and the areas AR can be squared, the velocities communicated in any other time AO may be found; and vice versâ.

558. Cor.

* Newt. Prin. Tom. I. Sect. VII. Euler's Mechan. Ch. II.

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558. Cor. 2. If the force F be finite, and the time Ap infinitely small, the increment of velocity, \dot{V} , is evanescent; because $AB \times Ap =$ a finite quantity multiplied into one that is evanescent $= 0$; and consequently the increment of velocity, or \dot{V} , is infinitely small, and no finite change of velocity can be generated by F in an instant.

559. Cor. 3. If the body ascend or descend, and the velocity be as the time in which it is acquired or destroyed, the force is inva-
riable; for, from this proposition, the change of velocity is as AR , or as AO (hypoth.), and consequently AR is a rectangular parallelogram, or OR is constant.

FIG.
CLXI.

560. PROP. If PO represent the space described by the action of a force tending to S , and the ordinate OR be the relative magnitude of the force at the point O , the velocity, V , acquired at O will vary in a subduplicate ratio of the corresponding area PR .

DEM. Let V, F, S , be the velocity, force, and space described, respectively, and because \dot{V} is as $F \times \dot{T}$ (554), and \dot{T} is as $\frac{\dot{S}}{V}$; $V\dot{V}$ is as $F \times \dot{S}$, or as $PB \times Pp$, and the fluent of $V\dot{V}$, or $\frac{V^2}{2}$ is as the fluent of $PB \times Pp$ or PR , and V is as \sqrt{PR} . Q. E. D.

561. Cor. 1. If x be the distance of the moving body from the center of force, S , and the force be as any power of the distance, whose exponent is $n - 1$; $V\dot{V}$ is as $F \times -\dot{x}$, or as $-x^{n-1}\dot{x}$, and $\frac{V^2}{2}$ is as $-\frac{x^n}{n} + \text{correction}$: but when $V = 0$ at P , $x = SP = p$, and $\frac{x^n}{n} = \frac{p^n}{n}$; therefore the correction $= \frac{p^n}{n}$, and V is as $\sqrt{p^n - x^n}$.

562. Cor.

562. Cor. 2. If the force be constant, or $n = 1$, then V being as $\sqrt{p^n - x^n}$ is as $\sqrt{p - x}$, or in a subduplicate ratio of the space described, either in acceding to S from a state of rest, or receding from S till the velocity be destroyed. The fluxion of the time \dot{T} is as the fluxion of the space directly and velocity inversely; or as $\frac{-\dot{x}}{\sqrt{p - x}} = -\dot{x} \times (p - x)^{-\frac{1}{2}}$, and T is as $\frac{(p - x)^{\frac{1}{2}}}{\frac{1}{2}}$, and when $T = 0$, $x = p$, and $\frac{(p - x)^{\frac{1}{2}}}{\frac{1}{2}} = 0$; therefore T is as the square root of the space described from rest, which coincides with art. 526.

563. Cor. 3. If the force vary directly as the distance, or $n = 2$, then V varies as $\sqrt{p^2 - x^2}$, or as the right sine of a circular arc, described from S as a center, and radius equal to SP or p , whose versed sine is the space described. And describing this circle, the fluxion of the time \dot{T} is as $\frac{\dot{S}}{V}$, or as $\frac{Oo}{OR}$, or as $\frac{Rn}{Sn}$ (sim. triang.), and the fluent T is as the fraction $\frac{PRD}{SD}$, which is a constant quantity. If the force therefore vary directly as the distance, the times of descent to S from a state of rest are equal wherever P be taken.

FIG.
CLXII.

564. Cor. 4. If the force vary inversely as any power of the distance, or n be negative, or less than 1, let this power be expressed by the number $-m - 1$; and V is as $\sqrt{\frac{1}{m \times x^m} - \frac{1}{m \times p^m}}$, or as $\sqrt{\frac{p^m - x^m}{m \times p^m x^m}}$. If the force be inversely as the distance, or $m = 0$, this expression does not shew the variation of velocity. If the force be inversely as the square of the distance, or $m = 1$, V is as $\frac{\sqrt{p - x}}{\sqrt{m \times p x}}$, or the velocity is as the square root of the space described directly, and inversely as the square root of the distance from S .

565. PROP.

FIG.
CLXIII.

565. PROP. *If the ordinates PB, AL, OR, be always proportional to the forces at those points, and the force at A, continued constant through the space AM, communicate a velocity, V, equal to that acquired by the body at O, descending by the variable force, the area PR will be equal to the rectangle AN, and the body must have fallen from P.*

DEM. Let \dot{v} and \dot{V} be the fluxions of the velocities communicated, by the variable and constant forces respectively, in any periods of their descent; and $\dot{v} : \dot{V} :: \frac{PB \times P\dot{O}}{v} : \frac{AL \times A\dot{M}}{V}$ (555),

and $v\dot{v} : V\dot{V} :: PB \times P\dot{O} : AL \times A\dot{M}$, and, taking the fluents, $v^2 : V^2 :: PR : AN$. If therefore $v = V$, the areas PR and AN are equal, and the body must have fallen from P . Q. E. D.

566. Cor. 1. If the areas PR and AN be equal, the velocities, acquired by the action of the constant and variable forces, whilst the bodies describe, from rest, the spaces AM and PO , will also be equal.

567. Cor. 2. Because the areas AN and PR are always equal, when the velocities at M and O are equal, their increments Or and Mn must be equal; and if F be the given force at A , and $y =$ the space described AM , $F \times \dot{y} = Or = OR \times Oo =$ the force at O multiplied into the fluxion of the space.

568. PROP. *If a body be attracted towards the point S, by forces which always vary as that power of the distance whose exponent is $n-1$, and begin to move at any given distance SP, it is required to assign the velocity, V, acquired in describing any space PO, supposing the magnitude of the force at any given distance SA to be known.*

Let

Let the magnitude of the force at A be to the force of gravity as $F : 1$, and let

$$SA = a,$$

$$SP = p,$$

$SO = x$; and, from the supposition, F : force at any point

$O :: a^{n-1} : x^{n-1}$, and the force at $O = \frac{F \times x^{n-1}}{a^{n-1}}$. But, if y be the

space through which a body must fall, when acted upon by a constant force equal to F , to acquire a velocity equal to that at O ,

from (567) $F \times y = \frac{F \times \overline{x^{n-1} - x}}{a^{n-1}}$, and the fluents are equal, or

$F \times y = \frac{-F x^n}{n \times a^{n-1}} + \text{correction}$; but when $F \times y$ vanishes, $x = p$,

and consequently the fluent corrected $= \frac{F \times \overline{p^n - x^n}}{n \times a^{n-1}}$. But if s

be the space described by the force of gravity in 1'', or any other given time 1, and the force of gravity be expressed by 1, $V =$

$$*\sqrt{4sFy} = \sqrt{\frac{4sF \times \overline{p^n - x^n}}{n \times a^{n-1}}} \text{ feet in 1''}. \quad Q. E. I.$$

569. Cor. 1. If the force be as some negative power of the distance, or $n - 1 = -m$, then F : the force at $O :: a^{-m} : x^{-m}$, and

the force at $O = \frac{F \times x^{-m}}{a^{-m}}$; the fluxion of $F \times y = \frac{F \times x^{-m} \dot{x}}{a^{-m}}$ and

its fluent $= \frac{F \times x^{1-m}}{1-m \times a^{-m}} + \text{cor.} = \frac{F \times a^m}{1-m \times x^{m-1}} - \frac{F a^m}{1-m \times p^{m-1}}$

$$= \frac{F \times a^m \times \overline{p^{m-1} - x^{m-1}}}{1-m \times p^{m-1} x^{m-1}}, \text{ and } V = \sqrt{\frac{4sF \times a^m \times \overline{p^{m-1} - x^{m-1}}}{1-m \times p^{m-1} x^{m-1}}}.$$

570. Cor.

* Let $f =$ the force of gravity,

$s =$ the space described by the constant action of f in 1'' = 16 feet nearly, and

$v =$ the velocity acquired in falling 1''; and

$V^2 : v^2 :: F \times y : f \times s$ and $V^2 = \frac{v^2 \times F \times y}{f \times s}$; but if $f = 1$, $v = 2s$ and $V^2 = 4sFy$ and

$$V = \sqrt{4sFy}.$$

RECTILINEAL MOTION OF BODIES.

570. Cor. 2. Because $4Fs$ and a^{n-1} are given quantities, V varies as $\sqrt{p^n - x^n}$.

571. Cor. 3. If the body descend to the center, or $x = 0$, and the force vary according to any direct law of the distance, or inverse law less than the simple, that is, if n be any affirmative, whole number, or fraction less than unity, $V = \sqrt{\frac{4sF \times p^n}{n \times a^{n-1}}}$ is a real finite quantity. Let the force be as that power of x , whose exponent is $-\frac{1}{2}, 1, 2, 3, \&c.$ or $n = \frac{1}{2}, 2, 3, 4, \&c.$ and $V = \sqrt{\frac{4sF \times p^{\frac{1}{2}}}{\frac{1}{2} \times a^{-\frac{1}{2}}}}, \sqrt{\frac{4sF \times p^2}{2 \times a}}, \sqrt{\frac{4sF \times p^3}{3 \times a^2}}, \sqrt{\frac{4sF \times p^4}{4 \times a^3}}, \&c.$

572. Cor. 4. If the force be inversely as the distance, or $n = 0$, and the body descend to the center, as before, $V = \sqrt{\frac{4sF \times p^0}{0 \times a^{-1}}}$ is infinitely great. If the force be inversely as the square of the distance, or $n = -1$, $F \times y = -\frac{F x^{-2} \dot{x}}{a^{-2}}$ and $Fy = \frac{Fa^2}{x} + \text{cor.} = \frac{Fa^2}{x} - \frac{Fa^2}{p} = \frac{Fa^2 \times \overline{p-x}}{px}$, and $V = \sqrt{4sFa^2 \times \frac{\overline{p-x}}{px}}$. From hence it appears, as before (564), that the velocity varies as $\sqrt{\frac{PO}{OS}}$.

573. PROP. Suppose a body to be impelled by a force varying as that power of the distance from the center of force S , whose exponent is $n-1$, it is required to assign the time of describing any space PO , supposing it to fall from a state of rest.

The fluxion of the space $\dot{S} = V \times \dot{T}$, and consequently \dot{T} is equal to the fluxion of the space divided by the velocity $= -\frac{\dot{x}}{V} =$

$$= \frac{\dot{x} \times \sqrt{n a^{n-1}}}{\sqrt{4 s F \times p^n - x^n}}; \text{ and the time is equal to the fluent of } \frac{\dot{x} \times \sqrt{n a^{n-1}}}{\sqrt{4 s F \times p^n - x^n}}. \text{ Or, if } n-1 \text{ be negative and } = -m, \text{ then}$$

$$\dot{T} = -\dot{x} \times \sqrt{\frac{1-m \times p^{m-1} x^{m-1}}{4 s F a^m \times p^{m-1} - x^{m-1}}}, \text{ and } T = \text{its fluent. Q. E. I.}$$

574. Cor. 1. Let the force vary as the distance from the center of force S , or $n=2$, and $T = \sqrt{\frac{2a}{4sF}} \times$ into the fluent of $\frac{-\dot{x}}{\sqrt{p^2 - x^2}} = \text{cor.} = \text{length of a circular arc, whose radius is unity and right sine } \frac{x}{p}$. If the time vanish, or $x=p$, the cor. = length of a circular arc whose radius is 1 and right sine is $\frac{p}{p}$, or 1, and therefore = a quadrantal arc; consequently the time of describing any space PO is equal to $\sqrt{\frac{2a}{4sF}} \times$ into a quadrant $— \sqrt{\frac{2a}{4sF}} \times$ into the arc whose right sine is $\frac{x}{p}$, the radius being unity. When the body descends to the center, or $\frac{x}{p} = 0$, the time is equal to $\sqrt{\frac{2a}{4sF}} \times$ into a quadrant. If the radius of the circle = r , and a quadrantal arc of that circle = \mathcal{Q} , the quadrant of a circle whose radius is unity = $\frac{\mathcal{Q}}{r}$, and the time of descent to $S = \sqrt{\frac{2a}{4sF}} \times \frac{\mathcal{Q}}{r} = a$ given quantity.

575. Cor. 2. If the force vary inversely as the square of the distance from S , then $V = \sqrt{\frac{4sFa^2 \times p - x}{p x}}$ and $\dot{T} = \dot{x} \times$

F f \sqrt{p}

FIG.
CLXVI.

$\sqrt{\frac{p}{4sFa^2}} \times \sqrt{\frac{x}{p-x}}$, and $T = \sqrt{\frac{p}{4sFa^2}} \times$ into the fluent of $\dot{x} \times \sqrt{\frac{x}{p-x}}$. With SP or p as a diameter, and from the center C de-

scribe a circle, and the fluent of $\dot{x} \times \sqrt{\frac{x}{p-x}}$ is as the area PRS ; for this area = sector PCR + the triangle SRC = $\overline{\text{arc } PR} + \text{fine } OR \times \frac{SP}{4}$, and the fluxion of the area = $\overline{Rr + mr} (P\dot{R} + O\dot{R}) \times \frac{SP}{4}$

= $\frac{\dot{x} \times CR}{OR} + \dot{x} \times \frac{CO}{OR}$ (sim. triang.), (or $\frac{\dot{x} \times x}{OR}$, or $\frac{x \times \dot{x}}{\sqrt{x \times p-x}}$, or $\dot{x} \times$

$\sqrt{\frac{x}{p-x}} \times \frac{SP}{4}$; and consequently the fluent of $\sqrt{\frac{p}{4sFa^2}} \times \dot{x}$

$\sqrt{\frac{x}{p-x}} = \sqrt{\frac{p}{4sFa^2}} \times \overline{\text{arc } PR} + \text{fine } OR$. And because $\frac{p}{4sFa^2}$ is given, the time of describing PO varies as $\overline{PR} + OR$, or as $\overline{PR} + OR \times \frac{SP}{4}$, or as the area PRS .

FIG.
CLXIII.

576. PROF. If a body begin to fall from A , it is required to determine the space AO through which it must descend, when impelled by a force varying as that power of the distance whose exponent is $n-1$, to acquire a velocity, V , equal to that communicated, whilst the body describes a space equal to $\frac{AS}{2}$, by the constant action of a force equal to that at A .

Let the force at $A = 1$

$SA = r$

$SO = x$. And the force at A , or 1 : the force at

$O :: r^{n-1} : x^{n-1}$, and the force at $O = \frac{x^{n-1}}{r^{n-1}}$, and $\frac{r}{2} \times$ the force at A

$$= \frac{r}{2} = \text{the fluent of } -\frac{x^{n-1} \dot{x}}{r^{n-1}} \quad (567) = -\frac{x^n}{n \times r^{n-1}} + \text{cor.} =$$

$$\frac{r^n - x^n}{n \times r^{n-1}}, \text{ and } r^n - x^n = \frac{n r^n}{2}, \text{ and } x = \frac{\sqrt[n]{2 - n}^{\frac{1}{n}} \times r}{2^{\frac{1}{n}}}. \quad \text{Q. E. I.}$$

577. Cor. 1. If the force vary directly as the distance from the center S , or $n = 2$, $x = \frac{\sqrt{2 - 2}^{\frac{1}{2}} \times r}{2^{\frac{1}{2}}} = 0$, and the body must fall to the center to acquire a velocity equal to that acquired in descending through $\frac{SA}{2}$, when acted upon by a constant force equal to that at A .

578. Cor. 2. If the force vary inversely as the distance, or $n = 0$, $x = \frac{\sqrt{2 - 0}^{\frac{1}{0}} \times r}{2^{\frac{1}{0}}}$, which expression does not shew the value of x : but it may be found by the following process. The force at A or 1 : force at $O :: x : r$, and the force at $O = \frac{r}{x}$, and $\frac{r}{2} = r \times \text{flu. } -\frac{\dot{x}}{x} = r \times -\text{hyp. log. of } x + \text{cor.} = r \times -\text{log. of } x + r \times \text{log. of } r = r \times \text{log. of } \frac{r}{x}$.

579. Cor. 3. If the force vary inversely as the square of the distance from S , or $n = -1$; $x = \frac{\sqrt{2 + 1}^{-1} \times r}{2^{-1}} = \frac{2}{3} \times r$, and the body must fall through a third of the distance SA , or to m . If the force vary according to any law therefore between the direct simple and inverse duplicate ratio of the distance from S , the body must fall to some intermediate space between S and m . If the force vary as that power of the distance inversely, whose exponent

ponent is 3, 4, &c. or $n = -2, -3, \&c.$; then $x = \frac{2 + 2)^{-\frac{1}{2}} \times r}{2^{-\frac{1}{2}}}$,
 $\frac{2 + 3)^{-\frac{1}{3}} \times r}{2^{-\frac{1}{3}}}$, &c. $= \frac{2^{\frac{1}{2}} \times r}{2^{\frac{1}{2}}}$, $\frac{2^{\frac{1}{3}} \times r}{5^{\frac{1}{3}}}$, &c.; and the space fallen
 through $= r - \frac{2^{\frac{1}{2}} \times r}{4^{\frac{1}{2}}}$, $r - \frac{2^{\frac{1}{3}} \times r}{5^{\frac{1}{3}}}$, &c.

FIG.
CLXIV.

580. PROP. Suppose the force at A to be 1, and in other places to vary as that power of the distance whose exponent is $n-1$, it is required to determine the height to which a body will ascend when acted upon by this variable force, and projected from A in the direction SA with a velocity equal to that acquired in falling through $\frac{SA}{2}$, as in the last proposition.

Let B be the place to which the body ascends, and the force at A or 1 : force at any other point O :: $r^{n-1} : x^{n-1}$, and the force at O $= \frac{x^{n-1}}{r^{n-1}}$. But if the body were impelled by a constant force equal to that at A, it would ascend through a space equal to $\frac{r}{2}$, and (567) $\frac{r \times 1}{2} =$ the fluent of $\frac{x^{n-1} \dot{x}}{r^{n-1}} = \frac{x^n}{n \times r^{n-1}} + \text{cor.} = \frac{x^n - r^n}{n \times r^{n-1}}$, and $\frac{n \times r^n}{2} = x^n - r^n$, and $x = \frac{n + 2)^{\frac{1}{n}} \times r}{2^{\frac{1}{n}}}$. Q. E. I.

581. Cor. 1. Suppose the force to vary directly as the distance from S, or $n = 2$; then $\frac{n + 2)^{\frac{1}{n}} \times r}{2^{\frac{1}{n}}} = \frac{4^{\frac{1}{2}} \times r}{2^{\frac{1}{2}}}$, and $x (= SB) : r :: 4^{\frac{1}{2}} : 2^{\frac{1}{2}} :: 2^{\frac{1}{2}} : 1$.

582. Cor. 2. If the force vary inversely as the distance from S, or $n = 0$, $x = \frac{0 + 2)^{\frac{1}{0}} \times r}{2^{\frac{1}{0}}}$, which expression does not shew the mag-

magnitude of x ; but it may be determined by the process in (578).

583. Cor. 3. If the force vary inversely as the square of the distance from S , or $n = -1$, then $x = \frac{\sqrt{2-1}^{\frac{1}{-1}} \times r}{\sqrt{2}^{\frac{1}{-1}}} = 2r = SD$.

If the force therefore vary according to any law between the direct simple and inverse duplicate ratio of the distance from S , the body will rise to some intermediate altitude between B and D . If the force vary inversely as the cube of the distance, or $n = -2$,

then $x = \frac{\sqrt{2-2}^{\frac{1}{-2}} \times r}{2^{-\frac{1}{2}}} = \frac{2^{\frac{1}{2}} \times r}{0}$, and is infinitely great. If therefore the force vary according to a law between the inverse duplicate and triplicate, the body will ascend to some intermediate space between D and infinity.

C H A P. XIII.

PENDULOUS MOTION*.

FIG. 584. DEF. *A* BODY, or number of bodies, connected to a right line
CLXV. *SP*, and moving about a point *S* to which it is suspended, by the force of gravity, is called a pendulum.

585. DEF. The motion of a pendulum in the same direction, from a state of rest, till it begins to return in an opposite direction, is one vibration or oscillation.

586. Cor. The velocity acquired in descending from *P* to *V*, supposing *SV* to be perpendicular to the horizon, will make the body *P* describe an arc *Vp*, whose perpendicular altitude is equal to that of *VP* (542), and the motion through *VP* is equal to one half of a vibration. And because a pendulum, composed of any number of bodies, will perform a given vibration, or equal parts of a vibration, in the same time as if they were collected in their center of oscillation (512); a pendulum composed of any number of bodies may be reduced to one, where one body is connected to a right line *SP*, and the length of a pendulum is always understood to mean the distance between the centers of oscillation and suspension.

587. PROP. If a body vibrate in a circular arc *PVp*, that part of the force of gravity, which accelerates and retards it, varies as the right sine of the arc intercepted between the body and the lowest point.

DEM.

* Keil's Phys. Lect. XV. Helsham, Lect. X. Muschenb. Ch. XII. CCXCVII. Mac-laurin's Newt. Book II. Ch. V. Emerson's Mechan. Prop. XL. Hugen. Horol. Oscil. Part II. Prop. XVI. Rohault's Notes, Part II. Ch. XXVIII.

DEM. Let the force of gravity be represented, in quantity and direction, by a given line LP perpendicular to the horizon, and be resolved into two forces, LM parallel to the string SP , and PM touching the circle in P , of which LM is the tension of the string and has no effect upon the body's motion, and it is accelerated and retarded by PM only; therefore the accelerating force (A) : force of gravity (G) :: $PM : PL :: PN : SP$ (sim. triang.), and $A = \frac{G \times PN}{SP}$, and consequently A varies PN . Q. E. D.

588. Cor. 1. The velocity at any point Q is equal to that acquired in falling through the perpendicular altitude ND (541), and varies as \sqrt{ND} , or $\sqrt{NV - VD}$, or $\sqrt{NV \times 2VS - VD \times 2VS}$, or $\sqrt{PV^2 - QV^2}$ (PV and QV being chords of the arcs PV and QV), or as the right sine of a circular arc, whose radius is the chord PV , and versed sine the difference between the chords PV and QV .

589. Cor. 2. If the arc mV be supposed equal to two inclined planes mn and nV , touching it at m and V , these planes are equal, and the velocity in the horizontal plane nV is uniform, and the time of describing it is equal to half the time of falling down mn (527): but the time of falling down mn (t) : time (T) of falling down mV ($= 2mn$), or down the diameter $2SV :: mn : mV :: 1 : 2$, and $t = \frac{T}{2}$; and consequently the times of describing mn and nV , or the time of half a vibration $= \frac{3T}{4}$, and the time of a whole vibration : $T :: 3 : 2$.

SCHOLIUM.

590. If the body P be acted upon by a force F , perpendicular to the horizon, which is to the force of gravity as the arc PV to its sine PN , all vibrations are isochronal; for, let SP represent the constant force of gravity, and $PL = F$; and by a resolution of PL

PENDULOUS MOTION.

PL into two, LM parallel to SP , and PM coincident with the tangent, this last is the only part of F that accelerates P . But $PL : PS :: PM : PN$ (sim. triang.), and

$PL : PS :: PV : PN$ (hypoth.); and consequently $PM = PV$, or the accelerating force is as the distance from the lowest point V , and P will always arrive at V in the same time (563).

The isochronal force in a circle, F , is therefore equal to $G \times \frac{PV}{PN}$; and, when a pendulum is urged by the force of gravity, the time of a vibration will be encreased with the arc of vibration, because the excess of F above G is encreased with that arc.

FIG.
CLXVII.

591. DEF. If a circle FPE revolve upon the right line BA , the curve line BVA described by any point P of the periphery in one revolution, is called a cycloid: BA is the base, VD bisecting BA at right angles is the axis, V the vertex, PM parallel to the base is an ordinate, and FPE is the generating circle, of the cycloid.

592. Cor. Because every point of the periphery of the generating circle has been applied to the base BA , it is evident that BA is equal to the periphery, and BD to the semiperiphery of the generating circle.

593. LEMMA. If a circle be described upon the axis as a diameter, and an ordinate PM be drawn from any point P , the part of this ordinate PL , contained between the point from whence it is drawn and the periphery of the circle, is equal to the circular arc VL , contained between the vertex and the intersection of the ordinate and periphery.

DEM. Let FPE be any position of the generating circle, and, because every point of the arc PE has been applied to BE , the right line $BE =$ the arc $PE =$ the arc LD , and the arc $LV = ED = MN = PL$. Q. E. D.

594. Cor.

FIG. CLVIII.

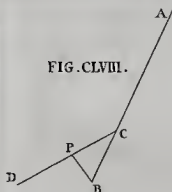


FIG. CLIX.

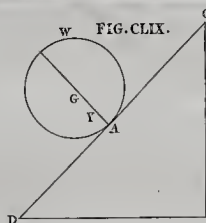


FIG. CLX.

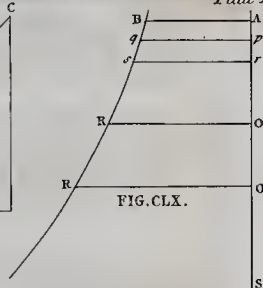


FIG. CLXI.

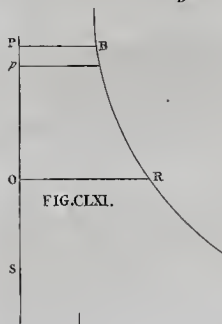


FIG. CLXII.

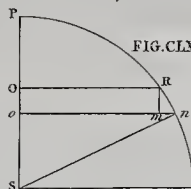


FIG. CLXIV.



FIG. CLXIII.

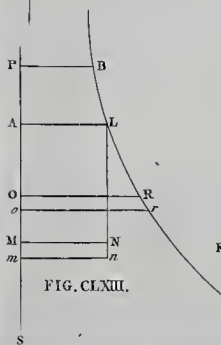
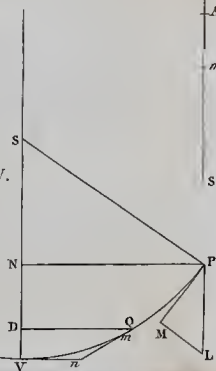


FIG. LXV.



594. Cor. Because PL is always equal to the arc LV or PF , their cotemporary increments or decrements are equal, that is, the initial motions of the point P , which traces out the cycloidal arc, one parallel to the base BA , and the other in the direction of the tangent to the circle at P , are equal to each other.

595. LEMMA. *An ordinate being drawn from any point P of the cycloidal arc, cutting the periphery of the generating circle, whose diameter is the axis VD , in L , the chord VL of the circular arc is parallel to the line touching the cycloid at P .*

DEM. Let GP be a tangent to the circle at P , and producing EP to C , the $\angle CPq = \angle GPE = \angle GEP$ (the tangents DP and DE being equal, EUC. B. III. p. 36.) $= \angle EPN$: but the initial motions of P , in the directions GP and PL , being equal (594), the path of P or Pb will bisect the angle qPp (composition of motion), and consequently the $\angle CPb = \angle EPb$ or Pb is at right angles to EP , and FP , which is parallel to LV , is also perpendicular to EP , and therefore is a tangent to the curve at P . Q. E. D.

596. Cor. A tangent to the cycloid at the vertex V is therefore perpendicular to the axis VD and parallel to the base BA .

597. LEMMA. *The generating circle being described upon the axis as a diameter, and an ordinate being drawn from any point O , cutting the periphery of the circle in R , the cycloidal arc VO is equal to twice the corresponding chord of the circular arc VR .*

DEM. Draw an ordinate ro infinitely near to RO , and rs perpendicular to VR produced, and $Vr = Vs$, the angle at V being evanescent; and $Oo (= Rv)$ and Rr are cotemporary increments of the arc VO and chord VR : drawing tangents to the circle at V and R , the triangles VRT and Rrv , having the angles at R vertical,

tical, and the angles TVR and Rvr alternate, are similar, and consequently $Rr = vr$ and Rv or $Oo = 2Rs$. The increment of the arc VO is therefore equal to twice the corresponding increment of the chord VR , and, because they are nascent and evanescent at the same time, $VO = 2VR$. Q. E. D.

FIG.
CLXVIII.

598. PROP. *To make a body oscillate in a given cycloid AVB.*

Produce the axis VD making $SD = DV$, and through S draw a line KSH parallel to BA , and let two circles, each equal to DLV , generate two femicycloids SA , SB , each equal to BV or AV ; and if one extremity of a string SCX , whose length is equal to SV or SCA , be fixed at S , a body collected in the other extremity X , will always be found in the cycloidal arc XVP ; for, because SCX is a tangent to the cycloid at C , CX is parallel to EA (595), and $CG = EA$; but $CX = 2AE$ (597) $= 2CG$ and $CG = GX$, and the ordinates XL and CE are equidistant from AD , or $AF = DN$, and the arc $AE =$ the arc LD , and the chord LD is parallel to AE or XG : but $AG = CE =$ the arc AE , and consequently the remaining arc $HE(=LV) = GD = XL$, and X is in the cycloidal arc AXV (593).

599. Cor. Because SCX touches the curve at C , it is evident that whilst X describes a very small arc, C may be deemed quiescent, or CX is always perpendicular to the cycloidal arc at X ; and very near the vertex V , where SV is perpendicular to the curve, a circular arc, whose radius is SV , will coincide for a small distance with the cycloidal arc.

600. PROP. *If a pendulum begin to vibrate from any point P, and a circle be described, whose radius is equal to the cycloidal arc VP, the velocity in different points will vary as the right sine of an arc of this circle, whose versed sine is the space described by P.*

DEM.

DEM. The velocity at any point \mathcal{Q} is equal to that acquired in falling through the same perpendicular altitude NR (545), which is as \sqrt{NR} , or $\sqrt{NV - RV}$, or $\sqrt{NV \times VD - VR \times VD}$, or $\sqrt{VF^2 - VE^2}$, or $\sqrt{VP^2 - V\mathcal{Q}^2}$, or $\sqrt{Vn^2 - Vq^2}$, or qn . Q.E.D.

601. Cor. 1. The velocity therefore encreases from P to V , where it is the greatest, being as Vv , and then it decreases in such a manner that the velocities at all equal distances from V are equal; and consequently the cotemporary changes, or increments as the body accedes to V , and decrements as it recedes from it, are equal at equal distances from V .

602. Cor. 2. If s be the number of inches through which a body descends from rest by the force of gravity in 1", and V be the velocity acquired in vibrating through $P\mathcal{Q}$, or in falling down the perpendicular NR , then $NR : s$ as the squares of the velocities acquired in falling through NR and s , or as $V^2 : 4s^2$, and $V = \sqrt{NR \times 4s}$

$$= \sqrt{4s \times NV - VR} = \sqrt{\frac{4s}{DV} \times \sqrt{FV^2 - EV^2}} = \sqrt{\frac{4s}{DV} \times}$$

into the right sine $\frac{qn}{2}$; and at the vertex the velocity is equal to

$$\sqrt{\frac{4s}{DV} \times \frac{Vv}{2}} \text{ inches in 1"}. \quad \text{Q.E.D.}$$

603. Cor. 3. If a pendulum begin to oscillate at different distances PV and $\mathcal{Q}V$ from the vertex V , the velocities acquired at V in these different vibrations, are as $\sqrt{NV} : \sqrt{RV}$, or as $VF : VE$, or as $VP : V\mathcal{Q}$; or the velocity acquired in vibrating from rest is as the distance from the vertex at which it begins to vibrate.

604. Cor. 4. If s be the space described from rest, as in cor. 2. the velocities acquired at V , in vibrating from different points P

PENDULOUS MOTION.

and \mathcal{Q} are equal to $\sqrt{\frac{4s}{DV}} \times \frac{VP}{2}$, and $\sqrt{\frac{4s}{VD}} \times \frac{qn}{2}$, or such as would carry the body over these numbers of inches in 1".

605. PROP. *The time in which the pendulum vibrates through any arc QH, is equal to the time in which a body would describe the corresponding circular arc no with the greatest velocity Vv, continued uniformly.*

Let M be infinitely near to \mathcal{Q} , so that $\mathcal{Q}M$ may be described uniformly, and taking Vp, Vq, Vm, Vh , respectively equal to $VP, V\mathcal{Q}, VM, VH$, and drawing the right line nk parallel to Vp ; nk (qm or $\mathcal{Q}M$) : nl :: qn : Vv , or $\mathcal{Q}M$ and nl are to each other as the velocities with which they are described, and are therefore described in the same time. The same may be proved of the other corresponding parts of $\mathcal{Q}H$ and no , which are, consequently, described in the same time. Q. E. D.

606. Cor. 1. The times of describing any arcs $P\mathcal{Q}, \mathcal{Q}V$ are to each other as the corresponding circular arcs pn, nv , because they are described with the same uniform velocity; and the time of a whole vibration is as the semiperiphery pvx .

607. Cor. 2. The velocity, with which $\mathcal{Q}M$ is described, is actually equal to $\sqrt{\frac{4s}{DV}} \times \frac{qn}{2}$, and the velocity with which nl is described, is equal to $\sqrt{\frac{4s}{DV}} \times \frac{Vp}{2}$. The time of a vibration is equal to the time in which a body would describe the semiperiphery pvx with an uniform velocity equal to $\sqrt{\frac{4s}{DV}} \times \frac{Vv}{2}$ inches in a second.

608. Cor.

608. Cor. 3. If a body begin to oscillate from different points P and Q , the times of their vibrations are equal to the times of describing the femiperipheries of circles, whose radii are VP and VQ , with uniform velocities, which are respectively equal to $\sqrt{\frac{4s}{DV}} \times \frac{VP}{2}$, and $\sqrt{\frac{4s}{DV}} \times \frac{VQ}{2}$; but these times are as the spaces or femiperipheries divided by the velocities, or divided by their radii VP and VQ , and are consequently equal: the times therefore of all vibrations, however different, are equal.

609. PROP. Supposing a pendulum to begin to oscillate from any point P , the time in which it performs one vibration, is to the time of descent down the axis as the periphery of a circle to its diameter.

DEM. The times of descent down DV and FV are equal (538), and in this time a body would describe a space equal to $2FV$ or PV or pV , with the velocity acquired in falling down FV or PV continued uniform; but this is to the time of describing pvx with the same velocity, or to the time of one vibration (607) as the spaces described, or as $pV : pvx$, or as the diameter of a circle to its periphery. Q. E. D.

610. Cor. 1. The periphery of a circle, being to its diameter in a given ratio, it appears again that the times of all vibrations are as the times of descent down the axis, and consequently are given.

611. Cor. 2. The velocity with which the femiperiphery pvx is described, when the time of describing it is equal to a vibration,

is equal to $\sqrt{\frac{4s}{DV}} \times \frac{PV}{2}$ (607), and the time is equal to $\frac{pvx}{\text{vel.}} =$

$\frac{2pvx}{PV} \times \sqrt{\frac{DV}{4s}}$; and the time of falling down $DV = \sqrt{\frac{DV}{s}}$; con-

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consequently the time of a vibration is to the time of falling down

$DV :: \frac{2 \times p v x}{PV} \times \sqrt{\frac{DV}{4s}} : \sqrt{\frac{DV}{s}} :: \frac{2 \times p v x}{PV} \times \frac{1}{\sqrt{4}} : 1 :: 2 p v x : 2 PV$, the same analogy as that in this proposition.

612. Cor. 3. The times of vibrations in different cycloids, being equal to the times of descent down their axes multiplied into the same given quantity, will vary as the times of descent down the axes, or lengths of the pendulums; or supposing L to be the length of a pendulum, and T the time of a vibration, and to be variable; T will be as \sqrt{L} , and L as T^2 .

613. Cor. 4. If the vibrations be very small, the string is not sensibly affected by the cycloidal arcs SA , SB , and the pendulum SP will describe a circular arc, and hence it appears that very small vibrations in circular arcs are performed in equal times. The times of descent down a circular arc and its chord are therefore unequal, the former being to the time of descent down the axis as half the periphery of a circle to its diameter; and the latter is equal to the time of descent down the diameter, or four times the axis, and is consequently to the time of descent down the axis as 2 : 1.

614. Cor. 5. If the pendulum begin to oscillate from B , half the time of a vibration is to the time of descent down the inclined plane BV as $BD : BV$; for half the time of a vibration is to the time of descent down the axis, as the semiperiphery of a circle to its diameter, or as $DFEV(BD)$ is to DV , and the time down DV : time down the inclined plane $BV :: DV : BV$, and, ex æquo, half the time of a vibration : time of descent down the inclined plane $BV :: BD : BV$.

615. Cor. 6. The space described by a falling body in 1", may be discovered, if the length of a pendulum performing one vibration

tion in a second be known: let the length of this pendulum be 39.2 inches, and (609) $1''$: time of descent through $\frac{39.2}{2} (t) ::$ periphery of a circle to its diameter $:: 3.14159, \&c. (p) : 1$ and $t = \frac{1}{p}$. But $\frac{39.2}{2} : \text{space fallen through in } 1'' (x) :: \frac{1}{p^2} : 1^2$ and $x = \frac{39.2}{2} \times p^2 = 193.1$ inches nearly $= 16.1$ feet nearly.

616. Cor. 7. The length of the pendulum performing one vibration in $1''$, may be discovered, if the space through which a body descends in $1''$ be known; for let this space $= 193.1$ inches, and (last cor.) the length of the pendulum $= \frac{2 \times 193.1}{3.14159}$ inches $= 39.2$ inches.

617. Cor. 8. The time of one vibration may be discovered, if the length of the pendulum be known; for let L be the given length of the pendulum in inches, and T be the time of one vibration, and T : time of descent through $\frac{L}{2} :: 3.14159, \&c. : 1 :: p : 1$; but the time of descent through $\frac{L}{2} = \sqrt{\frac{L}{2s}}$, and consequently $T = p \times \sqrt{\frac{L}{2s}}$ seconds, or parts of seconds. If $L = 39.2$ inches, and $s = 16.1 \times 12$ inches, $T = 3.14159 \times \sqrt{\frac{39.2}{2 \times 16.1 \times 12}} = 1''$ nearly.

618. Cor. 9. The number of vibrations in a given time, being inversely as the time of one vibration, will be as $\frac{1}{\sqrt{L}}$. If N be the

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the number of vibrations in a given time T , and t be the time of one vibration, then $N = \frac{T}{t}$; but, the length of the pendulum being L , $t = p \times \sqrt{\frac{L}{2s}}$, and consequently $N = \frac{T}{p} \times \sqrt{\frac{2s}{L}}$. If $T = 60 \times 60''$, then $N = \frac{60 \times 60}{3.14159} \times \sqrt{\frac{193 \times 2}{39.2}} = 60 \times 60$ very nearly.

619. PROP. *Let a pendulum, whose length is L inches, lose or gain any number of vibrations expressed by n , in any number of hours, h , to find the length of a pendulum that shall vibrate once in 1".*

Let Y = the length of the pendulum required, and let the number of vibrations performed by Y in h hours, or $h \times 60 \times 60'' = m$, and L performs $m \pm n$, vibrations in h hours; but $L : Y :: m^2 : m \pm n^2$ (618) and $Y = L \times \frac{m^2}{m \pm n^2}$.

620. PROP. *If the length of a pendulum and the force of gravity be variable, the time of an oscillation varies as the square root of the length of the pendulum L , divided by the square root of the force of gravity, G .*

DEM. The time of one oscillation varies as the time of descent down the axis, and that is as the square root of the axis directly and force inversely, or as $\sqrt{\frac{DV}{G}}$ (532), or as $\sqrt{\frac{L}{G}}$. Q. E. D.

621. Cor. 1. Therefore L is as $T^2 \times G$, and T being given, L is as G : if therefore the lengths of two pendulums, performing a vibration, or the same number of vibrations, in the same time be known,

PENDULOUS MOTION.

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known, the ratio of the forces of gravity, being the same with the ratio of these lengths, will consequently be known.

622. Cor. 2. Or, if the number of vibrations performed in the same time, in different places, by two pendulums whose lengths are L and l , be equal to m and $m + n$ respectively, the forces of gravity G and g , in the places of observation, may be found: for, the times of one vibration of the pendulums L and l are as $\sqrt{\frac{L}{G}}$:

$\sqrt{\frac{l}{g}}$, and $m + n : m$ inversely as the times of one vibration, or as $\sqrt{\frac{L}{G}} : \sqrt{\frac{l}{g}}$ (620), and consequently $G : g :: m^2 \times L : (m + n)^2 \times l$.

623. PROP. *That part of the force of gravity, which accelerates or retards a pendulum vibrating in a cycloid, varies directly as its distance from the vertex.*

DEM. Let the force of gravity be represented, in quantity and direction, by the line DV , and be resolved into two forces, DF parallel to the string, and VF parallel to the tangent at P , or direction in which P moves; and $VF (= 2VP)$, which only accelerates P , varies as VP . Q. E. D.

Otherwise; The velocities at Q and M are as qn , ml ; the time of describing QM is as the included circular arc nl ; and, supposing the time nl to be very small and given, the accelerating force at Q is as the change of velocity kl (453), or as $\frac{Vq \times nl}{Vn}$ (sim. triangles), or, because nl and Vn are given, as Vq or VQ . Q. E. D.

624. Cor. 1. The force accelerating or retarding the pendulum (A) : whole force of gravity (G) :: $VF : VD :: VP : VB$ and $A = \frac{G \times VP}{Hh}$

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$\frac{G \times VP}{VB}$; and the tension of the string : $G :: DF : DV$, and this

tension $= \frac{G \times DF}{DV}$. At the vertex where $DF = DV$ the tension of the string $= G$, and $A = 0$; and, at B , $VP = VB$ and $A = G$, and the tension $= 0$.

625. Cor. 2. The tension of the string : force accelerating or retarding $P :: DF : VF :: \sqrt{DN} : \sqrt{VN}$ (sim. triangles).

CHAP.

C H A P. XIV.

PROJECTILE MOTION.

626. DEF. *THE range or random of a projectile, is the rectilineal distance between the point of projection and impulse upon any obstacle; the horizontal range is called the amplitude; and the angle of elevation is the angle contained between the horizon and the direction of projection.*

627. PROP. *A body projected in any direction DN, and acted upon by a constant force, G, in a direction parallel to a right line DO inclined in any angle to DN, will describe a parabola.*

FIG.
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DEM. Let *DO* be described from rest, by the action of *G*, in the same time, *T*, in which *DE* is described by the velocity of projection, and, completing the parallelogram, the body will, at the end of the time *T*, evidently be at *R*; but *DO* varies as T^2 (526), or (because the velocity in *DN* is uniform from the first law of motion, and *T* varies as *DE*) as DE^2 or OR^2 , which is a property peculiar to the common parabola. Q. E. D.

628. Cor. 1. The parameter belonging to any diameter *DO* is equal to $\frac{OR^2}{DO}$ or $\frac{DE^2}{ER}$ (conic sect.); and if the parameter belonging to any point *D* be equal to $\frac{DE^2}{ER}$, the parabola will pass through the point *R*.

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629. Cor. 2. If the velocity of projection be the same, the parameter belonging to DO will be the same, whatever be the angle of projection; for the projectile velocity and G being given, DE , ER , and consequently $\frac{DE^2}{ER}$, or the parameter, are given.

630. Cor. 3. Let Y be the space through which a body falls from rest, when acted upon by the constant force G , to acquire a velocity, V , equal to that in any point D of the curve, and let v be the velocity acquired in falling from E to R , by the action of same force G , and $V : v :: DE : 2ER$ (106): but $Y : ER :: V^2 : v^2$ (526) $:: DE^2 : 4ER^2$, and $Y = \frac{DE^2}{4ER} = \frac{1}{4}$ of the parameter belonging to DO .

631. Cor. 4. If G be the force of gravity, the axis and all diameters of the parabola, are perpendicular to the horizon; the velocities are equal at equal distances from the principal vertex; the time of arriving at the vertex of the axis, or at the greatest altitude, is equal to half the time of the flight; the parabola cuts the horizon, and all lines parallel to the horizon, in equal angles; and DE , which is parallel to the ordinate OR , is a tangent to the parabola at D .

FIG.
CLXX.

632. Cor. 5. The velocity in different points varies in a subduplicate ratio of the parameter belonging to those points, or, if IX be the directrix and DI perpendicular to it, in a subduplicate ratio of DI or DS , S being the focus, or, in the same parabola, as the perpendicular SY upon the tangent DT . The horizontal velocity DH , determined by drawing perpendiculars from the body upon the horizon, varies as DE , and is consequently uniform; and the velocity, therefore, in any point D , being as DE , is as the secant of the angle of elevation; and at the principal vertex, the velocity in the curve and the horizontal velocity, are equal.

633. PROP.

633. PROP. To find the directions in which a body, being projected with any velocity, V , from a given point D , will pass through any point P .

FIG.
CLXXI.

Because V , or the number of feet uniformly described in 1", with the velocity of projection, is given, the space DE , described from rest to acquire this velocity, is known, being equal to $\frac{V^2}{64}$ (art. 529). Draw DA , equal to $4DE$, perpendicular to the horizon, BC perpendicular to DA through its bisection G , and DC perpendicular to DP ; and from C as a center, at the distance CD , describe a circle, and, if R, r , be the intersections of its periphery with a right line, passing through P perpendicular to the horizon, DR , and Dr will be the directions required; for (sim. triang.) $DA : DR :: DR : RP$, and $DA : Dr :: Dr : rP$, and DA (= the parameter) = $\frac{DR^2}{RP}$ or $\frac{Dr^2}{rP}$; and the parabola described by a body projected in either of the directions DR or Dr with the velocity V , will pass through P (628). Q. E. I.

634. Cor. 1. Because the arcs BR and Br are equal, the line DB makes equal angles with the corresponding directions DR and Dr . Drawing BQ a tangent to the circle at B , a body projected in the direction DB , will pass through Q , and DQ is the greatest distance upon that line to which the body can be projected with the velocity V .

635. Cor. 2. If DP be horizontal, the range upon the horizon = $DF = RN$, the sine of twice the angle of elevation FDR ; for, by the property of the circle $\angle FDR = \angle RAD = 2 \angle RCD$. The horizontal range is therefore the greatest, when the angle of elevation is 45° , the sine of twice that angle, or of 90° , being equal to $\frac{1}{2}DA$, or $\frac{1}{2}$ of the parameter at D ; or it is equal to the latus rectum, for that is always equal to $\frac{DM^2}{MV}$, which, in this case, is equal to $\frac{DC^2}{4 \times \frac{1}{4}DC} = DC$.

FIG.
CLXXII.

Because

PROJECTILE MOTION.

Because the sine of any angle varies as the radius or diameter, if the circle be described upon any other diameter DE or $\frac{1}{4}$ th of DA , the horizontal range will be equal to the sine of twice the angle of elevation of this circle multiplied into $\frac{DA}{DE}$, or into 4.

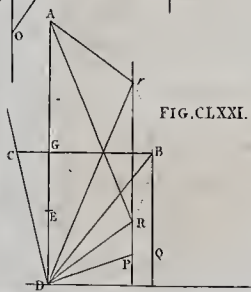
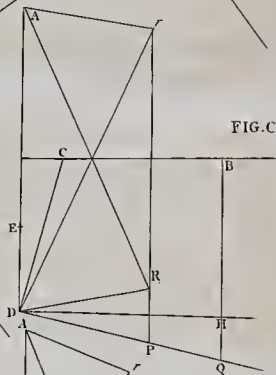
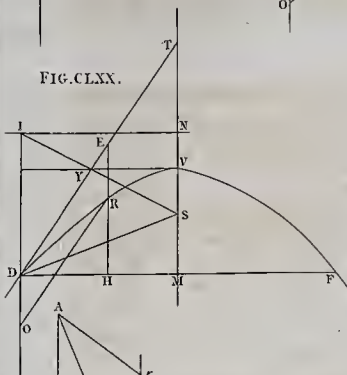
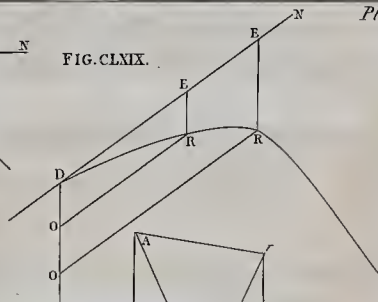
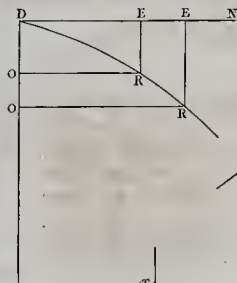
636. Cor. 3. If the axis MV , produced, be intersected by the tangent in T , the greatest altitude $MV = \frac{1}{2}MT$ (conic sect.) $= \frac{1}{4}FR = \frac{1}{4}$ th of the versed sine of twice the angle of elevation; and when this angle is right, $\frac{1}{4}$ th of the versed sine $= \frac{1}{4}DA$. Because the versed sines of any angles are as the diameters of the circles described, if the circle be described with a diameter equal to DE , the altitude is equal to the versed sine of twice the angle of elevation of this circle.

637. Cor. 4. The time of the flight is equal to the time of describing DR with the projectile velocity, and varies as DR , or as $\frac{1}{2}DR$, which is the sine of the angle of elevation; and consequently the time is the greatest when the angle $FRD = 90^\circ$, and it is then equal to the time of descent from rest through the parameter AD . The time is also equal to the time of falling from rest through RF , and therefore varies as \sqrt{RF} , or as the square root of the versed sine of twice the angle of elevation, or as the square root of $\frac{1}{4}RF$, or the greatest altitude.

638. PROP. *The velocity of projection, V , and the angle of elevation being given, to describe the path of the projectile.*

FIG.
CLXXIII.

Let V be the number of feet uniformly described in 1", and $\frac{V^2}{64} = DE = \frac{1}{4}$ th of the parameter, is the same whatever be the angle of elevation, and the periphery of a circle, whose radius is DE , and center D , will pass through the foci of all parabolas described by a body projected from D with the velocity V . Let DE be perpendicular to the horizon, and DR the direction of pro-



projection, and making the angle $SDR = \angle RDE$; S will be the focus: EX drawn perpendicular to DE is the directrix; V the bisection of SX drawn parallel to DE , is the principal vertex, and SV the axis; DM perpendicular to VM is an ordinate; SE bisects DT in R and is perpendicular to it, and RV is a tangent to the parabola at V (conic sect.).

639. Cor. Again, it appears that the amplitude DF is (as in art. 635) equal to four times the sine of twice the angle of elevation, for $\angle MDR = \angle RDC = 2\angle RCN$, and $DF = 2DM = 2NV = 4RN$.

640. PROP. To find the least velocity with which a body projected, from a given point D , will hit a given point P .

FIG.
CLXXIV.

Draw BPH perpendicular to the horizon, and, making PB equal to PD joined, bisect the angle DPB by the right line PC , meeting DC perpendicular to DP , in C ; and, because the triangles CBP and CDP are equal and similar, $CB = CD$, the \angle s CBP and CDP are right angles, and BP is a tangent to the circle described from C , as a center, at the distance CB . Bisect DG , drawn perpendicular to the horizon, in E , and a body projected in the direction DB , with a velocity equal to that acquired in falling through ED , will pass through P (634), and, because DP is the greatest range (634), the velocity must be the least possible. Q. E. I.

641. Cor. If P be in the horizon, CD and GD coincide, and the angle of elevation BDH is equal to 45° .

642. PROP. The range upon any plane DP , is equal to the product of the parameter and sines of the angles formed by the plane and direction of projection, and the plane and a perpendicular to the horizon, divided by the square of the cosine of the plane's elevation.

FIG.
CLXXI.

DEM.

PROJECTILE MOTION.

DEM. From trigon. $DP:DR::\sin.\angle DRP$ or $\angle ADR:\sin.\angle DPF$,
 $DR:DA::\sin.\angle DAR$ or $\angle PDR:\sin.ARD$ or $\angle RPD$ or DPF ;
 and consequently $DP:DA::\sin.\angle ADR\times\sin.\angle PDR:\sin.^2\angle DPF$ or
 $\cos.^2\angle PDF$, and $DP = \frac{DA \times \sin.\angle ADR \times \sin.\angle PDR}{\cos.^2\angle PDF}$.

Q. E. D.

FIG.
CLXXV.

643. Cor. 1. Let S = the sine of the angle contained between the plane and the direction; \mathcal{Q} = the sine of the angle contained between the direction and the parameter DA perpendicular to the horizon; and C = the cosine of the angle contained between the plane and the horizon; and the range $DP = \frac{DA \times S \times \mathcal{Q}}{C^2}$, and, because DA varies as the square of the velocity, or V^2 , DP is as $\frac{V^2 \times S \times \mathcal{Q}}{C^2}$.

This corollary may be demonstrated differently by the following process: let V be the velocity of projection, or number of feet described uniformly in 1", and $V:DR::1":\frac{DR'}{V}$ = the time of describing DR or RP , and $RP = 16 \times \frac{DR^2}{V^2}$ (529) = $DP \times \frac{S}{\mathcal{Q}}$; but $DR = DP \times \frac{C}{\mathcal{Q}}$, therefore $DP \times \frac{S}{\mathcal{Q}} = 16 \times \frac{DP^2 \times C^2}{V^2 \times \mathcal{Q}^2}$, and consequently $DP = \frac{S \times \mathcal{Q} \times V^2}{16 \times C^2}$.

644. Cor. 2. If DP be bisected in M and MV be drawn perpendicular to the horizon, the greatest altitude $MV = \frac{1}{2}MT = \frac{1}{4}PR$; but $PR:DP::S:\mathcal{Q}$ and $PR = \frac{DP \times S}{\mathcal{Q}}$ = (substituting the value of DP) $\frac{V^2 \times S^2}{C^2}$; and consequently MV , or the greatest altitude varies as $\frac{V^2 \times S^2}{C^2}$.

645. Cor.

PROJECTILE MOTION.

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FIG.
CLXXV.

645. Cor. 3. The time of the flight is equal to the time of describing DR , and is as $\frac{DR}{\text{vel.}}$; or, if $V =$ the number of feet described uniformly in 1" with the velocity of projection, the time is actually equal to $\frac{DR''}{V} = (\text{fig. 171.}) \frac{DA \times S}{C \times V}$; and because DA is as V^2 , the time varies as $\frac{V \times S}{C}$. The time is also equal to the time of descent from rest through RP , and consequently varies as \sqrt{PR} , or as \sqrt{MV} .

646. Cor. 4. Supposing V to be given, the range DP is the greatest when $S = 2$, or when the direction bisects the angle contained between the plane and the vertical DA ; and if the plane DP coincide with the horizon, the amplitude $DP = DF (= \frac{DA \times S \times 2}{C^2}) = \frac{DA \times FR \times DF}{DR^2}$ (making DR the radius) $= DF$ or RN the sine of twice the angle of elevation, DA being equal to $\frac{DR^2}{FR}$, and $\frac{DA \times FR}{DR^2}$ being given. The greatest altitude $= \frac{DA \times S^2}{4C^2} = \frac{DA \times FR^2}{4DR^2} = \frac{1}{4}FR$.

647. PROP. If both the velocity of projection, V , and the angle of elevation, vary, the horizontal range or amplitude will vary as the square of the velocity and sine of twice the angle of elevation.

DEM. The range, being generally as $\frac{V^2 \times S \times 2}{C^2}$ (643), will, when DP coincides with the horizon, be as $\frac{V^2 \times S \times 2}{\text{rad.}^2}$; and, if the radius be unity, $S \times 2 = \frac{1}{2}$ of the sine of twice the angle of elevation (trigonom.), and the amplitude consequently is as $V^2 \times$ sine of twice that angle. Q. E. D.

FIG.
CLXXVI.

I i

Another

PROJECTILE MOTION.

Another demonstration:

FIG.
CLXXVII.

Let V = the number of feet described uniformly in 1", by the velocity of projection, and $V : DL :: 1'' : \frac{DL''}{V}$, or the time of describing DL with the velocity of projection, which is equal to the time of describing LF from rest, and $LF = \frac{s \times DL^2''}{V^2}$, supposing s to be the descent from rest in the first second (529); but, describing a circle with any diameter DA , $LF = \frac{DL \times DR}{DA}$ (sim. triang.), and consequently $\frac{s \times DL^2}{V^2} = \frac{DL \times DR}{DA}$, and $DL = \frac{DR \times V^2}{s \times DA}$; and $DF (= \frac{DL \times RN}{DR} \text{ from sim. triang.}) = \frac{RN \times V^2}{s \times DA}$, and varies as $V^2 \times RN$, because s and DA are given. Q. E. D.

648. Cor. The greatest altitude is generally as $\frac{V^2 \times S^2}{C^2}$ (644), or when DP is horizontal, as $\frac{V^2 \times S^2}{\text{rad.}^2}$, or, the radius being unity, as $V^2 \times \text{versed sine of twice the angle of elevation}$; for $S^2 = \frac{1}{2}$ of this versed sine.* This corollary is also deducible from the second demonstration; for the greatest altitude $= \frac{LF}{4} = \frac{DL \times DR}{4DA} =$ (substituting the value of DL) $\frac{DR^2 \times V^2}{4s \times DA^2} = \frac{V^2 \times DN}{4s \times DA}$, and consequently varies as $V^2 \times DN$, s and DA being given.

649. PROP.

FIG.
CLXXVI.

* Let the arcs DF , EF , and the angles DCF and ECF , be equal, and, FCD being the angle of elevation, EL drawn perpendicular to the radius CD , is the sine, and DL the versed sine, of twice that angle; and, joining ED and drawing GH from G , the intersection of CF and ED , parallel to DL , $EH = \frac{1}{2}$ the sine, and $GH = \frac{1}{2}$ the versed sine. But the triangles EGH and ECG are similar, the angles at H and G being right, and the $\angle EGH = \angle EDC = \angle CED$; therefore $EG (= FN = S) : EH :: CE (= 1) : CG (= CN = 2)$, and $EH = S \times 2$. And EG or $S : GH (= \frac{1}{2} LD) :: CE : EG$, and consequently $GH = S^2$.

649. PROP. *The velocity of projection, V, or number of feet described uniformly in 1" with that velocity, and the angle of elevation, E, being given, to find the amplitude, altitude, and time of flight.*

1. The space described from rest to acquire the velocity of projection is equal to $\frac{V^2}{4s}$, and the amplitude, when the angle of elevation $= 45^\circ$, is equal to $\frac{V^2}{2s}$ (635); but the amplitudes are as the times of twice the angles of elevation, and consequently $\sin. \text{ of } 90^\circ : \sin. \text{ of } 2E :: \frac{V^2}{2s}$, or $\frac{1}{2}$ parameter : amplitude, which is therefore known.

2. When the angle of elevation $= 45^\circ$, the altitude $= \frac{1}{2}$ of the parameter $= \frac{V^2}{8s}$, and, the altitude being always as the versed sine of twice the angle of elevation, the versed sine of $90^\circ : \text{versed sine of } 2E :: \frac{V^2}{8s} : \text{altitude}$, which is therefore known.

3. When the angle of elevation is 90° , the time of the flight is equal to the time of falling, from rest, down the parameter, $\frac{V^2}{s}$, and $= \sqrt{\frac{V^2}{s^2}} = \frac{V}{s}$, and the times of flight being as the sines of angles of elevation (637), the sine of $90^\circ : \sin. \text{ of } E :: \frac{V}{s} : \text{time of the flight}$, which is therefore known. Q. E. I.

650. PROP. *The amplitude and the angle of elevation, E, being given, to find the velocity of projection, and the altitude.*

PROJECTILE MOTION.

1. Let the parameter $= P$, and, when the angle of elevation is 45° , the amplitude $= \frac{P}{2}$; therefore, the $\sin.$ of $2E : \sin.$ of $90^\circ ::$ given amplitude $: \frac{P}{2}$, which is therefore known; and the velocity of projection, being equal to that acquired in falling from rest through $\frac{P}{4}$, is also known, being equal to $\sqrt{4s \times \frac{P}{4}}$ feet in 1" (529).

2. The altitude is found when the parameter is known, by the last proposition.

FIG.
CLXXVII.

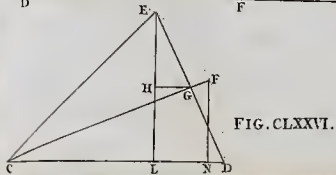
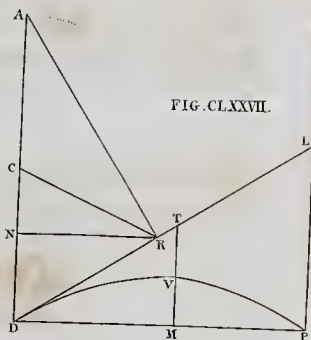
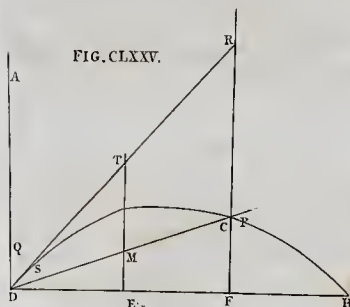
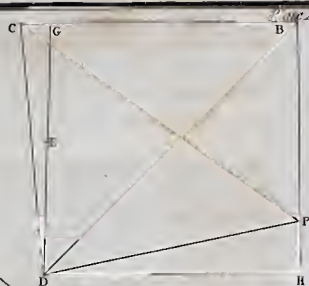
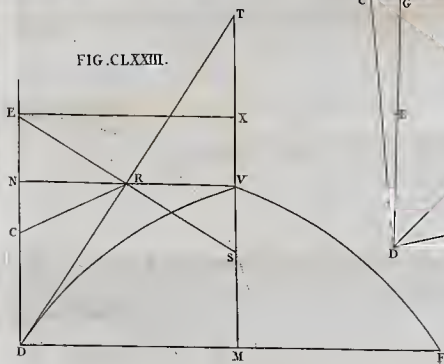
Or, The altitude is easily found by the following process: Bisect the amplitude DE in M , and, drawing MT perpendicular to DE , the angles in the triangle DMT , and the side DM , are known, and MT may be found from trigonometry; and consequently $MV = \frac{1}{2}MT$ is known. Q. E. I.

SCHOLIUM.

651. By a similar process the converse of these propositions are easily solved, that is, 1. the amplitude and altitude being given, to find the velocity of projection and angle of elevation; 2. the velocity of projection and altitude being given, to find the amplitude and angle of elevation; and, 3. and the angle of elevation and altitude being given, to find the amplitude and velocity.

SCHOLIUM II.

652. In the theory of projectiles, the medium is supposed to be void of resistance, which supposition does not obtain in practice, the resistance of the air being variable according to the different velocities and magnitudes of the projectiles, and always considerable.



able. 1. If a musket ball, of $\frac{3}{4}$ of an inch diameter, be fired, from a piece forty-five inches long, with half its weight of powder, its velocity is nearly 1700 feet in 1", and its horizontal range, when the angle of elevation is 45° , ought to be about seventeen miles; but from practical writers it appears, that the range is less than half a mile. An iron ball of 24 lb. weight, discharged with a full charge of powder, has a velocity of 1650 feet in 1", and its horizontal range at 45° would be about sixteen miles, were the path described a parabola; but, from experiments, it appears not to be three miles, and not $\frac{1}{5}$ th part of the space investigated by the theory. When the velocity of the shot is about 400 feet in 1", the resistance of the air is still considerable, and neither the amplitude, height or time of flight, correspond with the theory. This resistance is confirmed by the constant observation of all conversant in the projection of bombs; for the ranges at elevations equally distant from 45° , ought, according to the theory, to be equal; but the ranges of a shell projected at an elevation of 15° or 20° , are always found to be greater than those projected at elevations equal to 60° or 65° , though equally distant from 45° .

Projectiles, whose motion is sensible to the eye, are seen to descend in a curve obviously shorter, and inclined in a greater angle to the horizon, than that in which they ascended; and whoever views, in a proper situation, the flight of stones, arrows, shells, &c. projected to any considerable distance, may see evidently that the vertex of the curve or greatest altitude, divides the path described into two unequal parts, and is more remote from the point of projection, than where the projectile falls to the ground.

SCHOLIUM III.

653. A body, projected with a considerable velocity, is often not only affected by the resistance of the air, and deflected from a parabolic path in a direction perpendicular to the horizon, but is made to deviate laterally and change the plane of motion. The path of a tennis-ball, struck with great force, is plainly observed to be incurvated sideways as well as downwards, and to move in a different plane from that arising from the combined.

PROJECTILE MOTION.

bined action of gravity and force of projection. Bullets are not only depressed beneath the line of projection, but deflected to the right or left of that direction by the resistance of the air, or action of some other force. Mr. Robins fixed a barrel, carrying a ball of $\frac{3}{4}$ th of an inch diameter, and, firing at a mark, one foot and $\frac{1}{7}$ th square, at the distance of 180 feet, missed it only once in sixteen successive trials; but when the same barrel was fired with a smaller quantity of powder, he found the ball to be deflected 100 yards to the right or left of the line aimed at and placed at the distance of 760 yards; and its direction in the perpendicular was equally uncertain, its range differing sometimes 200 yards. Because the force of gravity acts always in a direction perpendicular to the horizon, a body projected in any direction would, if unresisted, be always in the same vertical plane, which, as is observed, is not true in fact; and, therefore, besides the resistance of the air to the progressive motion of the body, there is another lateral force producing a deviation from the plane of projection, which is probably the inequality of resistance upon its surface: for, if the surface of a projectile, protruding the particles of a fluid, be equally resisted in every point, it is obvious that the plane in which it first moves is not altered. But supposing the resistances upon different sides of a body to be unequal, it will be impelled, by the greater resistance, towards those parts where the resistance is least; and thus either diverge laterally from the first plane of motion, or ascend, or descend in that plane, beyond the altitude investigated by the theory. A motion of rotation round an axis will produce this inequality of resistance upon the different sides of a projectile, and consequently its irregular motions; for the parts of a revolving body are exposed to the air, which is protruded, in different angles of obliquity and with different forces, and must therefore be differently resisted. If the axis of rotation be perpendicular to the direction, the body will be deflected naturally from the vertical plane; and because the velocity of the body, and consequently resistance of the air, are variable, this deflection will be different in different points, or it will be a curve line. That this lateral deviation is produced by a motion of rotation seems to be confirmed by experiments. Let a sphere of wood, suspended freely by a cord in a

current

current of air, be made to revolve round a string as an axis, and the parts of its surface, on opposite sides, opposing or conspiring with the motion of the particles of the fluid, are unequally impelled by it, and the sphere is always deflected towards that side whose resistance is least. Or, let a wooden ball, loaded with lead, be suspended freely in a stream of water by a twisted cord; and, as the cord returns to its natural state, the ball revolves round it, and moves gradually towards the side where the resistance is least, or where the parts conspire with the motion of the water. When the ball arrives at its utmost extent, it is quiescent for a moment, and returns gradually to its first situation, and is again quiescent till the motion of the ball twist the cord the contrary way, and it then moves also towards the other side; and in this manner the ball continues to vibrate till its motion be destroyed by friction, and it remain quiescent in its first situation. The magnitude and direction of resistance of the air, and consequent deviations from theory, can only be ascertained by a series of experiments, which, resulting from the operation of causes apparently fluctuating and uncertain, produce an inconsistency in the same experiments, and render the subject at least very complicated and difficult *.

* See Robins's *Traacts of Gunnery*, p. 183, &c. and Euler's *True Principles of Gunnery*, by Mr. Brown.



ERRATA ET ADDENDA.

Pag. 2. lin. 7. *for of that, read that.*

10. art. 15. *for the existence of ratios, read the existence of the measures of ratios.*

24. art. 48. ex. 1. *for LA, read Lm, and for La, read LA.*

29. ex. 2. *dele or closeness.*

37. art. 78. *for all infinitely great, &c. read all infinitely great or small magnitudes compared with each other, admit of the same inequality of ratios with finite magnitudes.*

42. art. 85. *dele the remoteness of the elementary particles.*

67. lin. 8. *after velocity, read when in motion.*

fig. 26. *the letters L, M, wanting.*

81. art. 194. *for MBN, read PBN.*

82. art. 200. *after BD, put BE.*

83. art. 204. *dele in the same direction.*

fig. 48. *the letter D wanting.*

111. ex. 3. *for fire, read fibre.*

115. line the last. *after velocities, put in opposite directions.*

148. *for fig. 16, read fig. 106.*

152. l. . *by the direction of the resistances is understood the right line in which the resisting forces act, and by the quantities of the resistances are understood, not the quantities really exerted upon the wedge, but which would be exerted if no part were lost.*

186. *the following corollary wanting to art. 470. If the circle does not cut or touch the line DH, the bodies will never meet.*

197. art. 500. *for \pm velocity of $V \times SA$, read $= \frac{\text{velocity of } V \times SA}{SV}$.*

198. art. 500. l. 2. *for $m \times SV$, read $m \times SV \times SC$.*

231. art. 589. *for arc, read chord, in which corollary, the body is supposed to vibrate in the chords.*

